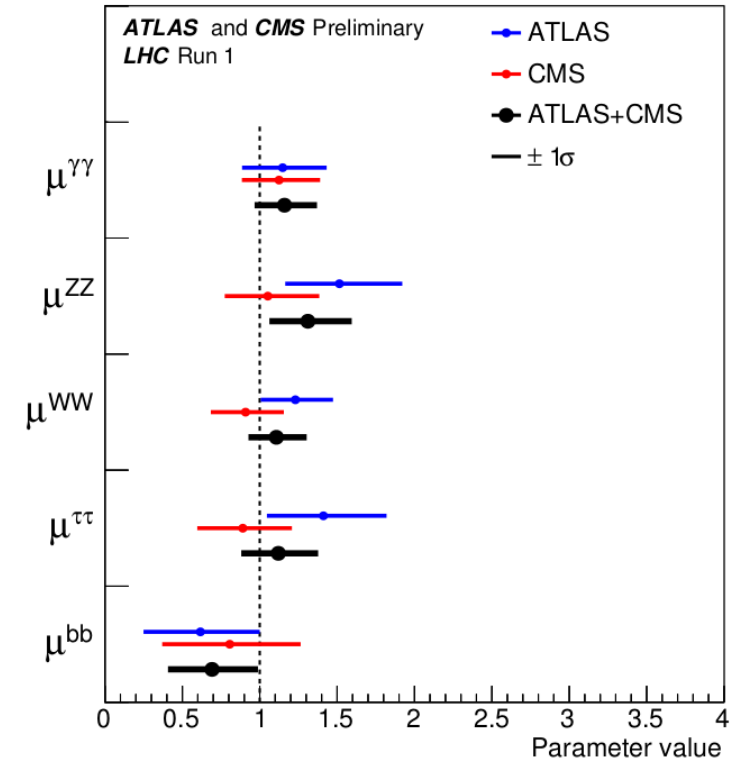
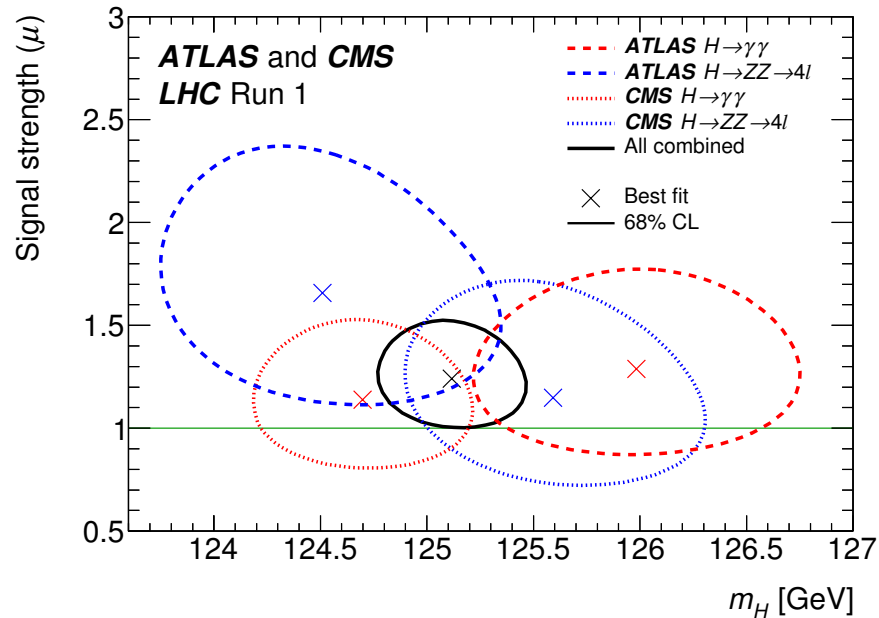


Supersymmetric models in view of recent LHC data

Werner Porod

Universität Würzburg

- Higgs discovery and LHC BSM results: implications
- 'Natural' SUSY: MSSM and extensions



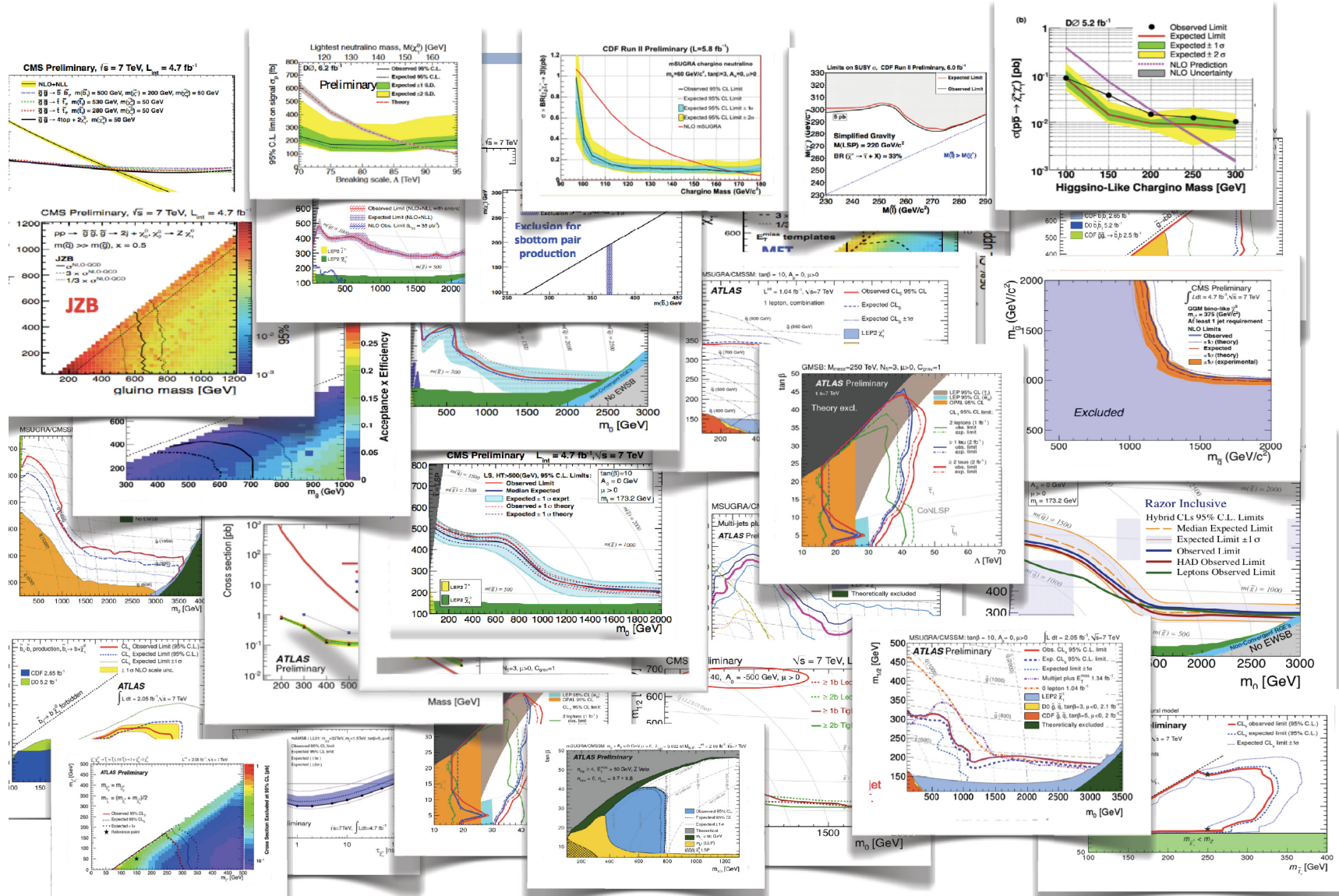
$m_H = 125.09 \pm 0.21$ (stat) ± 0.11 (sys) GeV
run 1, PRL 114 (2015) 191803

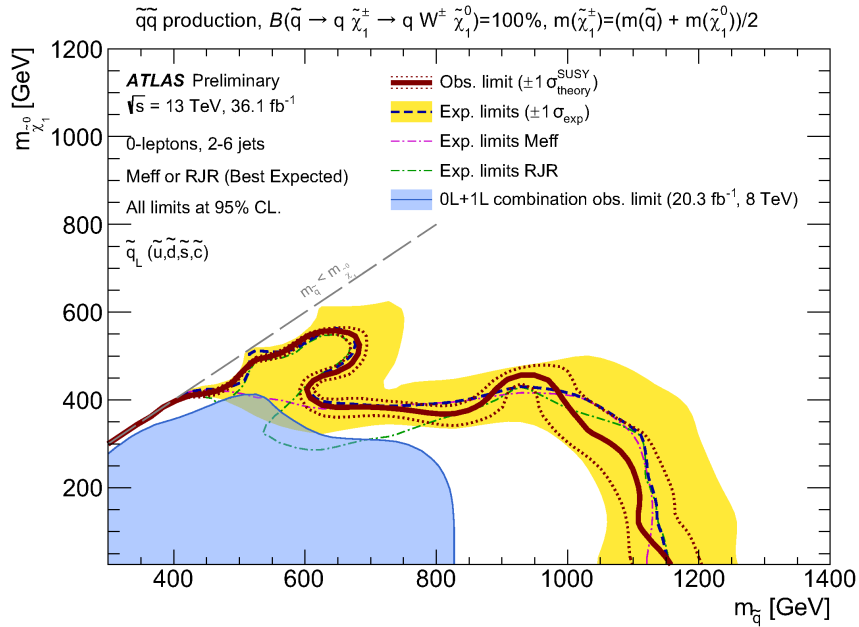
$m_H = 125.25 \pm 0.20$ (stat) ± 0.08 (sys) GeV
run 2, see talk by R. Nicolaidou

ATLAS-CONF-2015-044

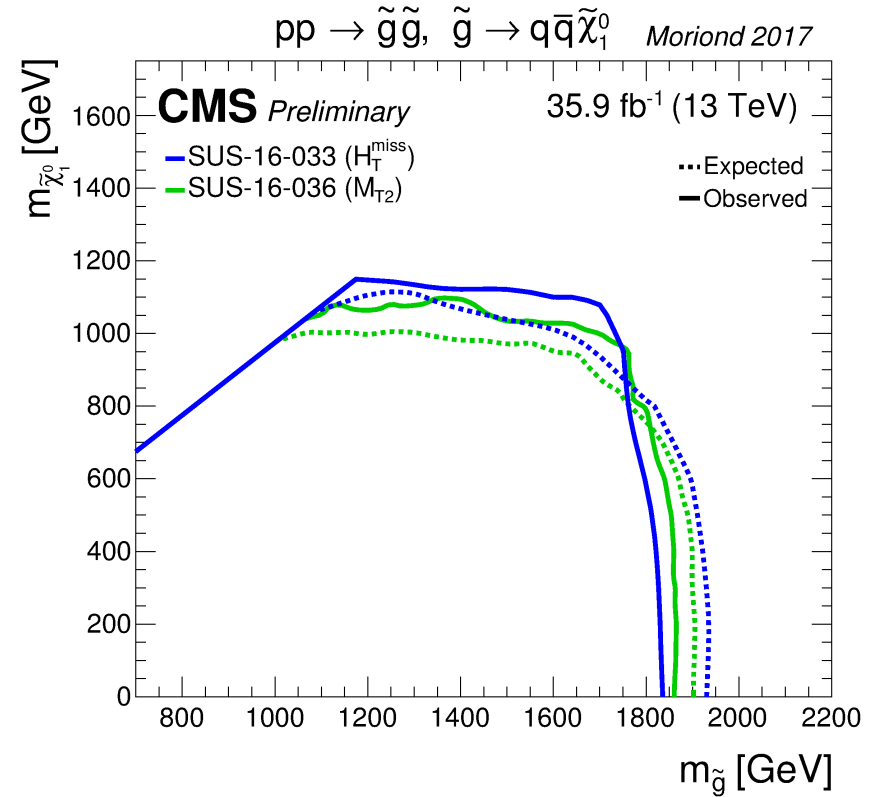
CMS-PAS-HIG-15-002

$$(125 \text{ GeV})^2 \simeq m_Z^2 + (86 \text{ GeV})^2 \Rightarrow \text{large corrections within MSSM}$$





ATLAS-CONF-2017-022

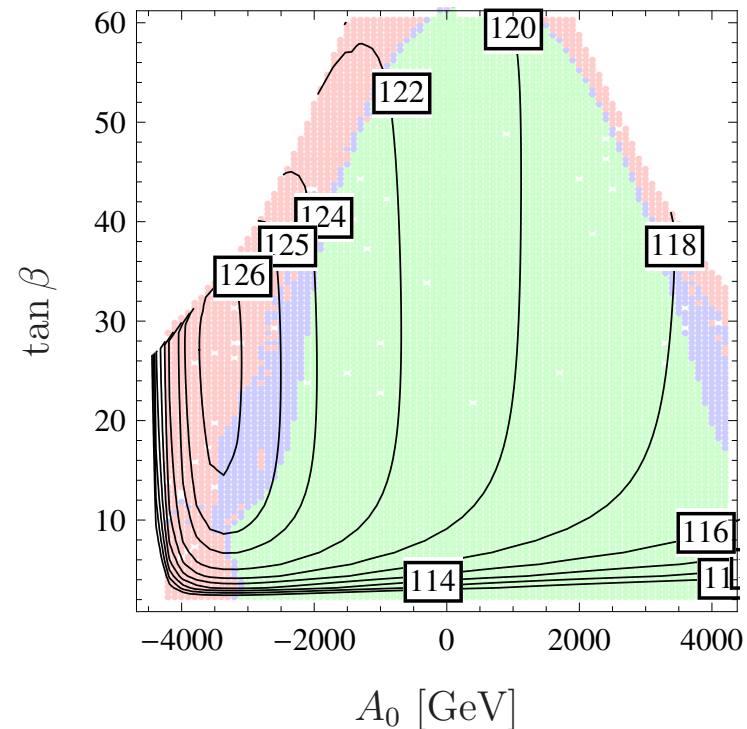
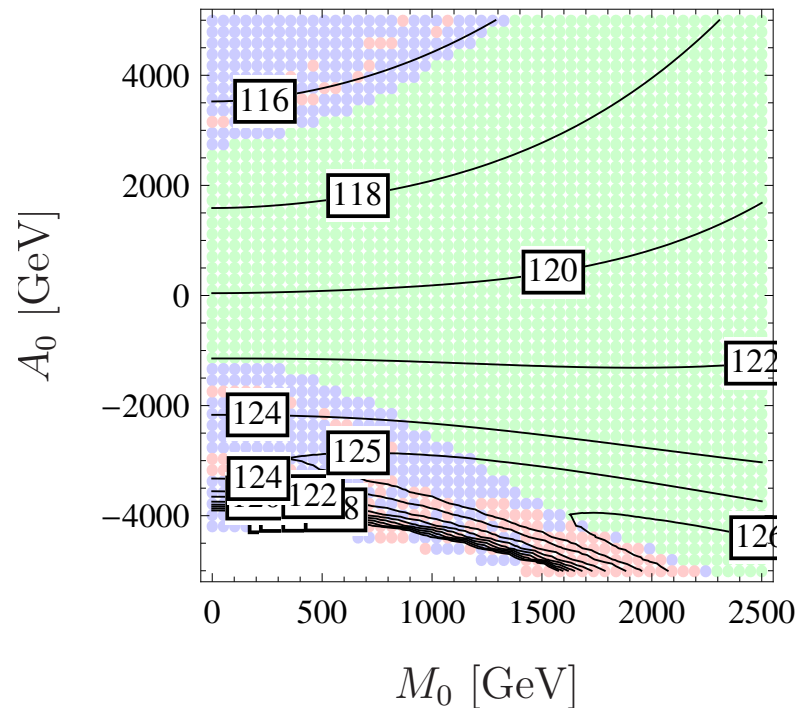


see talks by Simone Amoroso & Rishi Gautam Patel

$m_h = 125.2 \text{ GeV} \Rightarrow$ large loop contributions
 \Rightarrow heavy stops and/or large left-right mixing for stops

- GMSB: $m_{\tilde{t}_1} \gtrsim 6 \text{ TeV}$,
 M. A. Ajaib, I. Gogoladze, F. Nasir, Q. Shafi, arXiv:1204.2856
 more complicated models based on P. Meade, N. Seiberg and D. Shih,
 arXiv:0801.3278 \Rightarrow allow additional terms
 e.g. S. Knappen, D. Redigolo, arXiv:1606.07501 $m_{\tilde{t}_1} \simeq m_{\tilde{b}_1} \gtrsim 1 \text{ TeV}$ if
 $M_{\text{mess}} \gtrsim 10^{15} \text{ GeV}$
- CMSSM, NUHM models: $|A_0| \simeq 2m_0$,
 H. Baer, V. Barger and A. Mustafayev, arXiv:1112.3017; M. Kadastik *et al.*,
 arXiv:1112.3647; O. Buchmueller *et al.*, arXiv:1112.3564; J. Cao, Z. Heng, D. Li,
 J. M. Yang, arXiv:1112.4391; L. Aparicio, D. G. Cerdeno, L. E. Ibanez,
 arXiv:1202.0822; J. Ellis, K. A. Olive, arXiv:1202.3262; ...
 CMSSM fit to data P. Bechtle *et al.*, arXiv:1508.05951: best fit point with
 $m_{\tilde{g}}, m_{\tilde{q}} \gtrsim 2 \text{ TeV}$, $m_{\tilde{l}_R} \simeq 600 \text{ GeV}$, $m_{\tilde{\chi}_1^0} \simeq 450 \text{ GeV}$
- general high scale models: $A_0 \simeq -(1-3) \max(M_{1/2}, m_{Q_3}, m_{U_3}) @ M_{GUT}$
 among other cases, details in F. Brümmer, S. Kraml and S. Kulkarni, arXiv:1204.5977

- SUSY models contain many scalars \Rightarrow complicated potential
- usually some parameters (μ, B) are chosen to obtain correct EWSB
- does not exclude the existence of other minima breaking charge and/or color!



$$M_{1/2} = 1 \text{ TeV}, \tan \beta = 10, \mu > 0$$

$$M_{1/2} = M_0 = 1 \text{ TeV}$$

J.E. Camargo-Molina, B. O'Leary, W.P., F. Staub, arXiv:1309.7212

several studies, see e.g. S. Sekmen et al., arXiv:1109.5119; A. Arbey, M. Battaglia, A. Djouadi and F. Mahmoudi, arXiv:1211.4004; M. Cahill-Rowley, J. Hewett, A. Ismail and T. Rizzo, arXiv:1308.0297

- generic signatures are well known: multi-lepton, multi-jets + missing E_T

- sub-class of general MSSM: ‘natural SUSY’

see e.g. M. Papucci, J. T. Ruderman and A. Weiler, arXiv:1110.6926;

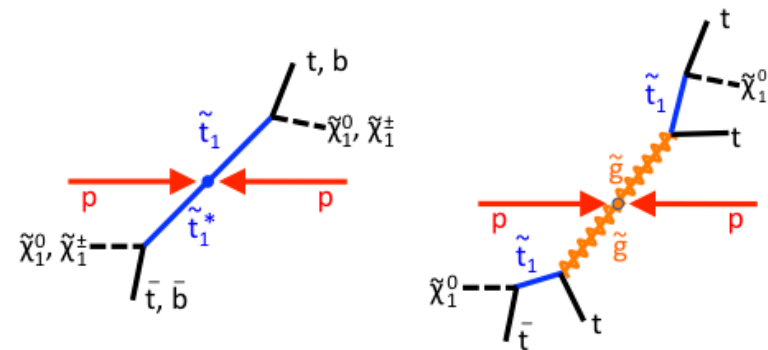
H. Baer, V. Barger, P. Huang, A. Mustafayev, X. Tata, arXiv:1207.3343

keep only SUSY particles light needed for ‘natural Higgs’:

$$\tilde{t}_1, \tilde{b}_1, \tilde{g}, \tilde{\chi}_{1,2}^0 \simeq \tilde{h}_{1,2}^0, \tilde{\chi}_1^+ \simeq \tilde{h}^+$$

$$\Rightarrow 100 \text{ MeV} \lesssim m_{\tilde{\chi}_1^+} - m_{\tilde{\chi}_1^0} \simeq m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0} \lesssim 5 - 10 \text{ GeV}$$

$$\begin{aligned} \tilde{g} &\rightarrow \tilde{t}_1 t, \tilde{b}_1 b \\ \tilde{t}_1 &\rightarrow t \tilde{\chi}_{1,2}^0, b \tilde{\chi}_1^+, W^+ \tilde{b}_1 \\ \tilde{b}_1 &\rightarrow b \tilde{\chi}_{1,2}^0, t \tilde{\chi}_1^-, W^- \tilde{t}_1 \end{aligned}$$



BRs depend on the nature of \tilde{t}_1 and \tilde{b}_1

Higgsino mass: $\mu + \mu'$ with soft SUSY breaking parameter: $\mathcal{L} = -\mu' \tilde{H}_d \tilde{H}_u$

- additional D-term contributions to m_h at tree-level

$$\text{extra } U(1)_\chi: m_{h,tree}^2 \leq m_Z^2 + \frac{1}{4}g_\chi^2 v^2$$

- Origin of R -parity $R_P = (-1)^{2s+3(B-L)}$

$$\Rightarrow SO(10) \rightarrow SU(3)_C \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

$$\rightarrow SU(3)_C \times SU(2)_L \times U(1)_R \times U(1)_{B-L}$$

$$\cong SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_\chi$$

$$\text{or } E(8) \times E(8) \rightarrow SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_{B-L}$$

- Neutrino masses

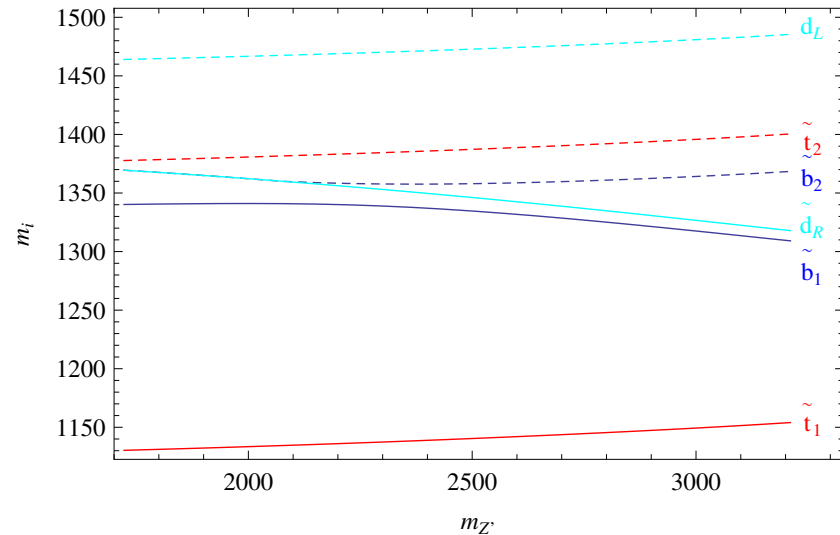
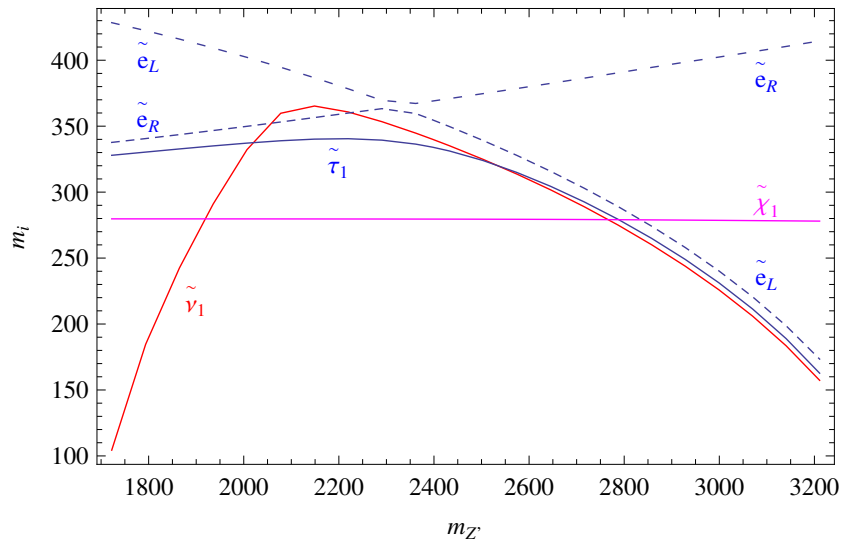
$B - L$ anomaly free $\Rightarrow \nu_R$

usual seesaw, inverse seesaw

$$M_{\tilde{l}}^2 = \begin{pmatrix} M_{\tilde{L}}^2 + D_L + m_l^2 & \frac{1}{\sqrt{2}} (v_d T_l - \mu Y_l v_u) \\ \frac{1}{\sqrt{2}} (v_d T_l - \mu Y_l v_u) & M_{\tilde{E}}^2 + D_R + m_l^2 \end{pmatrix},$$

$$D_L \simeq \left(-\frac{1}{2} + \sin^2_{\theta_W}\right) m_Z^2 c_{2\beta} - \frac{5}{4} m_{Z'}^2 c_{2\beta_R} \quad \text{and} \quad D_R \simeq -\sin^2_{\theta_W} m_Z^2 c_{2\beta} + \frac{5}{4} m_{Z'}^2 c_{2\beta_R}$$

neglecting gauge kinetic effects; similarly for squarks



$$m_0 = 100 \text{ GeV}, m_{1/2} = 700 \text{ GeV}, A_0 = 0, \tan \beta = 10, \mu > 0$$

$$\tan \beta_R = 0.94, m_{A_R} = 2 \text{ TeV}, \mu_R = -800 \text{ GeV}$$

effective model with $\tilde{t}_1, \tilde{b}_1, \tilde{h}_{1,2}^0, \tilde{h}^+, \tilde{\nu}_R$

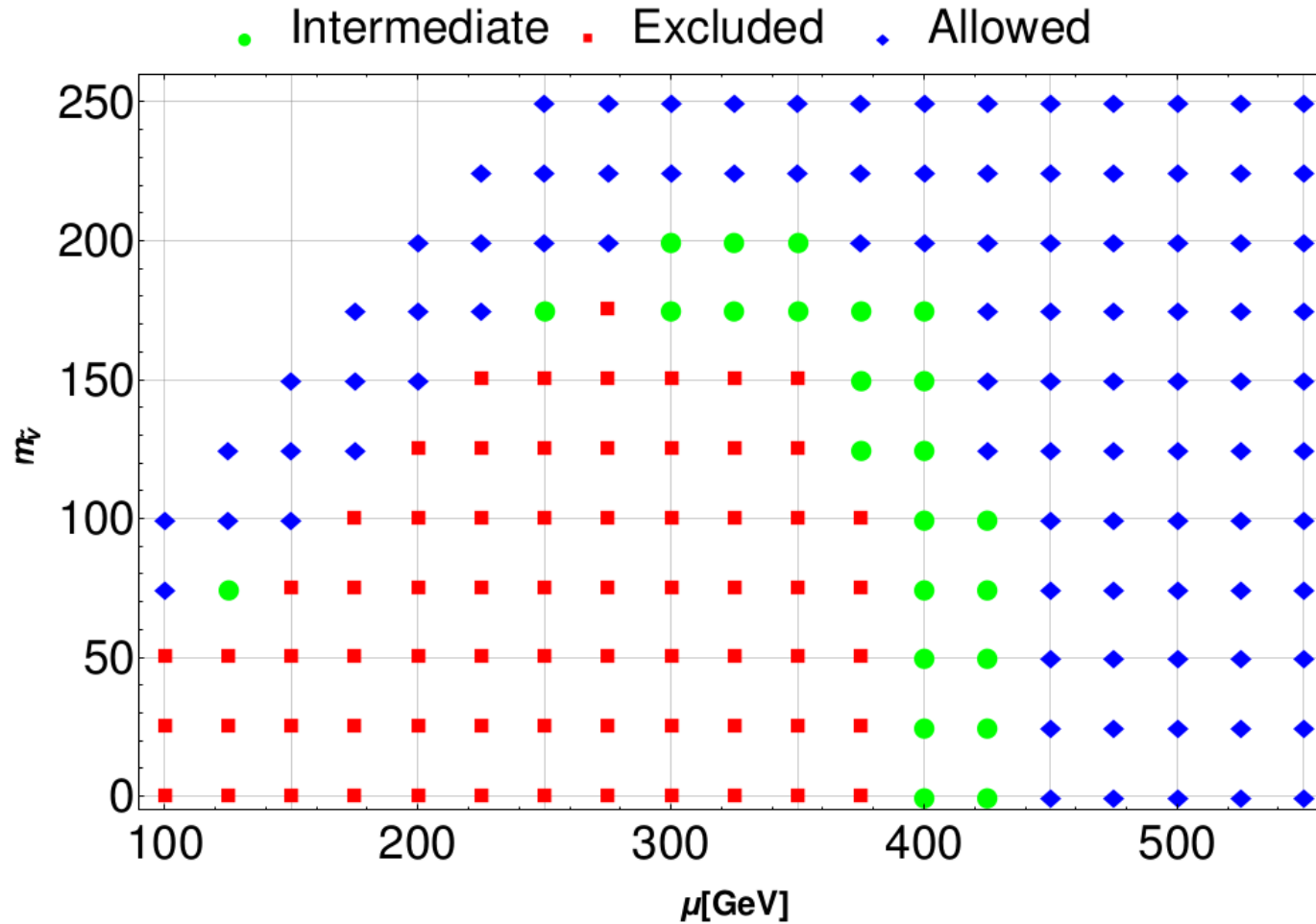
- $m_{\tilde{t}_1}$ in GeV: 300, 400, 500, 600, 700, 800, 900, 1000
- $m_{\tilde{b}_1}$ in GeV: 300, 400, 500, 600, 700, 800, 900, 1000
- $m_{\tilde{\nu}_R}$ in GeV : 60, 100, 200, 300, 400, 500
- μ in GeV: 110, 190, 290, 390, 490, 590 and require $m_{\tilde{\nu}_R} < \mu$
- $\tan \beta$: 10, 50
- $\theta_{\tilde{t}}$: $0^\circ, 45^\circ, 90^\circ$
- $\theta_{\tilde{b}}$: $0^\circ, 45^\circ, 90^\circ$
- $M_1 = M_2 = 1$ TeV
- everything else, including \tilde{t}_2 , and $m_{\tilde{g}}$: 2 TeV, \tilde{b}_2 calculated

$$m_W^2 \cos 2\beta = m_{\tilde{t}_1}^2 \cos^2 \theta_{\tilde{t}} - m_{\tilde{t}_2}^2 \sin^2 \theta_{\tilde{t}} - m_{\tilde{b}_1}^2 \cos^2 \theta_{\tilde{b}} - m_{\tilde{b}_2}^2 \sin^2 \theta_{\tilde{b}} - m_t^2 + m_b^2$$

$$\Rightarrow m_{\tilde{b}_2} \leftrightarrow m_{\tilde{b}_1} \text{ if necessary}$$

$$m_{\tilde{t}_2} \leftrightarrow m_{\tilde{t}_1} \text{ (if } \cos \theta_{\tilde{b}} = 1)$$

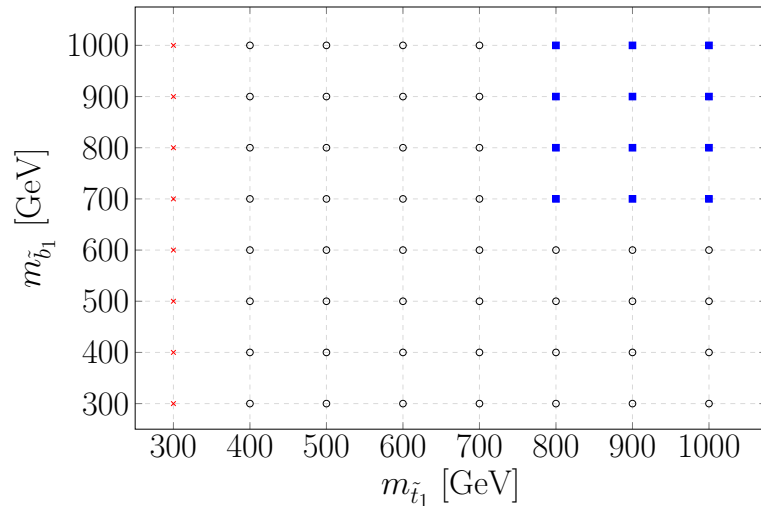
$$pp \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow \ell^+ \ell^- \tilde{\nu}_R \tilde{\nu}_R^*$$



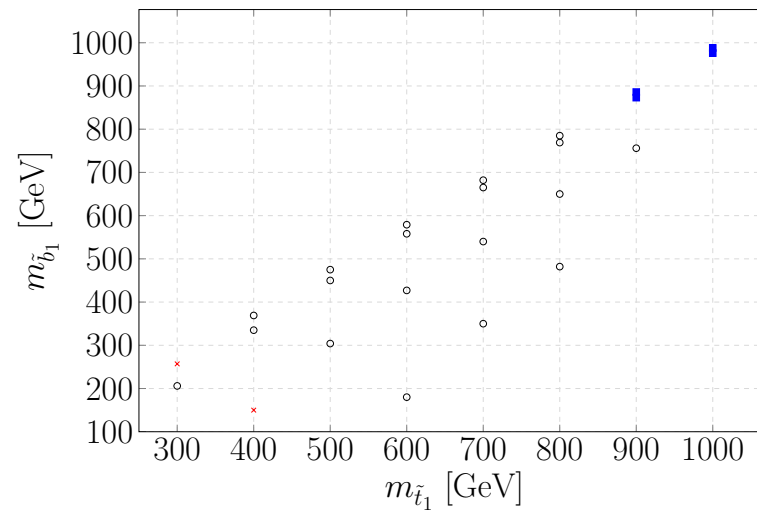
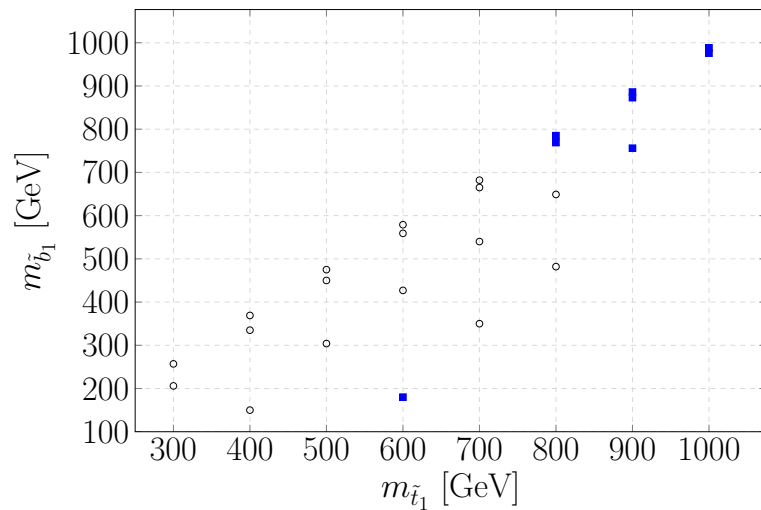
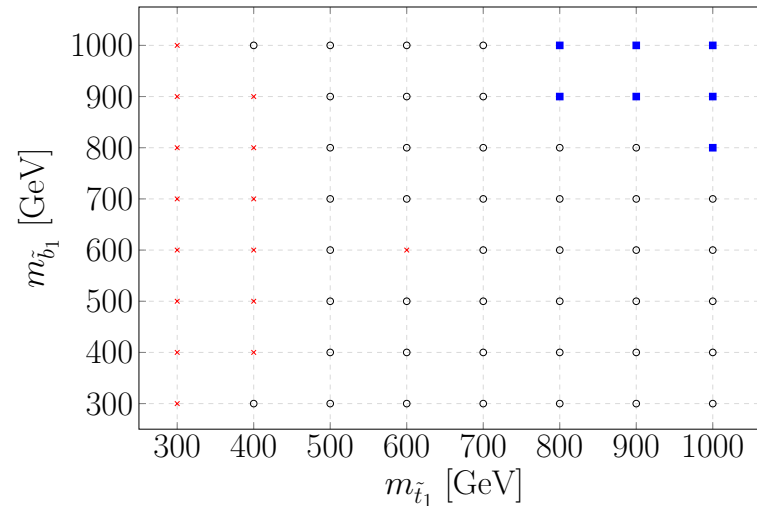
using CheckMATE 2.0

Th. Faber, J. Jones, Nh. Cerna-Velazco, WP arXiv:1704.xxxxx

ambiguous as allowed



ambiguous as forbidden



× excluded for all parameters considered, ○ exclusion depends on parameters

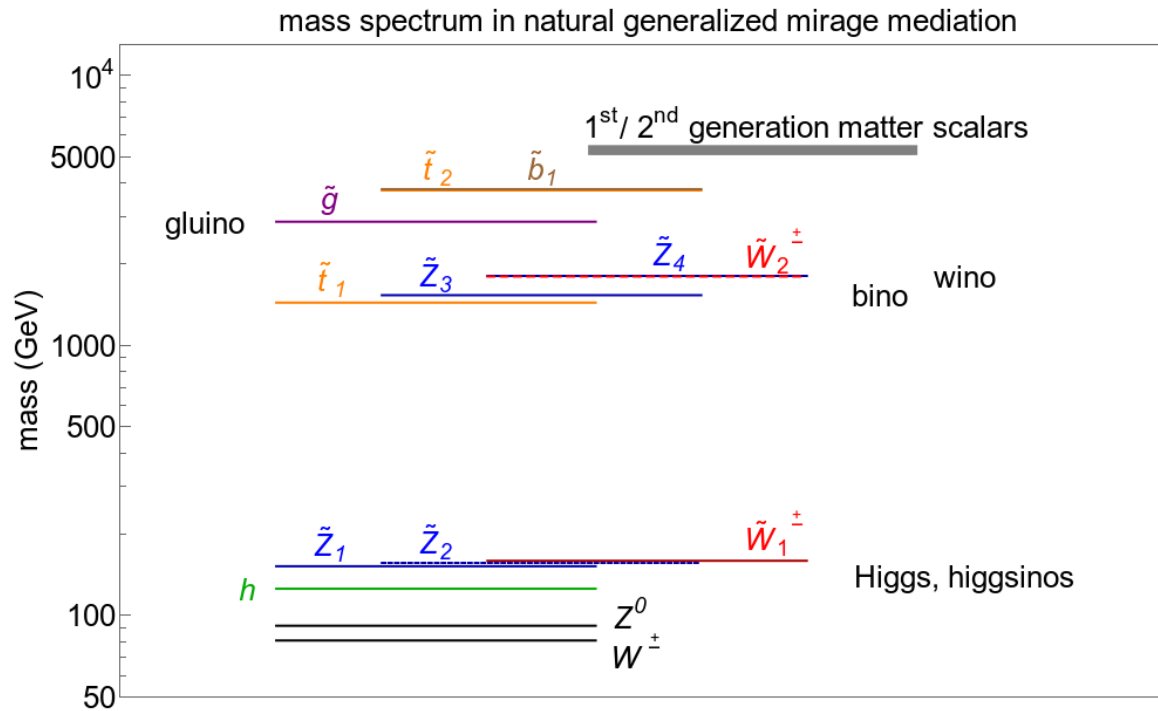
■ allowed for all parameters, using CheckMATE 1.0

L. Mitzka, WP arXiv:1603.06130

- LHC: $m_h \simeq 125$ GeV, no conclusive BSM physics found \Rightarrow
 - GMSB, CMSSM, NUHM: $m_{\tilde{g}}, m_{\tilde{q}} \gtrsim 2$ TeV
 - CMSSM, NUHM: large A_0 , danger of color and charge breaking minima
- general MSSM: SUSY particles with masses of few 100 GeV still allowed if spectra compressed, in particular light \tilde{t}_1 still allowed
- ‘Natural SUSY’: take only those states light which contribute to EWSB: $\tilde{h}^{0,\pm}, \tilde{t}_1, \tilde{g}, \tilde{b}_i$
disadvantage: cannot explain dark matter relic density
- ‘Natural SUSY’ + $\tilde{\nu}_R$
 - $\tilde{\nu}_R$ LSP: compatible with DM, no direct DM constraint apply
 - $m_{\tilde{h}^+} \lesssim 400$ GeV excluded if $m_{\tilde{h}^+} - m_{\tilde{\nu}_R} \gtrsim 150$ GeV
 - independent of other parameters: $m_{\tilde{t}_1} \lesssim 300$ GeV excluded except for very compressed spectra (8 TeV data only)
 - for 300 GeV $\lesssim m_{\tilde{t}_1} \lesssim 800$ GeV: exclusion depends on parameters, in particular on $\cos \theta_{\tilde{t}}$ (8 TeV data only)

Different sources for soft SUSY breaking: moduli & AMSB

main consequence: gaugino masses unify at a (vastly) different scale than gauge couplings

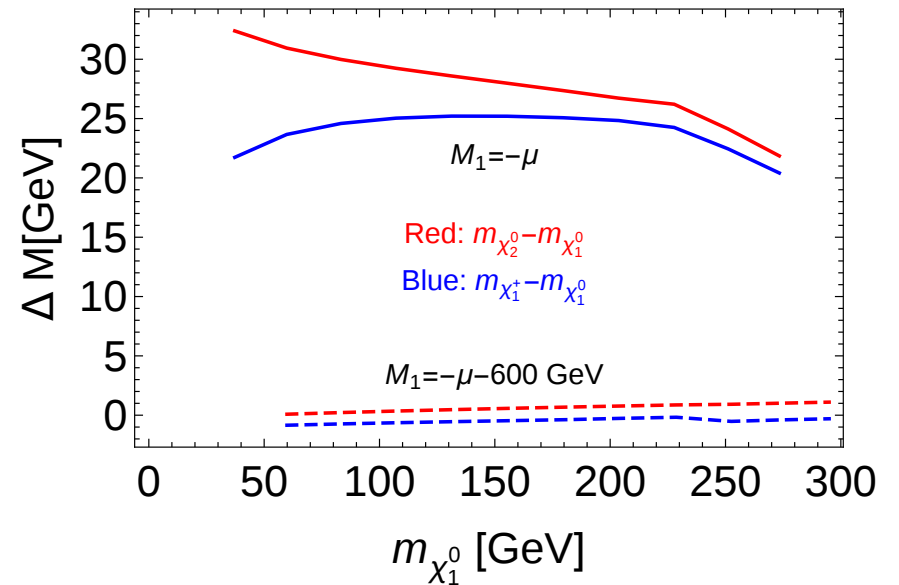
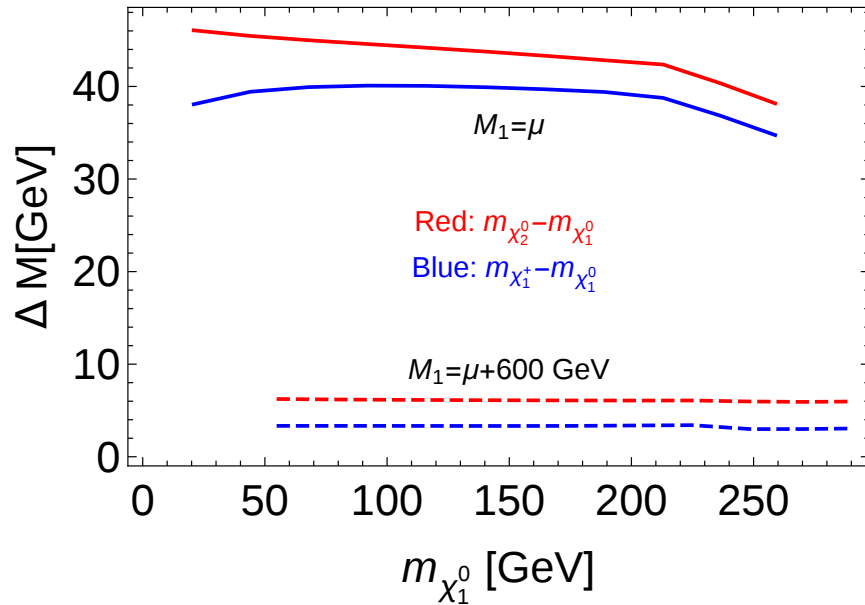


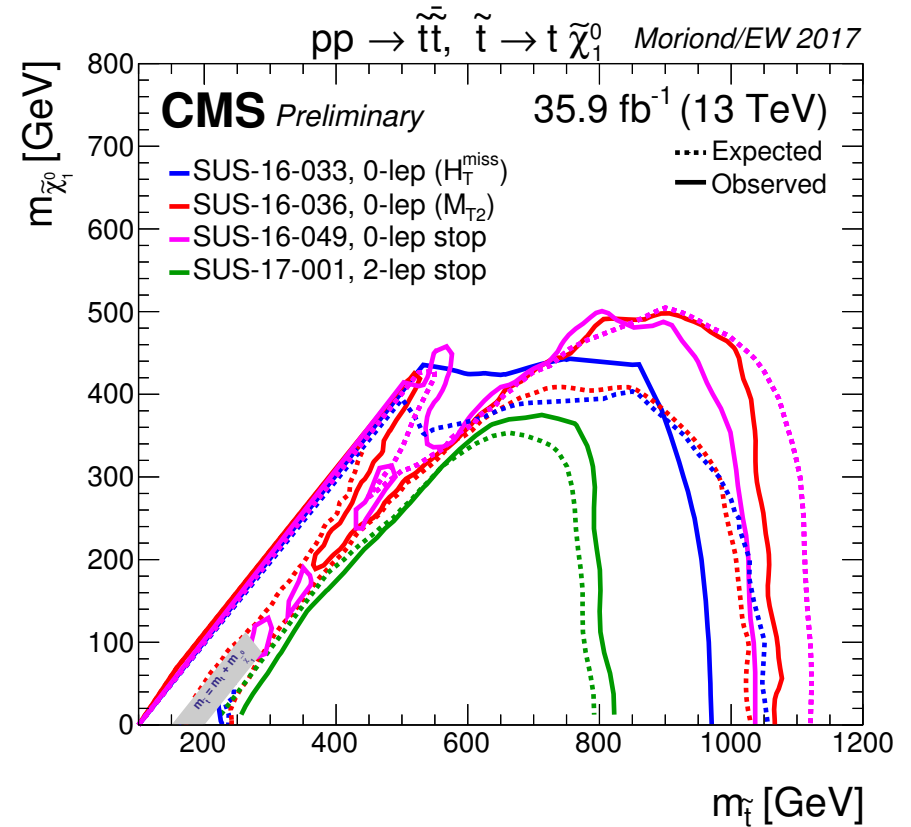
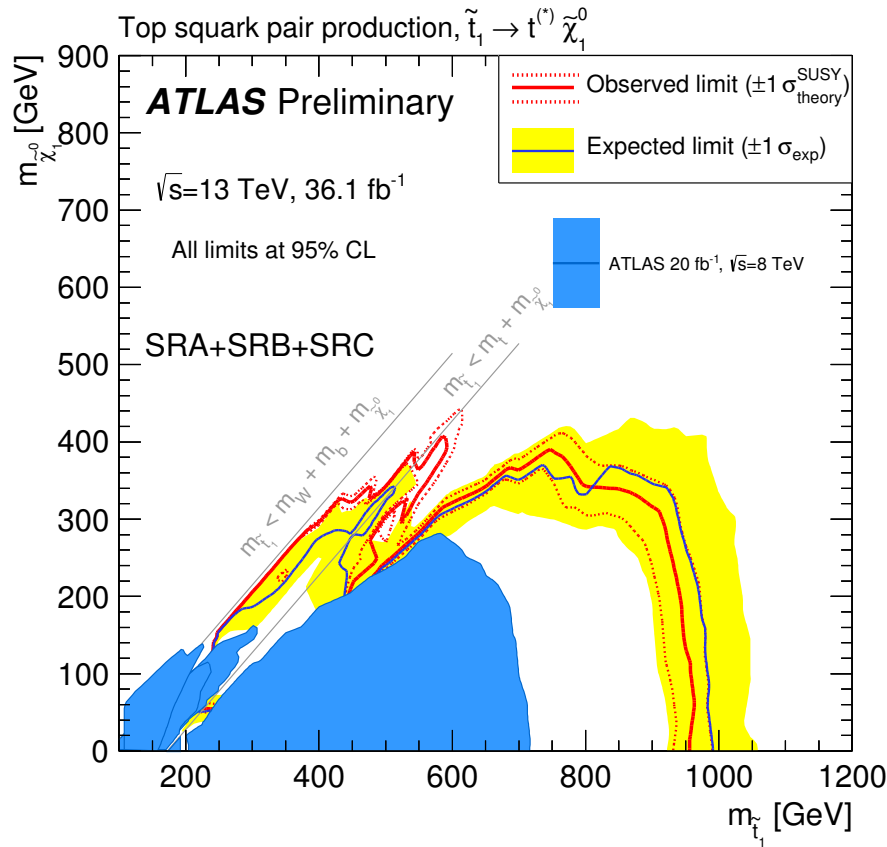
H. Baer, V. Barger, H. Serce and X. Tata, arXiv:1610.06205

limit $|\mu| \ll |M_1|, |M_2|$:

$$\Delta m_0 = m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0} \simeq m_Z^2 \left(\frac{s_\omega^2}{M_1} + \frac{c_\omega^2}{M_2} \right)$$

$$\Delta m_\pm = m_{\tilde{\chi}_1^\pm} - m_{\tilde{\chi}_1^0} \simeq \frac{\Delta m_0}{2} + |\mu| \frac{\alpha(m_Z)}{\pi} \left(2 + \ln \frac{m_Z^2}{\mu^2} \right)$$

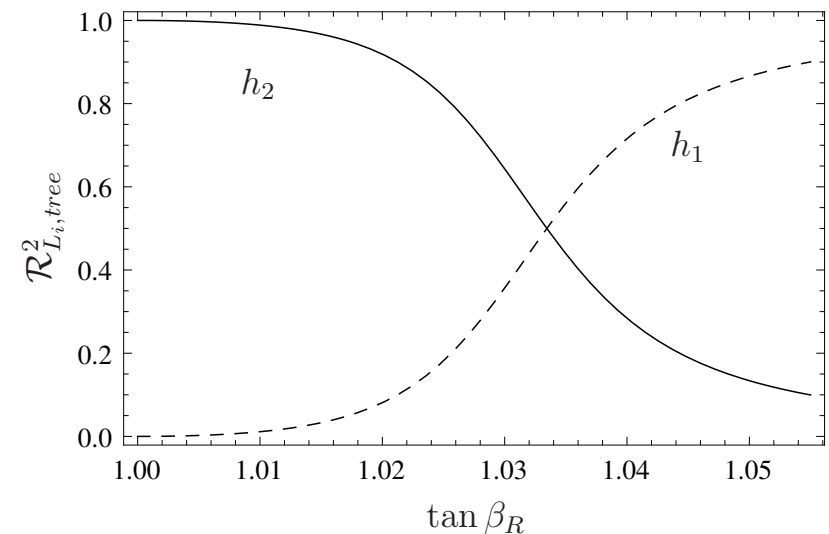
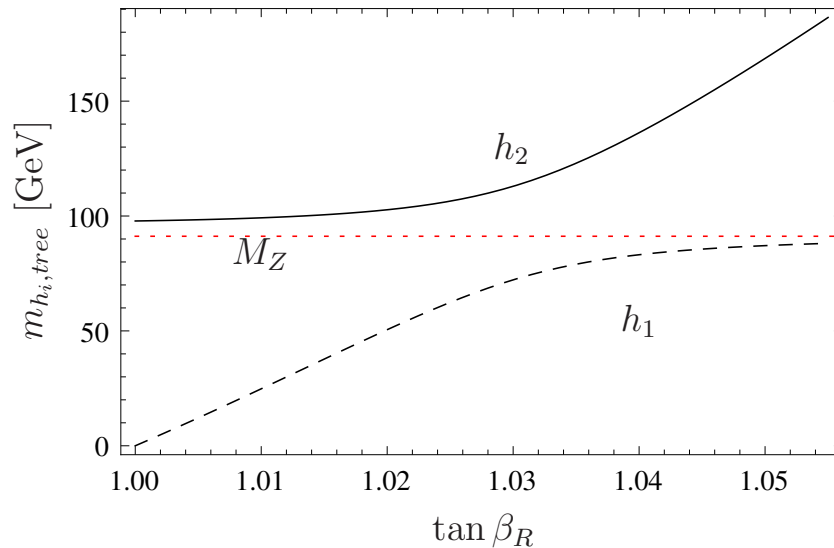




$$\frac{d}{dt} \begin{pmatrix} m_{H_u}^2 \\ m_{\tilde{t}_R}^2 \\ m_{\tilde{Q}_L^3}^2 \end{pmatrix} = -\frac{8\alpha_s}{3\pi} M_3^2 \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix} + \frac{Y_t^2}{8\pi^2} \left(m_{\tilde{Q}_L^3}^2 + m_{\tilde{t}_R}^2 + m_{H_u}^2 + A_t^2 \right) \begin{pmatrix} 3 \\ 2 \\ 1 \end{pmatrix}$$

extra $U(1)_\chi$ with new D-term contributions at tree-level: $m_{h_i,tree}^2 \leq m_Z^2 + \frac{1}{4}g_\chi^2 v^2$

H.E. Haber, M. Sher, PRD 35 (1987) 2206; M. Drees, PRD 35 (1987) 2910; M. Cvetič et al., hep-ph/9703317; E. Ma, arXiv:1108.4029; M. Hirsch et al., arXiv:1110.3037



$n = 1$, $\Lambda = 5 \cdot 10^5$ GeV, $M = 10^{11}$ GeV, $\tan \beta = 30$, $\text{sign}(\mu_R) = -$, $\text{diag}(Y_S) = (0.7, 0.6, 0.6)$, $Y_\nu^{ii} = 0.01$, $v_R = 7$ TeV

M.E. Krauss, W.P., F. Staub, arXiv:1304.0769

$$\mathcal{W}_{eff} = \mu \hat{H}_u \cdot \hat{H}_d + Y_t \hat{t}_R \hat{H}_u \cdot \hat{Q} + Y_b \hat{b}_R \hat{Q} \cdot \hat{H}_d + \sum_k \left(Y_{\nu,k} \hat{\nu}_{R,k} \hat{H}_u \cdot \hat{L}_k + M_k \hat{S}_k \hat{\nu}_{R,k} \right) ,$$

$$\mathcal{V}^{soft} = \frac{1}{2} M_3 \tilde{g} \tilde{g} + \sum_S m_S^2 |S|^2 + B_\mu H_u \cdot H_d + \sum_k \left(B_{M_k} \tilde{S}_k \tilde{\nu}_{R,k} + T_{\nu k} \tilde{\nu}_{R,k} \tilde{H}_u \cdot \tilde{L}_k \right) \\ + T_t \tilde{t}_R H_u \cdot \tilde{Q} + T_b \tilde{b}_R \tilde{Q} \cdot H_d$$

$$S = H_u, H_d, \tilde{Q}, \tilde{t}_R, \tilde{b}_R, \tilde{\nu}_R$$

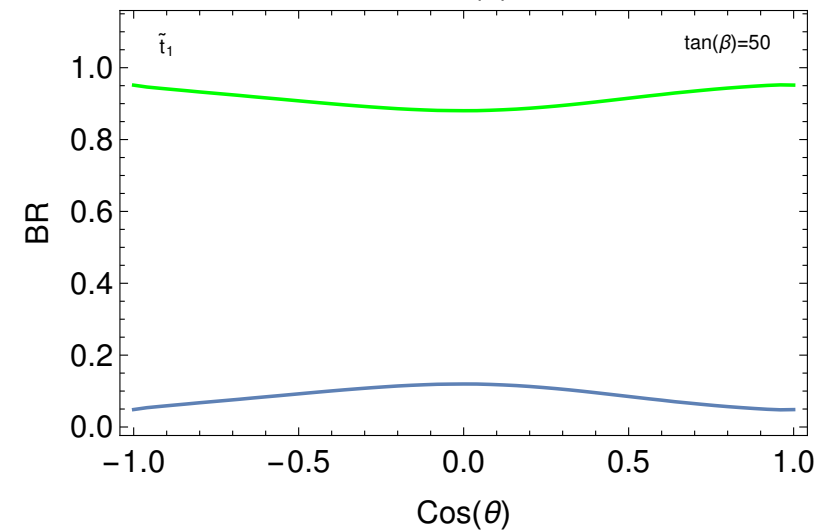
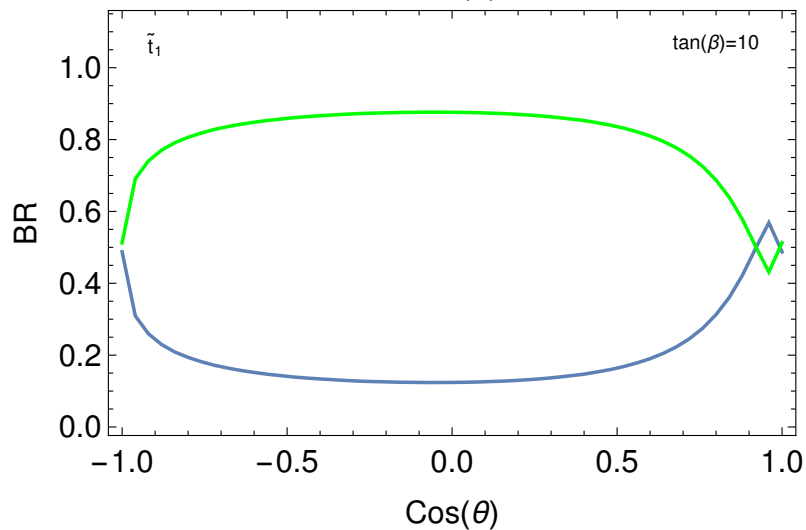
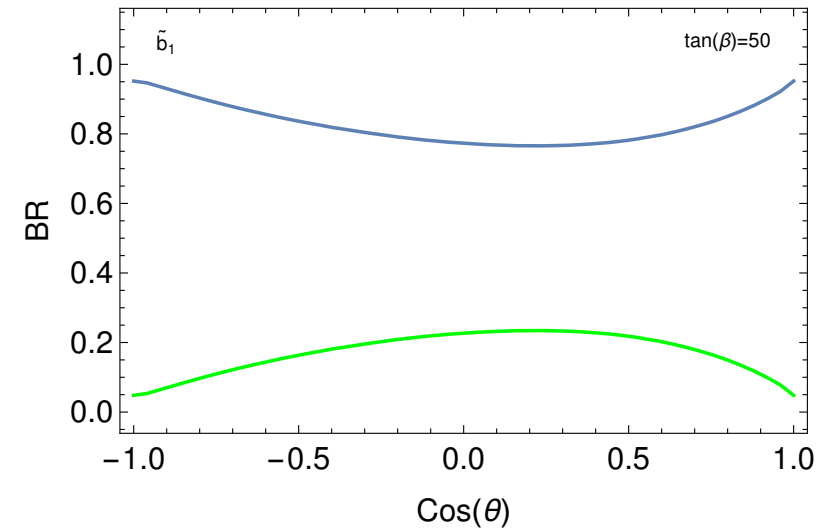
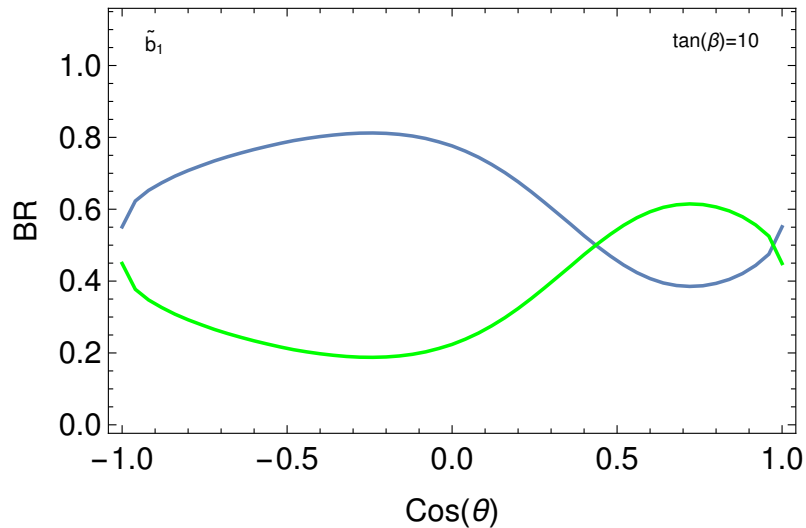
assume $Y_{\nu,k} = Y_\nu$; tree level relation

$$m_W^2 \cos 2\beta = m_{\tilde{t}_1}^2 \cos^2 \theta_{\tilde{t}} + m_{\tilde{t}_2}^2 \sin^2 \theta_{\tilde{t}} - m_{\tilde{b}_1}^2 \cos^2 \theta_{\tilde{b}} - m_{\tilde{b}_2}^2 \sin^2 \theta_{\tilde{b}} - m_t^2 + m_b^2$$

simplified $\tilde{\nu}_R, \tilde{S}$ mass matrix (one generation):

$$M_{\nu_R, \tilde{S}}^2 = \begin{pmatrix} |M_k|^2 & B_{M_k} \\ B_{M_k} & |M_k|^2 \end{pmatrix} \Rightarrow m_{1,2}^2 = |M_k|^2 \pm |B_{M_k}|$$

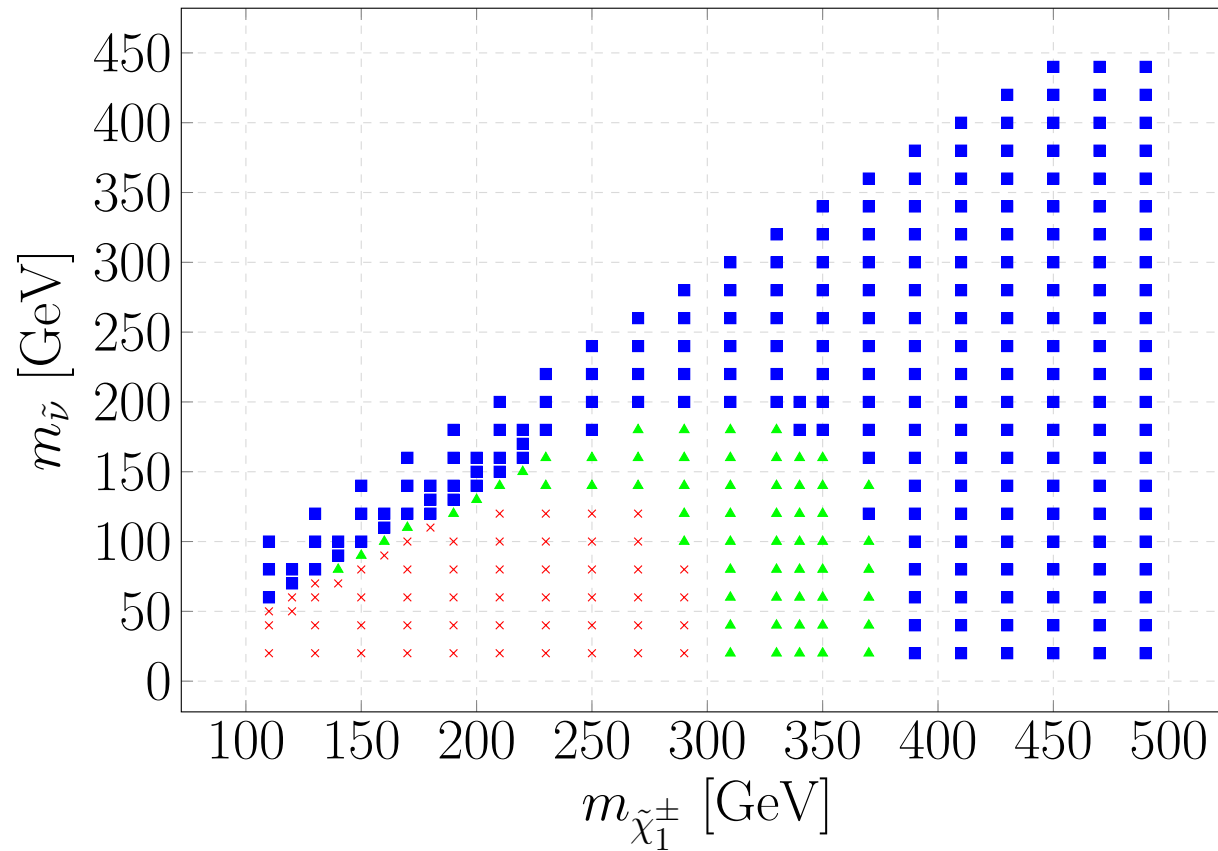
\Rightarrow expect lightest 'sneutrino' as LSP,



$m_{\tilde{q}_1} = 500 \text{ GeV}$ ($q = b, t$), $m_{\tilde{\nu}_R} = 100 \text{ GeV}$, $\mu = 590 \text{ GeV}$, $M_1 = M_2 = 1 \text{ TeV}$. blue line:
 $\tilde{q}_1 \rightarrow q\nu\tilde{\nu}_R$, green line $\tilde{q}_1 \rightarrow q'l\tilde{\nu}_R$ summing over l ; L. Mitzka, WP arXiv:1603.06130

atlas_1403_2500	\tilde{g} and \tilde{q}	jets, 2SS/3 leptons
atlas_conf_2013_036	RPV & RPC SUSY	four or more leptons
atlas_1402_7029	$\tilde{\chi}^\pm$ and $\tilde{\chi}^0$	3 leptons and E_T^{miss}
atlas_1403_4853	\tilde{t}	two leptons and 2 b jets
atlas_1403_5294	$\tilde{\ell}, \tilde{\chi}^{0,\pm}$	two leptons and E_T^{miss}
atlas_conf_089	\tilde{t}	two leptons via the razor variable
atlas_conf_2013_049	$\tilde{\chi}^{0,\pm}, \tilde{\ell}$	two leptons
atlas_conf_2013_014	\tilde{t}	2 b jets, two leptons (via τ), E_T^{miss}
atlas_1407_0583	\tilde{t}	1 lepton, jets and E_T^{miss}
atlas_conf_2013_062	\tilde{t}, \tilde{g}	1 lepton, jets and E_T^{miss}
atlas_conf_2013_104	\tilde{t}	1 lepton, jets and E_T^{miss}
atlas_conf_2013_061	\tilde{g}	three b -jets and E_T^{miss}
atlas_1308_2631	\tilde{b}, \tilde{t}	2 b jets and E_T^{miss}
atlas_conf_2013_047	\tilde{q}, \tilde{g}	jets and E_T^{miss}
atlas_conf_2013_024	\tilde{t}	hadronic $t\bar{t}$ final states

$$pp \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow \ell^+ \ell^- \tilde{\nu}_R \tilde{\nu}_R^*$$



× excluded, ▲ ambiguous, ■ allowed

using CheckMATE 1.0

L. Mitzka, WP arXiv:1603.06130