

# *Light particle dark matter and MeV gamma-ray observation*

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*Papers: JHEP01 (2023) 106, JHEP07 (2019) 050.*

- ✓ *MeV gamma-ray observations, e.g., COSI, are now being developed.*
- ✓ *These enable us to probe light particle dark matter candidates.*
- ✓ *CMB observation severely constrains such candidates, and it requests some discussions to be a target of the observations.*
- ✓ *We propose a framework to discuss the candidates at future MeV gamma-ray observations, using some concrete examples.*

# Dark matter problem & what we know

## ○ The dark matter problem:

*We know that dark matter (DM) exists in our universe.*

*We know how the DM is distributed in our universe.*

*We know little about the microscopic nature of the DM.*

## ○ What we know about the DM:

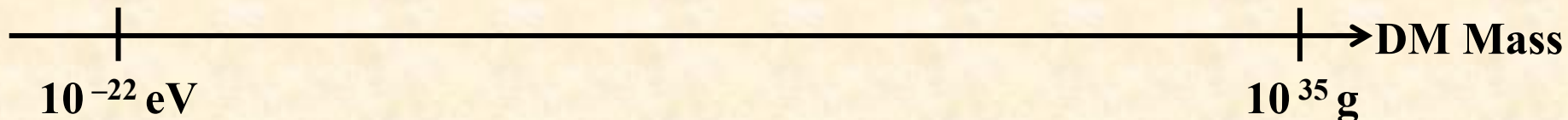
*The DM must be (almost) electrically neutral.*

*The DM must be (enough) stable. (Its lifetime  $\gg$  Age of U.)*

*The DM must be (enough) cold (non-relativistic) at present.*

*The DM must be (enough) weak-interacting.*

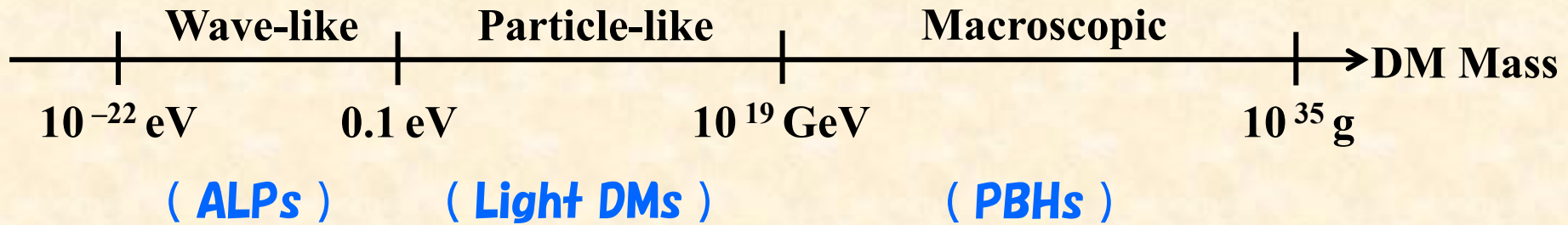
*The mass of the DM must be between  $10^{-22} \text{eV}$  and  $10^{35} \text{g}$ .*



*$m_{DM} > 10^{-22} \text{eV}$ :  $\lambda_D$  (De Broglie W. L.) =  $2\pi / (mv) <$  Galaxy size.*

*$m_{DM} < 10^{35} \text{g}$ : DM must be lighter enough than a host galaxy.*

# DM candidates & MeV- $\gamma$ observation



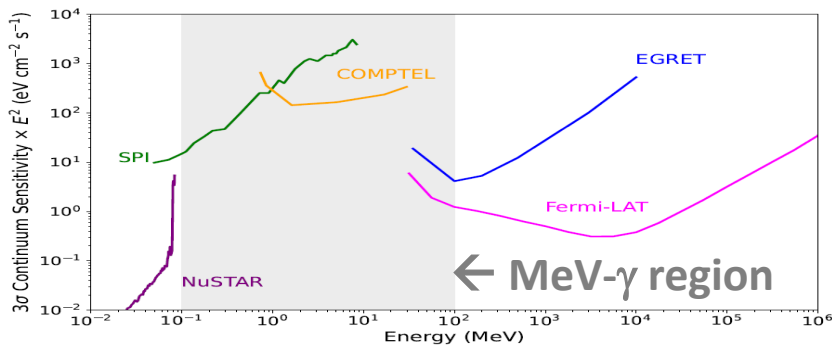
○ Three DM mass regions:

$m_{DM} < 10^{-1}$  eV: The occupation number of DM in a galaxy  $> 0(1)$ .

$m_{DM} > 10^{19}$  GeV: DM cannot be a particle,  $\lambda_c = 2\pi/m > r_s = 2m/m_{pl}^2$ .

DM can be a particle if its mass is  $10^{-1}$  eV  $< m_{DM} < 10^{19}$  GeV.

○ DM targets @ MeV- $\gamma$  region:



Observing MeV gamma will be developed shortly @ COSI, etc.  
 ✓ ALP (Wave-like cand.)

→ Fabrizio's talk tomorrow

✓ pBH (Macroscopic ".)

[T. Siebert, et al, MNRAS 511, (2002)]

✓ Light DM (Particle ".)

We will focus on the particle dark matter candidates in this talk.

# Dark matter processes producing MeV- $\gamma$

○ *Dark matter decay:*

≡ *Many well-motivated candidates producing MeV gamma-rays.*

✓ *Sterile neutrino DM ( $N$ ):*

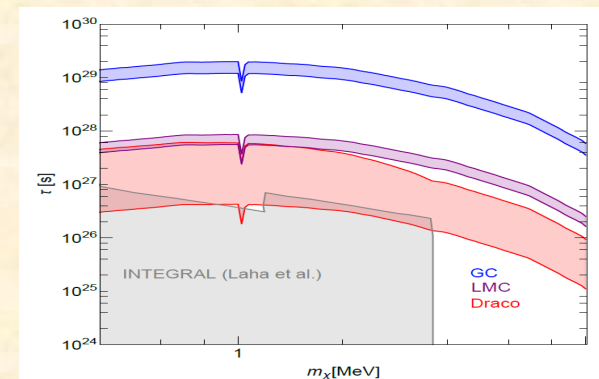
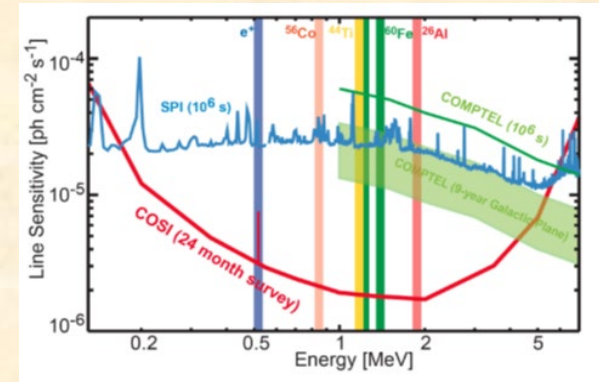
*Motivated by the generation of the neutrino masses/mixings based on  $B-L$  gauge sym. (Split seesaw, etc.)*

✓ *Gravitino DM ( $\psi_\mu$ ):*

*Motivated by the supersymmetric SM, including a (leptonic)  $R$ -parity violation due to some motivations.*

✓ *Heavy ALP DM ( $a$ ):*

*Motivated by the cosmology predicted by the string theory. (Axiverse).*



[A. Caputo, M. Negro, M. Regis, M. Taoso, JCAP02, 006, 2023.]

*Their masses are predicted to be around the MeV scale and are not absolutely stable, giving a monochromatic gamma-ray signal.*

# Dark matter processes producing MeV- $\gamma$

## Dark matter annihilation:

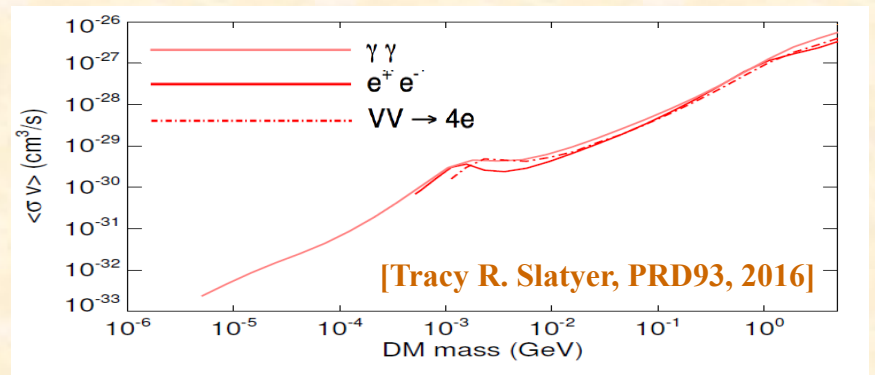
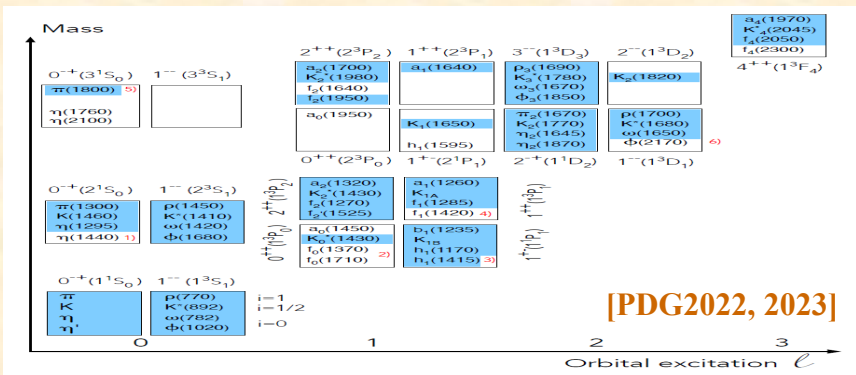
Well-motivated candidate = **Thermal DM**, a natural extension of the WIMP, i.e., the one that was in thermal equilibrium with SMs.

- Free from the initial condition problem of DM abundance.
- Detectable based on interactions maintaining the equilibrium.
- Its mass is predicted to be between  $\sim 1\text{MeV}$  and  $\sim 1\text{PeV}$ .

Light thermal DM ( $\ll \text{EW scale}$ ) is similar to WIMP but different!

✓ About hadronic final states  
When the light DM couples to quarks, we must treat final states with various hadrons.

✓ About the cross-section  
Annihilation cross-section of the light DM at recombination epoch is severely constrained.



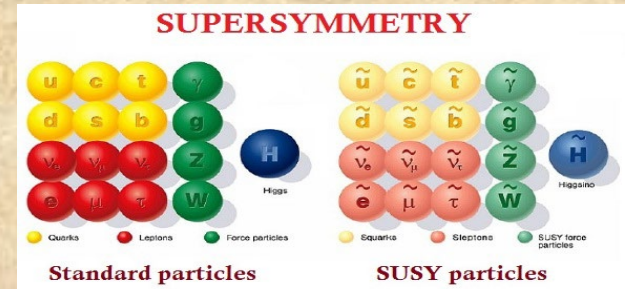
# Models for MeV- $\gamma$ from DM annihilation?

*Which strategy is being adopted for the case of WIMP (GeV- $\gamma$ )?*

**O Based on concrete models:**

*WIMP (GeV- $g$ ) signals are computed in concrete UV models (SUSY, etc.)*

*→ No vanilla models for the light DM.*



**O Based on annihilation processes:**

*The signal is computed at each process:  $e^-e^+$ ,  $qq$ ,  $\gamma\gamma$ ,  $W^-W^+$ , ...*

*→ Not good for the light DM, as stated in the previous slide. (It may work for the DM whose mass is below  $O(10)$  MeV.)*

**O Based on simplified (minimal renormalizable) models:**

*This is the strategy between the above two.*

*✓ Simplified models can be embedded in UV models.*

*✓ The models describe some parameter regions of UV models.*

*→ This strategy works for light DMs: Hadronic decays are well-defined, and various observables are well-computed.*

# Simplified models for the light DM

- ✓ **Renormalizability:** The spin of the DM should be 0,  $\frac{1}{2}$ , or 1.
- ✓ **Quantum numbers:** The DM should be singlet under SM sym.
- ✓ **Stability:** We assume that the DM is stable due to a  $Z_2$  sym.

## • Spin = 0 case:

The simplest (SM+DM) one has been ruled out for the light DM.



The next one is (SM+DM+MED), with MED being a bosonic med.

## • Spin = $\frac{1}{2}$ case:

No renormalizable int. in the (SM+DM) model.



The simplest model is the SM+DM+MED one, with MED being again a bosonic mediator.

## • Spin = 1 case:

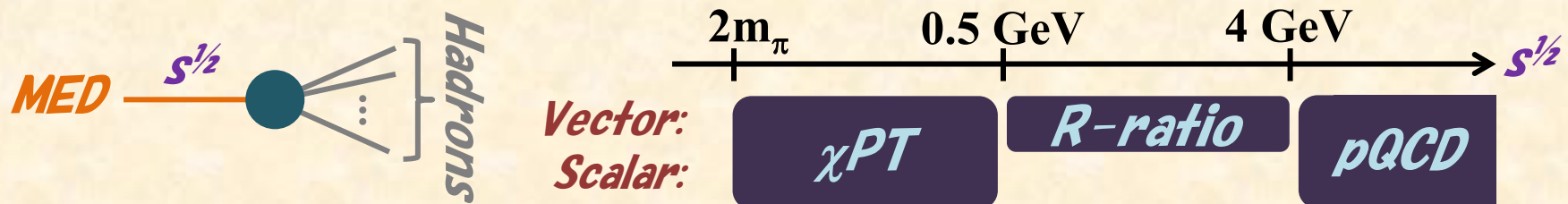
The minimal model is the so-called U(1) extension of the SM.



An Abelian Higgs field which breaks the U(1) plays the role of med.

The light DM couples to SM via a bosonic (scalar, vector) med!

○ Hadronic final states in light DM annihilation:



# Simplified models for the light DM

## ○ Avoiding CMB constraint

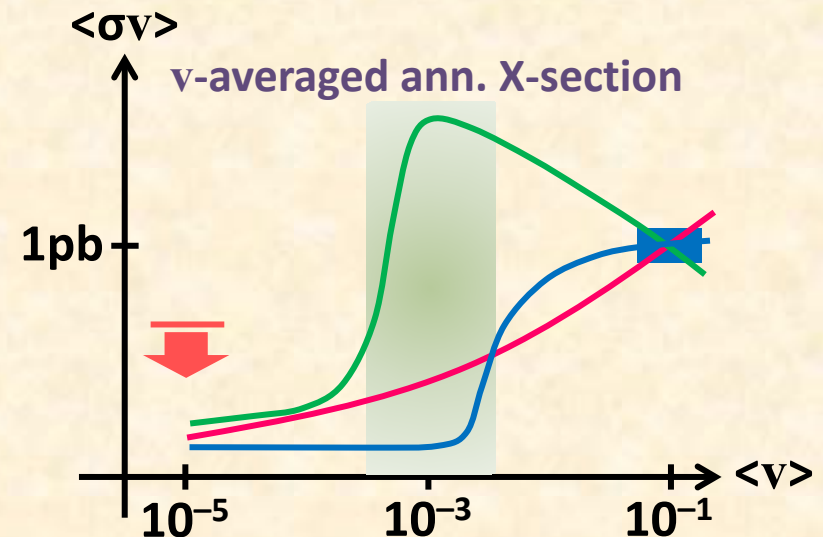
**Light DM with a velocity-dependent annihilation!**

With the averaged DM velocity  $\langle v \rangle$ ,

$\langle v \rangle \ll 10^{-3}$ : Avoiding CMB constraint.

$\langle v \rangle \sim 10^{-3}$ : MeV- $\gamma$  observation window.

$\langle v \rangle \sim 10^{-1}$ : For relic abundance cond.



✓  **$p$ -wave annihilation:**  $\sigma v \propto v^2$  @  $v \rightarrow 0$ .

$s$ -wave one is suppressed by angular and CP conservations.

✓ **Resonant annihilation:**  $\sigma v \propto [(v^2 - v_R^2)^2 + 16 \Gamma_{Res}^2 / M_{Res}^2]^{-1}$ .

Annihilation is dominated by an  $s$ -channel resonant diagram.

✓ **Forbidden annihilation:**  $\sigma v \propto (v^2 - v_{Cut}^2)^{l+1/2} \theta(v^2 - v_{Cut}^2)$ .

DM annihilation is slightly forbidden kinematically at  $v = 0$ .

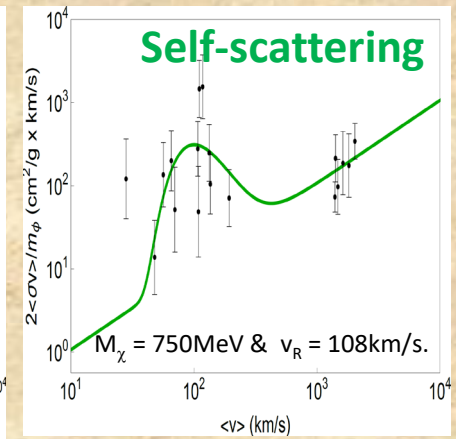
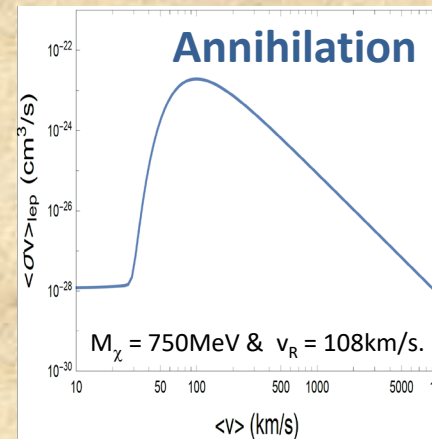
**Light DM having a  $v$ -dependent annihilation often becomes SIDM, giving a solution to the small-scale (core-cusp) crisis of our U!**



# Example: Scalar DM + Scalar MED

[T. Binder, S. Chakraborti, S. M., Y. Watanabe, JHEP01, 2022.]

At  $m_{DM} \sim m_{MED}/2$  with appropriate choosing parameters, annihilation and self-scattering processes are



[Data: M. Kaplinghat, S. Tulin, H. Yu, PRL116, 2016] ↗

## Global analysis

### Conditions/Constraints

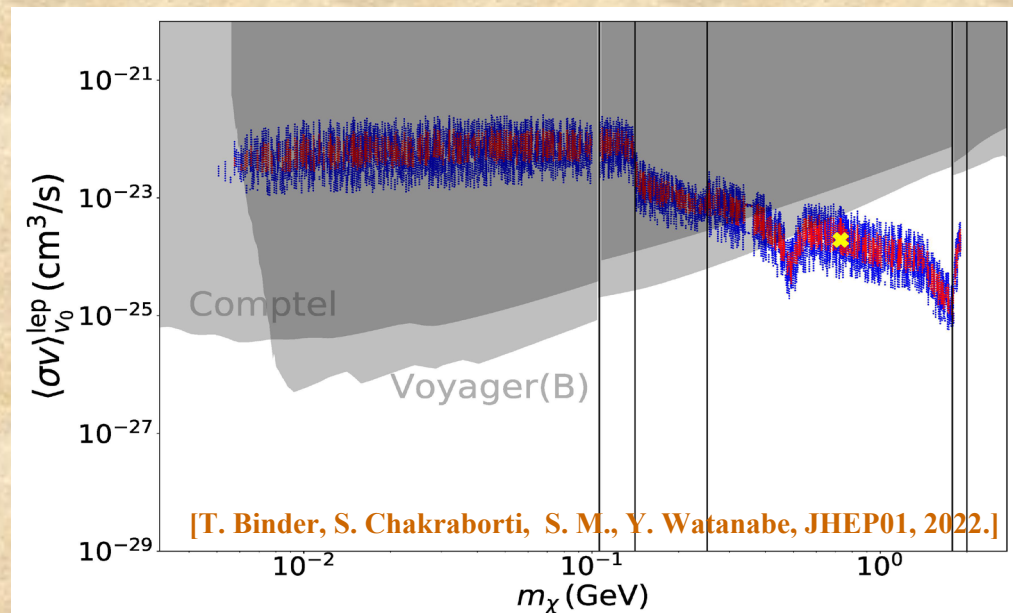
- ✓ Relic abundance condition,
- ✓ CMB constraints on  $\langle\sigma v\rangle$ ,
- ✓ Self-scattering condition.

### Indirect DM detection?

- ✓ COMPTEL ( $\gamma$ ) & Voyager ( $e$ ).

### Future MeV- $\gamma$ observation?

- ✓ COSI, AMEGO, GECCO, etc.



[T. Binder, S. Chakraborti, S. M., Y. Watanabe, JHEP01, 2022.]

# Example: Majorana DM + Scalar MED

[S.M., Y. Tsai, P. Tseng, JHEP07, 2019.]

When CP is conserved by interactions that DM participates in, the DM annihilation proceeds in the  $p$ -wave, i.e., in the  $s$ -wave is suppressed.

Attractive regions predicting a large self-scattering are as follows:



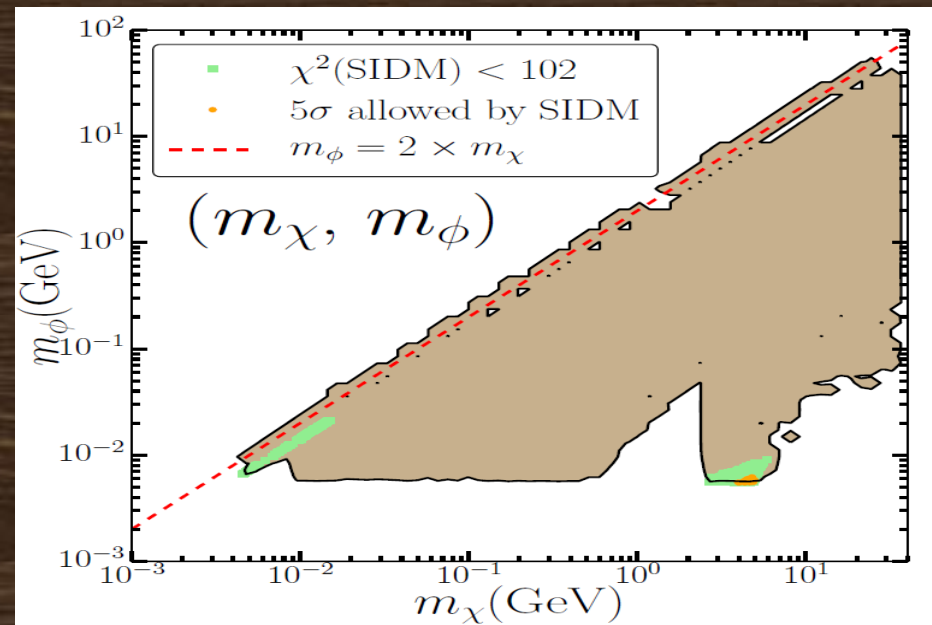
## Global analysis

### Conditions/Constraints

- ✓ Relic abundance condition,
- ✓ CMB constraints on  $\langle \sigma v \rangle$ ,
- ✓ Direct detection constraint,
- ✓ Various collider constraints,
- + Self-scattering condition.

Light MED ...  $m_{DM} = \mathcal{O}(1)$  GeV.

Forbidden ...  $m_{DM} = \mathcal{O}(10)$  MeV.



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Light MED region ( $m_{MED} \ll m_{DM}$ )



Forbidden region ( $m_{MED} \ll m_{DM}$ )

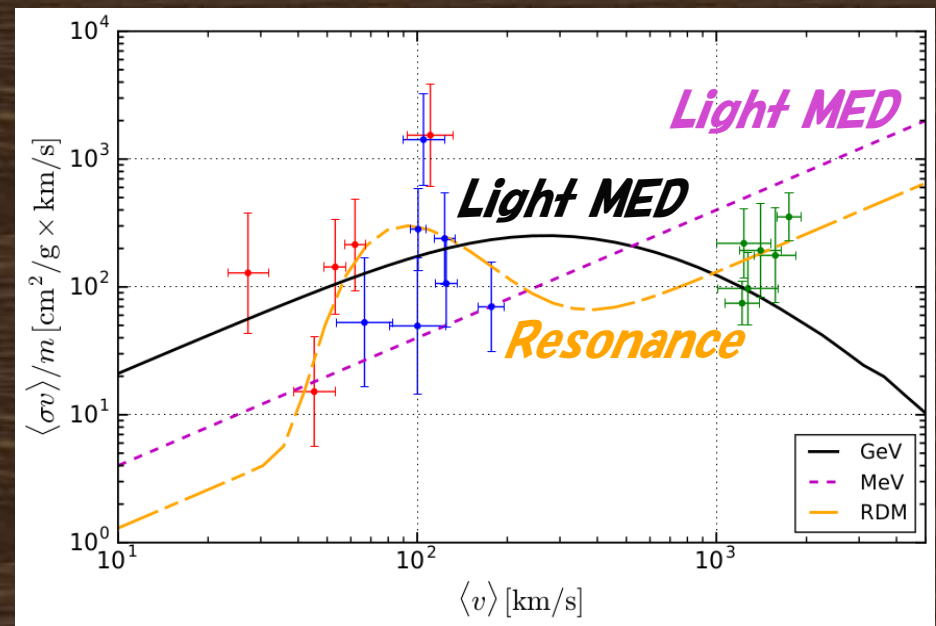
## Global analysis

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## Summary

*We discussed a framework of the light thermal dark matter aimed at quantitatively discussing the capability of future MeV- $\gamma$  observations.*

- ✓ It seems best to use the so-called simplified models (as in the case of LHC experiments for WIMP) involving the light dark matter and the bosonic mediator. Both are singlet under SM gauge interactions.*
- ✓ Annihilation into hadrons is systematically described (using chiral perturbation theory,  $R$ -ratio observation, and perturbative QCD), except for a scalar mediator case with its mass,  $m_{\text{MED}} = 0.5\text{--}4 \text{ MeV}$ . Candidates with a velocity-dependent (e.g.,  $p$ -wave, resonant, and forbidden) annihilation cross-section at the NR region should be considered to avoid the CMB constraint. Such candidates often predict a large self-scattering cross-section and become SIDM.*
- ✓ Because the energy of photons probed by COSI is 100keV to 5MeV and line gamma-rays are searched for with excellent sensitivity, we should focus on the light DM, whose mass is  $O(\text{MeV})$ , annihilating into two photons (via 1-loop, etc.) and/or (low-energy)  $e^-e^+$  ( $\gamma$ ).*