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Field-theory aspects of condensed matter physics

Examples IV

To be discussed Tuesday 17th December in the examples class

Berry curvature in quantum field theory

Based on <https://arxiv.org/pdf/1701.05587.pdf> by Baggio, Niarchos and Papadodimas

For the Schrödinger equation

$$H(\lambda)|n(\lambda)\rangle = E_n(\lambda)|n(\lambda)\rangle, \quad (1)$$

the Berry connection is given by

$$\mathbf{A}_i^n = \langle n(\lambda) | \frac{\partial}{\partial \lambda^i} | n(\lambda) \rangle, \quad (2)$$

and the Berry curvature by

$$\mathbf{F}_{ij}^{(n)} = \partial_i \mathbf{A}_j^n - \partial_j \mathbf{A}_i^n. \quad (3)$$

a) Show that using an orthonormal basis $|m\rangle$, the Berry curvature may be written as

$$\mathbf{F}_{ij}^{(n)} = \sum_{n \neq m} \frac{\langle n | \partial_i H | m \rangle \langle m | \partial_j H | n \rangle - (i \leftrightarrow j)}{(E_m - E_n)^2}. \quad (4)$$

b) Consider the Lagrangian of vacuum electrodynamics,

$$\mathcal{L} = -\frac{1}{4e^2} F_{\mu\nu} F^{\mu\nu} + \frac{\theta}{32\pi^2} F_{\mu\nu} \tilde{F}^{\mu\nu}, \quad \tilde{F}_{\mu\nu} = \frac{1}{2} \varepsilon_{\mu\nu\sigma\rho} F^{\sigma\rho}. \quad (5)$$

Show that with $F_{\mu\nu}^\pm = F_{\mu\nu} \pm i\tilde{F}_{\mu\nu}$, (5) may be rewritten as

$$\mathcal{L} = \frac{i}{64\pi} \tau (F^+)^2 - \frac{i}{64\pi} \bar{\tau} (F^-)^2, \quad \tau = \frac{\theta}{2\pi} + i \frac{4\pi}{e^2}. \quad (6)$$

By considering photon states $|\vec{p}, \epsilon\rangle$ with momentum \vec{p} and helicity $\epsilon = \pm$, use (4) to show that the Berry curvature is

$$\mathbf{F}_{\tau\bar{\tau}}^{(\vec{p}, \vec{p}', \epsilon, \epsilon')} = \frac{\epsilon}{8} \frac{1}{(\text{Im}\tau)^2} \delta_{\vec{p}, \vec{p}'} \delta_{\epsilon, \epsilon'}. \quad (7)$$

Using this result, deduce that for a multi-photon state containing n_+ photons of positive and n_- photons of negative helicity, the Berry curvature is

$$\mathbf{F}_{\tau\bar{\tau}}^{(n_+, n_-)} = \frac{1}{8} (n_+ - n_-) \frac{1}{(\text{Im}\tau)^2}. \quad (8)$$