

String Theory II - Exercise Sheet 01

Due to Thu Dec 21 2017 10-12 Uhr M

Problem 1.1: Supersymmetry and Stability

The $\mathcal{N} = 1$ supersymmetry algebra in 3+1 dimensions is generated by a single Weyl spinor $Q_\alpha = (Q_1, Q_2)^T$ with complex conjugate $\bar{Q}_{\dot{\alpha}} = Q_\alpha^\dagger$, translation generator P^μ and Lorentz generator $M^{\mu\nu}$. The algebra reads (nonvanishing conditions only)

$$\{Q_\alpha, \bar{Q}_{\dot{\beta}}\} = 2\sigma_{\alpha\dot{\beta}}^\mu P_\mu \quad (1)$$

$$[M^{\mu\nu}, Q_\alpha] = \frac{1}{2}(\sigma^{\mu\nu})_\alpha{}^\beta Q_\beta \quad (2)$$

$$[M^{\mu\nu}, M^{\rho\sigma}] = \eta^{\nu\rho} M^{\mu\sigma} - \eta^{\mu\rho} M^{\nu\sigma} + \eta^{\mu\sigma} M^{\nu\rho} - \eta^{\nu\sigma} M^{\mu\rho} \quad (3)$$

$$[M^{\mu\nu}, P^\rho] = \eta^{\mu\rho} P^\nu - \eta^{\nu\rho} P^\mu. \quad (4)$$

Here

$$\sigma_{\alpha\dot{\beta}}^\mu = (-\mathbb{1}, \sigma^i)_{\alpha\dot{\beta}} \quad (5)$$

$$\bar{\sigma}^{\mu\dot{\alpha}\beta} = \epsilon^{\dot{\alpha}\gamma} \epsilon^{\beta\delta} \sigma_{\delta\dot{\gamma}}^\mu \quad (6)$$

$$\epsilon_{21} = \epsilon^{12} = 1, \quad \epsilon_{12} = \epsilon^{21} = -1, \quad \epsilon_{11} = \epsilon_{22} = \epsilon^{11} = \epsilon^{22} = 0, \quad (7)$$

$$\sigma^{\mu\nu} = \frac{1}{4}(\sigma^\mu \bar{\sigma}^\nu - \sigma^\nu \bar{\sigma}^\mu). \quad (8)$$

$\sigma^{\mu\nu}$ is the generator of the Lorentz algebra (3) in the Weyl spinor representation. Show that

1. $\bar{\sigma}^\mu = (\sigma^0, -\sigma^i)$
2. $\text{Tr}(\sigma^\mu \bar{\sigma}^\nu) = -2\eta^{\mu\nu}$
3. that $P^\mu = -\frac{1}{4}\{Q_\alpha, \bar{Q}_{\dot{\beta}}\} \bar{\sigma}^{\mu\dot{\beta}\alpha}$ and from here that P^0 is a positive operator, $\langle \Psi | P^0 | \Psi \rangle \geq 0 \quad \forall \quad | \Psi \rangle$.

The last relation implies that the supersymmetric system has a stable ground state.

Problem 1.2: Space-Time Symmetries of the Polyakov action

Use the trick discussed in the exercise session on Dec. 15 to derive the conserved currents for Poincaré and Lorentz transformations

$$P_\mu^a = -T\eta^{ab}\partial_b X_\mu \quad (9)$$

$$M_{\mu\nu}^a = X_\mu P_\nu^a - X_\nu P_\mu^a \quad (10)$$

as well as the corresponding conserved charges for a string of length ℓ described by the Polyakov action in conformal gauge

$$S_{Pol} = -\frac{T}{2} \int d^2\sigma (\partial_a X^\mu)(\partial^a X_\mu). \quad (11)$$

Use the Poisson bracket structure between canonical coordinates X^μ and momenta $\Pi_\mu = P_\mu^\tau$

$$\{A[X, \Pi], B[X, \Pi]\}_{P.B.} = \int d\sigma'' \left(\frac{\delta A}{\delta X^\mu(\sigma'')} \frac{\delta B}{\delta \Pi_\mu(\sigma'')} - \frac{\delta A}{\delta \Pi^\mu(\sigma'')} \frac{\delta B}{\delta X_\mu(\sigma'')} \right) \quad (12)$$

to show that the conserved charges of Poincaré and Lorentz symmetry fulfill the classical version of the algebra (3) and (4),

$$\{.,.\} = -i[.,.] \quad (13)$$