

Unbounded Operators / Dynamics of States

1.1 CCRs vs. Boundedness

Julius-Maximilians-

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Consider two bounded operators A and B on a Hilbert space \mathcal{H} , i.e.

$$\exists c_A \in \mathbf{R} : \forall \psi \in \mathcal{H} : \|A\psi\| \le c_A \|\psi\|$$
(1a)

$$\exists c_B \in \mathbf{R} : \forall \psi \in \mathcal{H} : \|B\psi\| \le c_B \|\psi\|.$$
(1b)

Show that the canonical commutation relations

$$[A, B] = AB - BA = \mathbf{i} \tag{2}$$

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are inconsistent with the assumption of boundedness for the operators A and B.

NB: it is *not* necessary to find an original proof. It suffices to find, understand and present a proof from the literature.

1.2 Classical Dynamics on the 2-Torus

Consider a classical dynamical system with the 2-Torus $T^2 = S^1 \times S^1$ as phase space Γ (this is not a cotangent bundle, but it has the technical advantage of being compact).

Using standard coordinates $(\theta_1, \theta_2) \in [0, 2\pi)^2$, a consistent Poisson bracket is given by

$$\{f,g\} = \left(\frac{\partial f}{\partial \theta_1}\frac{\partial g}{\partial \theta_2} - \frac{\partial f}{\partial \theta_2}\frac{\partial g}{\partial \theta_1}\right).$$
(3)

Assume that the Hamiltonian is

$$H: \Gamma \to \mathbf{R} (\theta_1, \theta_2) \mapsto H(\theta_1, \theta_2) = c \cos \theta_1$$
(4)

In order to be well defined globally, the Hamiltonian must be periodic in θ_1 and θ_2 . This is the simplest choice.

1. Derive the equations of motion

- 2. Determine the flow Φ of a phase space point $(\theta_1, \theta_2) \in \Gamma$
- 3. Determine the time evolution of the state ω , where

$$\omega(f) = \int_{\Gamma} d^2\theta \,\omega(\theta) f(\theta) \tag{5}$$

with

$$\omega: \Gamma \to \mathbf{R}$$

$$(\theta_1, \theta_2) \mapsto \omega(\theta_1, \theta_2) = \frac{1}{\pi^2} \sin^2 \theta_1 \sin^2 \theta_2.$$
(6)

1.3 Classical Dynamics on the 2-Sphere

Consider a classical dynamical system with the 2-Sphere S^2 as phase space Γ (this is again not a cotangent bundle, but it has the technical advantage of being compact and is highly symmetric).

Using standard spherical coordinates $(\theta, \phi) \in [0, \pi) \times [0, 2\pi]$, a consistent Poisson bracket is given by

$$\{f,g\} = \frac{1}{\sin\theta} \left(\frac{\partial f}{\partial\theta} \frac{\partial g}{\partial\phi} - \frac{\partial f}{\partial\phi} \frac{\partial g}{\partial\theta} \right)$$
(7)

Assume that the Hamiltonian is

$$H: \Gamma \to \mathbf{R} (\theta, \phi) \mapsto H(\theta, \phi) = c \cos \theta$$
 (8)

In order to be well defined globally, the Hamiltonian must be periodic in θ and ϕ . This one of the simplest choices.

- 1. Show that the Poisson bracket satisfies all requirements.
- 2. Determine the flow Φ of a phase space point $(\theta, \phi) \in \Gamma$
- 3. Determine the time evolution of the state ω , where

$$\omega(f) = \int_{\Gamma} \sin\theta \,\mathrm{d}\theta \mathrm{d}\phi \,\omega(\theta,\phi) f(\theta,\phi) \tag{9}$$

with

$$\omega: \Gamma \to \mathbf{R}$$

$$(\theta, \phi) \mapsto \omega(\theta, \phi) = \frac{2}{\pi^2} \sin \theta \, \cos^2 \phi \,. \tag{10}$$