Karlsruher Institut für Technologie

Particle Physics Phenomenology
after the Higgs Discovery

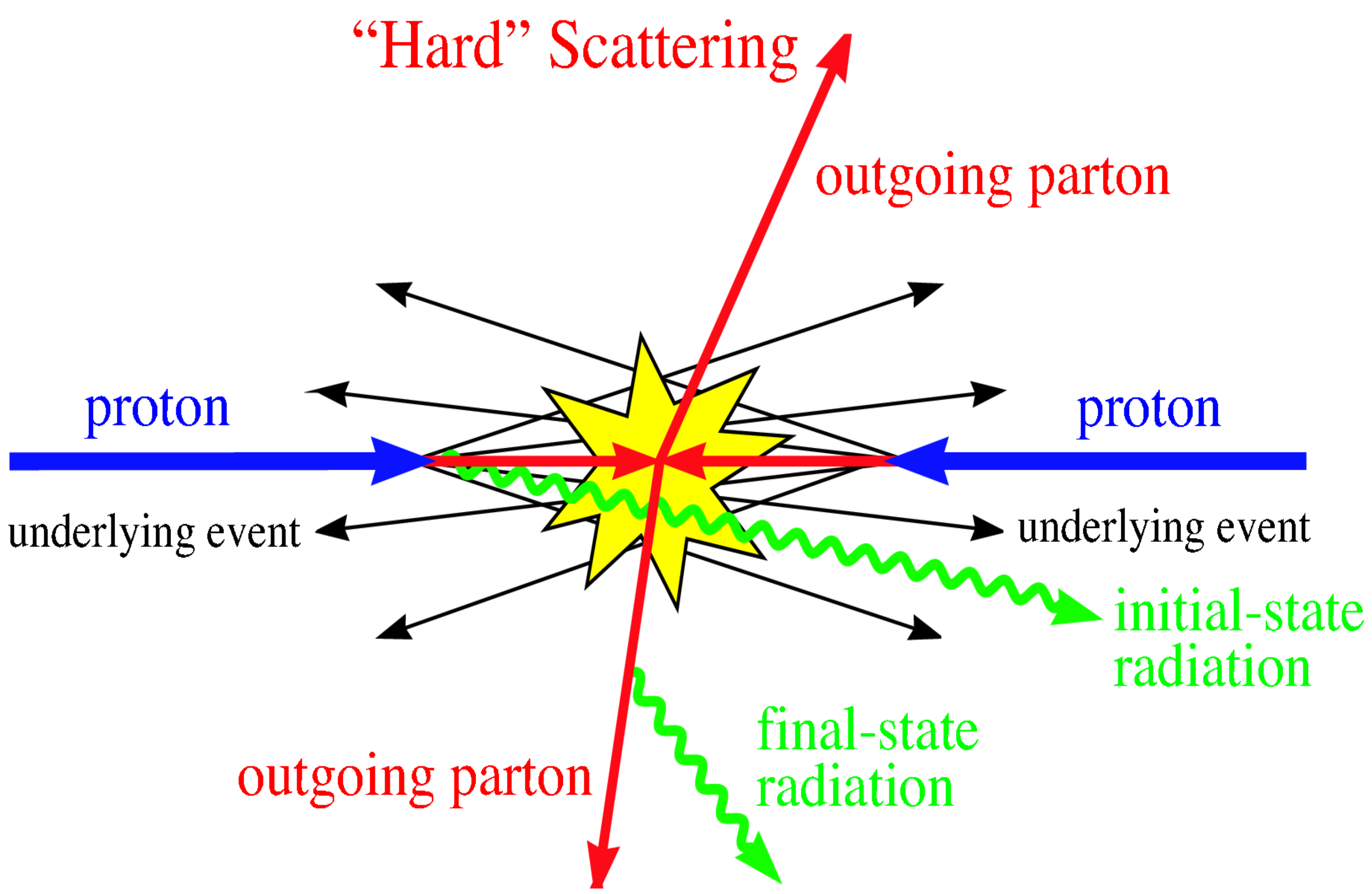
KIRILL MELNIKOV

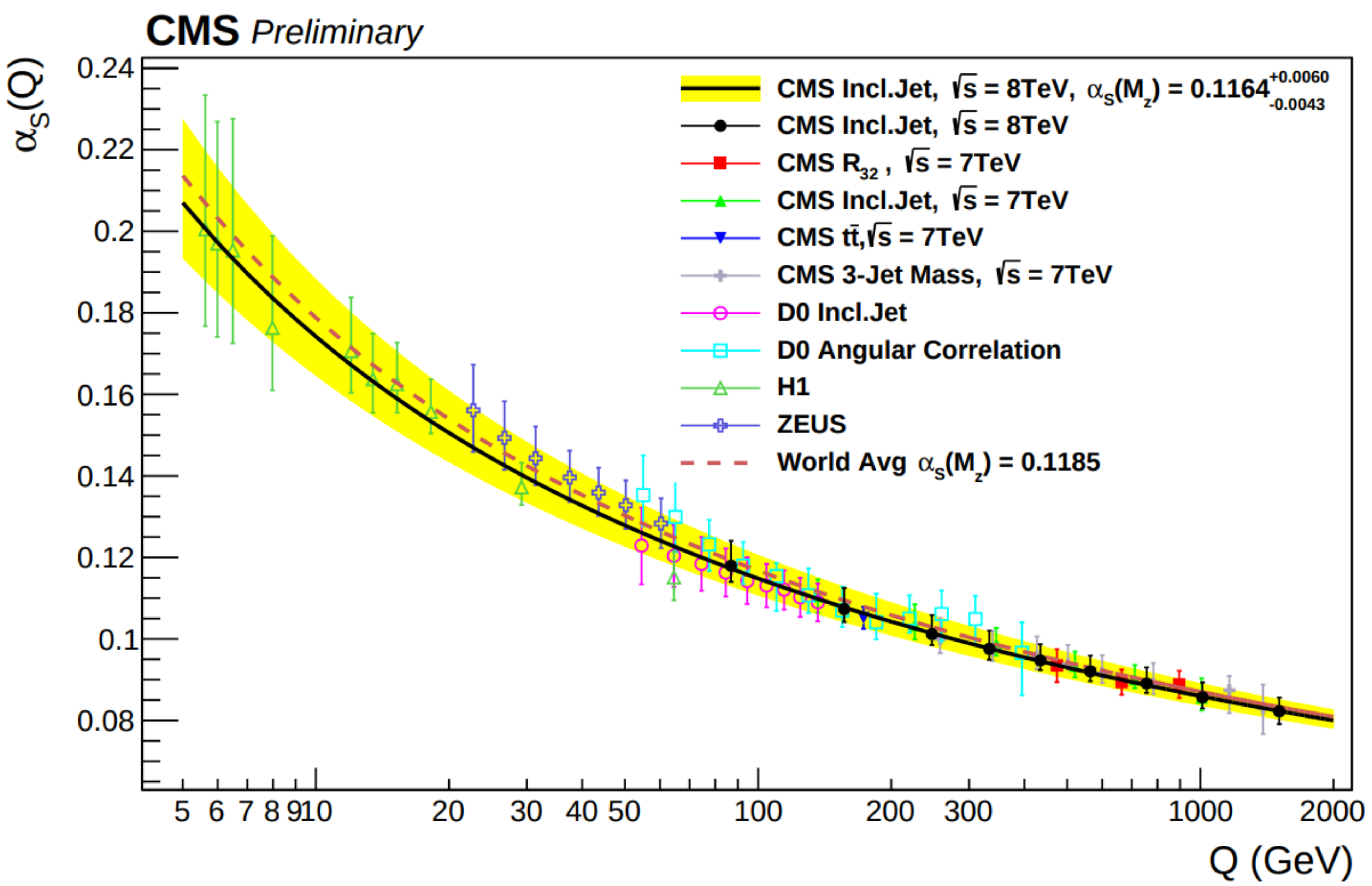
PRECISION QCD AT THE LHC: CHALLENGES AND OPPORTUNITIES

Inaugural Workshop of the RTG "Particle Physics at colliders at the LHC precision era", March 17-18 2025, Würzburg

HARD COLLISIONS CAN BE DESCRIBED FROM FIRST PRINCIPLES

Studies of hadron collisions with large momentum transfer allow us to explore heaviest particles in the Standard Model and search for new particle and interactions. Such collisions are amenable to a rigorous theoretical descriptions based on first principles.





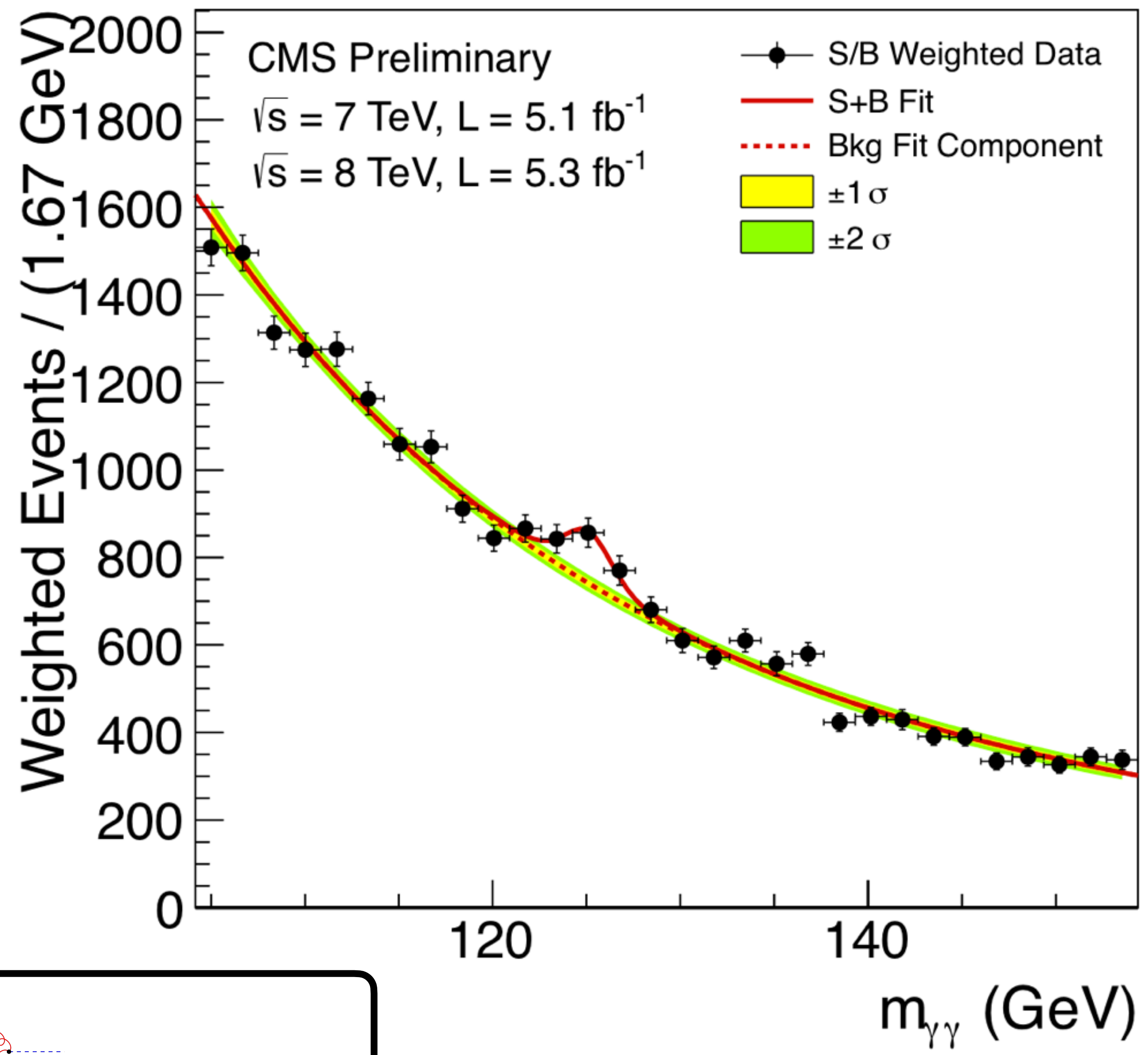
$$\mathcal{L}_{\rm QCD} = \sum \bar{q}_j \left(i\hat{D} - m_j \right) q_j - \frac{1}{4} G^a_{\mu\nu} G^{a,\mu\nu}$$

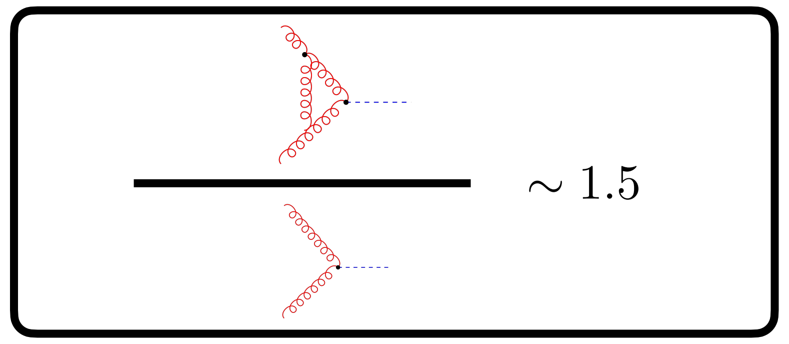
$$\mathrm{d}\sigma_{\rm hard} = \sum_{ij\in\{q,g\}} \int \mathrm{d}x_1\, \mathrm{d}x_2 f_i(x_1) f_j(x_2)\, \mathrm{d}\sigma_{ij}(x_1, x_2), \{p_{\rm fin}\})\, O_J(\{p_{\rm fin}\}).$$

$$\mathrm{d}\sigma_{ij} = \mathrm{d}\sigma_{ij,\rm LO} \left(1 + \alpha_s\, \Delta_{ij,\rm NLO} + \alpha_s^2\, \Delta_{ij,\rm NNLO} + \ldots\right)$$

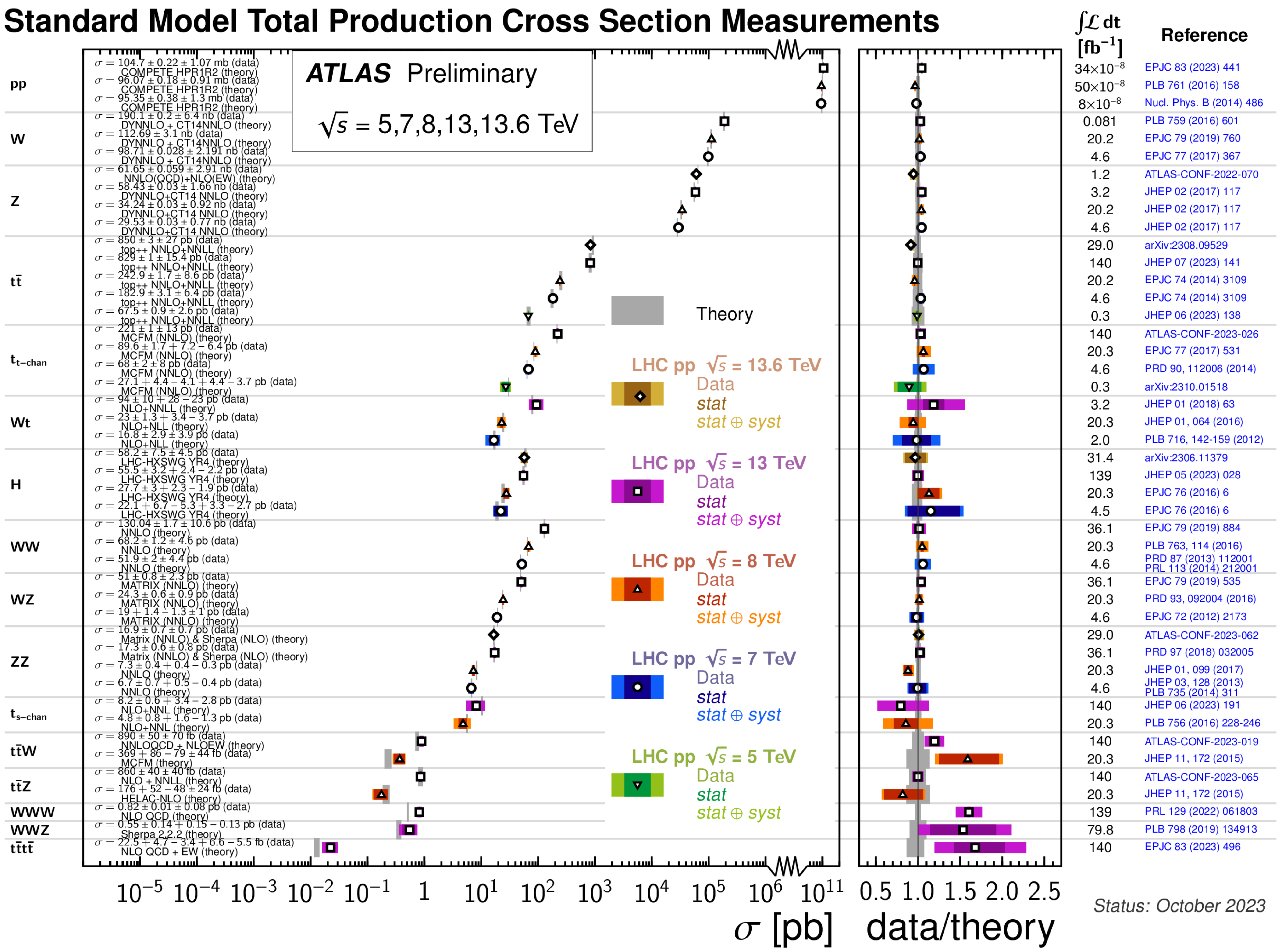
PERTURBATIVE QCD FACILITATES INTERPRETATION OF LHC MEASUREMENTS

Perturbative QCD is very well understood by now. We use it for an unambiguous interpretation of all LHC measurements.





$$
\begin{aligned}
48.58\,\text{pb} = \; & 16.00\,\text{pb} & (+32.9\%) & \quad (\text{LO, rEFT}) \\
& + 20.84\,\text{pb} & (+42.9\%) & \quad (\text{NLO, rEFT}) \\
& - \;\; 2.05\,\text{pb} & (-4.2\%) & \quad ((t,b,c),\ \text{exact NLO}) \\
& + \;\; 9.56\,\text{pb} & (+19.7\%) & \quad (\text{NNLO, rEFT}) \\
& + \;\; 0.34\,\text{pb} & (+0.7\%) & \quad (\text{NNLO}, 1/m_t) \\
& + \;\; 2.40\,\text{pb} & (+4.9\%) & \quad (\text{EW, QCD-EW}) \\
& + \;\; 1.49\,\text{pb} & (+3.1\%) & \quad (\text{N}^3\text{LO, rEFT})
\end{aligned}
$$



Anastasiou, Duhr, Dulat, Furlan, Gehrmann, Herzog, Lazopoulos, Mistlberger

PERTURBATIVE QCD AS QUANTUM FIELD THEORY

Demands of the LHC physics program have shaped the development of QCD as a perturbative Quantum Field Theory and kept QCD practitioners on their toes.

Higher order cross sections:The first wishlist

An experimenter's wishlist

Run II Monte Carlo Workshop

Single Boson	Diboson	Triboson	Heavy Flavour
$W+\leq 5j$	$WW+\leq 5j$	$WWW+\leq 3j$	$t\bar{t}+\leq 3j$
$W+b\bar{b}\leq 3j$	$W+b\bar{b}+\leq 3j$	$WWW+b\bar{b}+\leq 3j$	$t\bar{t}+\gamma+\leq 2j$
$W+c\bar{c}\leq 3j$	$W+c\bar{c}+\leq 3j$	$WWW+\gamma\gamma+\leq 3j$	$t\bar{t}+W+\leq 2j$
$Z+\leq 5j$	$ZZ+\leq 5j$	$Z\gamma\gamma+\leq 3j$	$t\bar{t}+Z+\leq 2j$
$Z+b\bar{b}+\leq 3j$	$Z+b\bar{b}+\leq 3j$	$ZZZ+\leq 3j$	$t\bar{t}+H+\leq 2j$
$Z+c\bar{c}+\leq 3j$	$ZZ+c\bar{c}+\leq 3j$	$WZZ+\leq 3j$	$t\bar{b}\leq 2j$
$\gamma+\leq 5j$	$\gamma\gamma+\leq 5j$	$ZZZ+\leq 3j$	$b\bar{b}+\leq 3j$
$\gamma+b\bar{b}\leq 3j$	$\gamma\gamma+b\bar{b}\leq 3j$		single top
$\gamma+c\bar{c}\leq 3j$	$\gamma\gamma+c\bar{c}\leq 3j$		
	$WZ+\leq 5j$		
	$WZ+b\bar{b}\leq 3j$		
	$WZ+c\bar{c}\leq 3j$		
	$W\gamma+\leq 3j$		
	$Z\gamma+\leq 3j$		

circa 20 years ago, NLO calculations are requested

Wishlist, the 2022 edition

process	known	desired
$pp\to t\bar{t}$	$\mathrm{NNLO_{QCD}}+\mathrm{NLO_{EW}}$ (w/o decays) $\mathrm{NLO_{QCD}}+\mathrm{NLO_{EW}}$ (off-shell effects) $\mathrm{NNLO_{QCD}}$ (w/ decays)	$\mathrm{N^3LO_{QCD}}$
$pp\to t\bar{t}+j$	$\mathrm{NLO_{QCD}}$ (off-shell effects) $\mathrm{NLO_{EW}}$ (w/o decays)	$\mathrm{NNLO_{QCD}}+\mathrm{NLO_{E}}$… …ays)
$pp\to t\bar{t}+2j$	$\mathrm{NLO_{QCD}}$ (w/o decays)	$\mathrm{NLO_{QCD}}$ … decays)
$pp\to V$	$\mathrm{N^3LO_{QCD}}$ $\mathrm{N^{(1,1)}LO_{QCD\otimes EW}}$ $\mathrm{NLO_{EW}}$	$\mathrm{N^3LO_{QCD}}+\mathrm{N^{(1,1)}LO_{QCD\otimes EW}}$ $\mathrm{N^2LO_{EW}}$
$pp\to VV'$	NNLO… …W + … (gg channel)	$\mathrm{NLO_{QCD}}$ (gg channel, w/ massive loops) $\mathrm{N^{(1,1)}LO_{QCD\otimes EW}}$
$pp\to V+j$	…$\mathrm{NLO_{QCD}}+\mathrm{NLO_{EW}}$	hadronic decays
pp…	$\mathrm{NLO_{QCD}}+\mathrm{NLO_{EW}}$ (QCD component) $\mathrm{NLO_{QCD}}+\mathrm{NLO_{EW}}$ (EW component)	$\mathrm{NNLO_{QCD}}$
$pp\to 2\,\text{jets}$	$\mathrm{NNLO_{QCD}}$ $\mathrm{NLO_{QCD}}+\mathrm{NLO_{EW}}$	$\mathrm{N^3LO_{QCD}}+\mathrm{NLO_{EW}}$
$pp\to 3\,\text{jets}$	$\mathrm{NNLO_{QCD}}+\mathrm{NLO_{EW}}$	

A need for N3LO and even N4LO results

and many many more…

Huss, Huston, Jones, Pellen

CHALLENGES IN DESCRIBING HARD COLLISIONS AT THE LHC

Although the evolution of perturbative QCD in the past decade was remarkable, we should not forget that our goal is to describe LHC events which start and end with hadrons. For this reason, two very different challenges need to be addressed for improving theoretical framework that we use to describe hard hadron collisions:

1) technical problems: develop efficient methods to describe quark and gluon collisions to higher and higher orders in QCD perturbation theory

$$\mathrm{d}\sigma_{ij} = \mathrm{d}\sigma_{ij,\mathrm{LO}} \left(1 + \alpha_s\, \Delta_{ij,\mathrm{NLO}} + \alpha_s^2\, \Delta_{ij,\mathrm{NNLO}} + ?.\right)$$

2) conceptual problems: find systematically-improvable description of proton-to-partons and partons-to-hadrons transitions, which are relevant for initial and final stages of the process. This problem can only be addressed if a better understanding of non-perturbative power $\mathcal{O}(\Lambda_{\mathrm{QCD}})$ corrections in collider processes is achieved.

$$\mathrm{d}\sigma_{\mathrm{hard}} = \sum_{ij\in\{q,g\}} \int \mathrm{d}x_1\, \mathrm{d}x_2 f_i(x_1) f_j(x_2)\, \mathrm{d}\sigma_{ij}(x_1, x_2), \{p_{\mathrm{fin}}\})\; O_J(\{p_{\mathrm{fin}}\}) \left(1 + \mathcal{O}(\Lambda^n_{\mathrm{QCD}}/Q^n)\right)$$

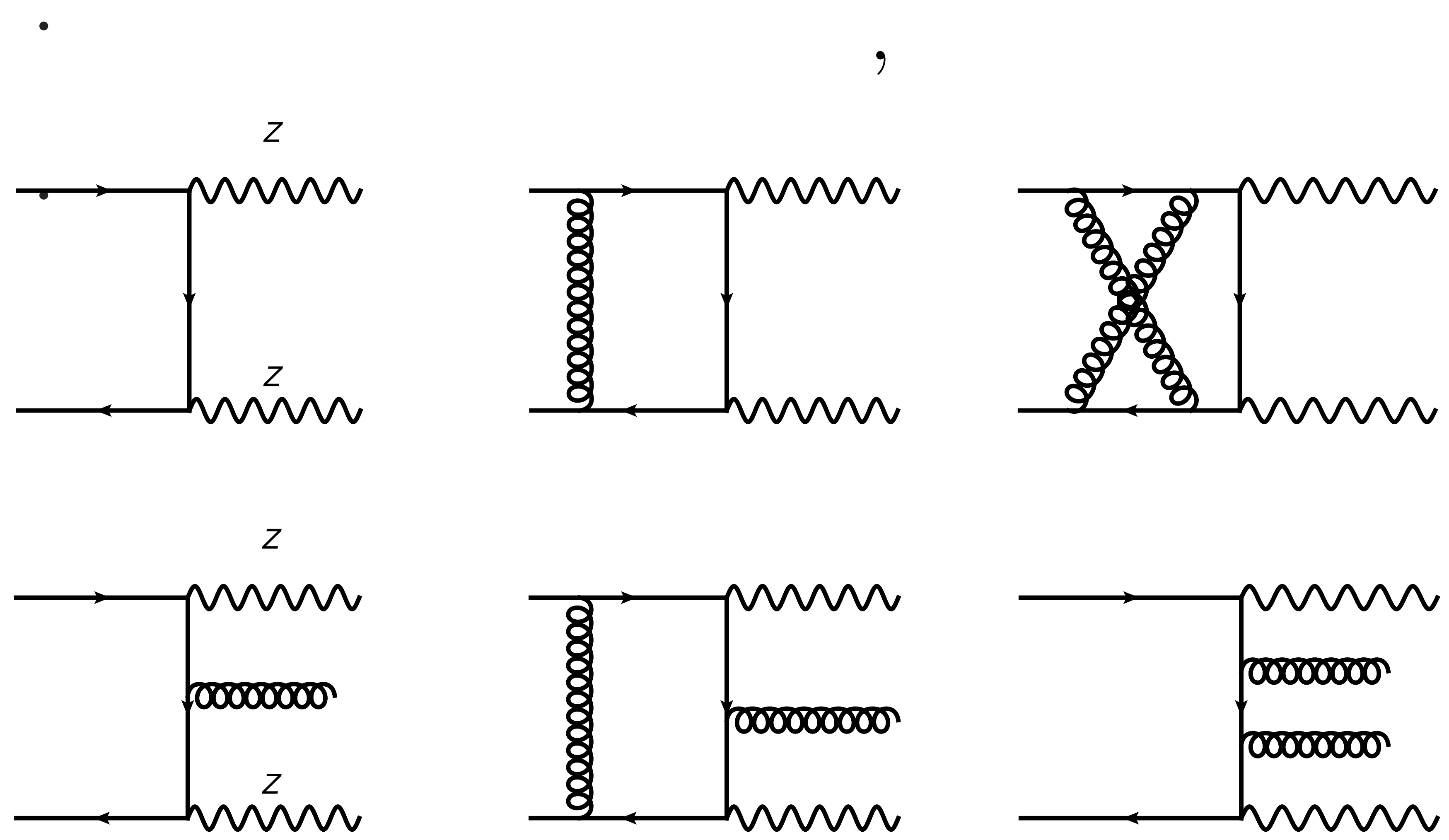
$$N_c \left(\frac{\alpha_s}{\pi}\right)^2 \sim \frac{\Lambda_{\mathrm{QCD}}}{Q}, \quad \Lambda_{\mathrm{QCD}} \sim 0.3\ \mathrm{GeV}, \quad Q \sim 30\ \mathrm{GeV}$$

PERTURBATIVE CHALLENGES

FIXED ORDER CHALLENGES

Any perturbative computation in higher orders of QCD requires calculation of loop amplitudes and real-emission contributions.Taken separately, they are divergent and only their combination is finite.

To perform phenomenologically-relevant computations, we need to:



1) figure out how to compute virtual loop amplitudes;

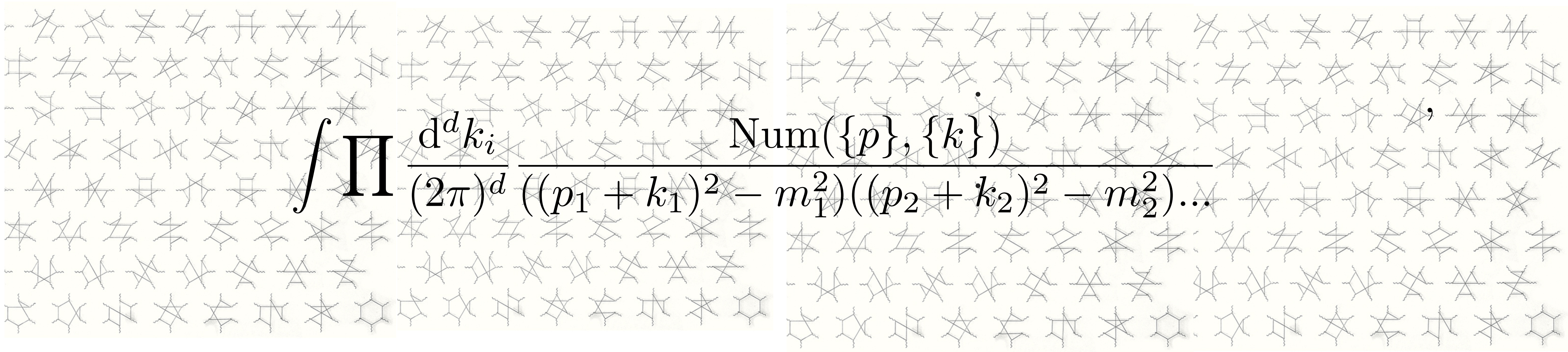
2) understand how to integrate infra-red divergent contributions over partonic phase spaces;

3) implement the emerging procedure into efficient numerical codes;

$$d\sigma_{ij} = d\sigma_{ij,\mathrm{LO}}\left(1 + \alpha_s\,\Delta_{ij,\mathrm{NLO}} + \alpha_s^2\,\Delta_{ij,\mathrm{NNLO}} + ...\right)$$

LOOP AMPLITUDES

The problem of computing loop amplitudes is the problem of calculating divergent integrals of rational functions in Minkowski space.



$$\int \prod \frac{\mathrm{d}^d k_i}{(2\pi)^d} \frac{\mathrm{Num}(\{p\},\{k\})}{((p_1+k_1)^2-m_1^2)((p_2+k_2)^2-m_2^2)\ldots},$$

www.edwardtufte.com

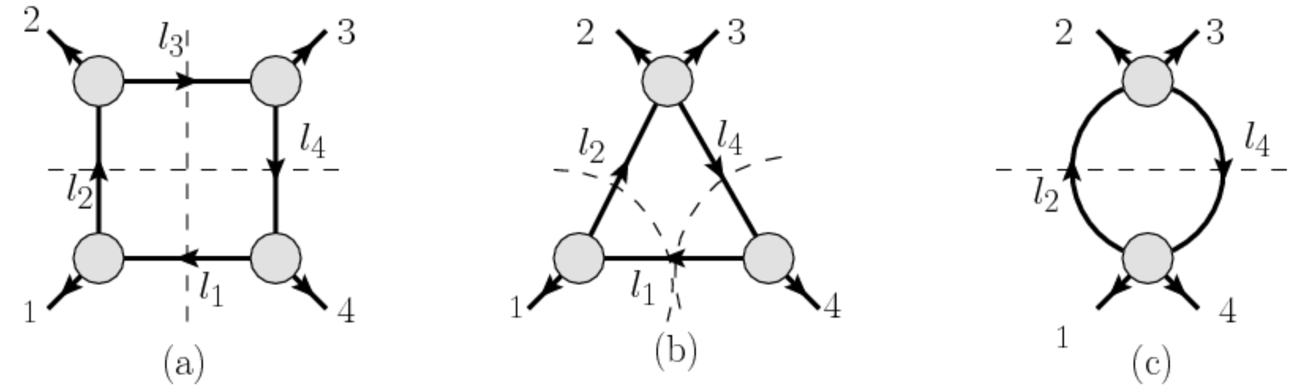
Different techniques to address this problem were developed over time, from analytic to numerical.

$$\sum c_n(s_{ij}, m_k^2) I_k = 0$$

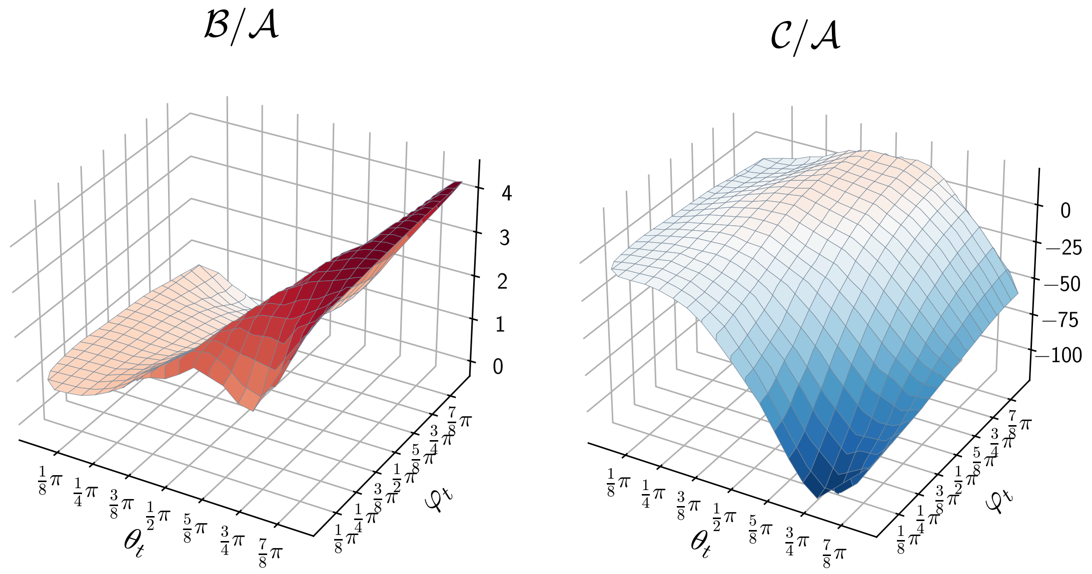
Integration-by-parts

$$s_i \frac{\partial}{\partial s_i} \vec{I} = \epsilon \hat{A}(\{s\}) \vec{I}$$

Differential equations



Generalized unitarity



$$G(\{a_n, \vec{a}_{n-1}\}, x) = \int_0^1 \frac{\mathrm{d}t}{t-a_n} G(\{\vec{a}_{n-1}, t) \qquad K(x,a) = \int_0^x \frac{\mathrm{d}t}{\sqrt{(1-t^2)(1-at^2)}}$$

Classes of functions, from Goncharov polylogarithmis, to elliptic integrals.

Numerics: integration, solution of differential equations

Chetyrkin, Tkachov, Laporta, Smirnov, von Manteufffel, Lee, Maierhoefer, Usovitsch, Uwer, Abreu, Cordero, Ita, Page, Zeng;, Badger, Hartano, Peraro, Sotnikov, Zola, Gehrman, Henn, Chicherin, Tancredi, Caola, Buncioni, Devoto, Chen, Czakon, Poncelet, Greiner, Heinrich, Kerner, Jones, Liu, Ma, C.Y.Wang, Moriello, Steinhauser, Schönwald, Anastasiou, Sterman, Hirschi

INTEGRATING REAL EMISSION CONTRIBUTIONS

Real emission contributions are integrated over partonic phase spaces with the help of subtraction and slicing methods.

$$\int |\mathcal{M}|^2 \, F_J \, \mathrm{d}\phi_d = \int \left[|\mathcal{M}|^2 \, F_J - S \right] \, \mathrm{d}\phi_4 + \int S \mathrm{d}\phi_d$$

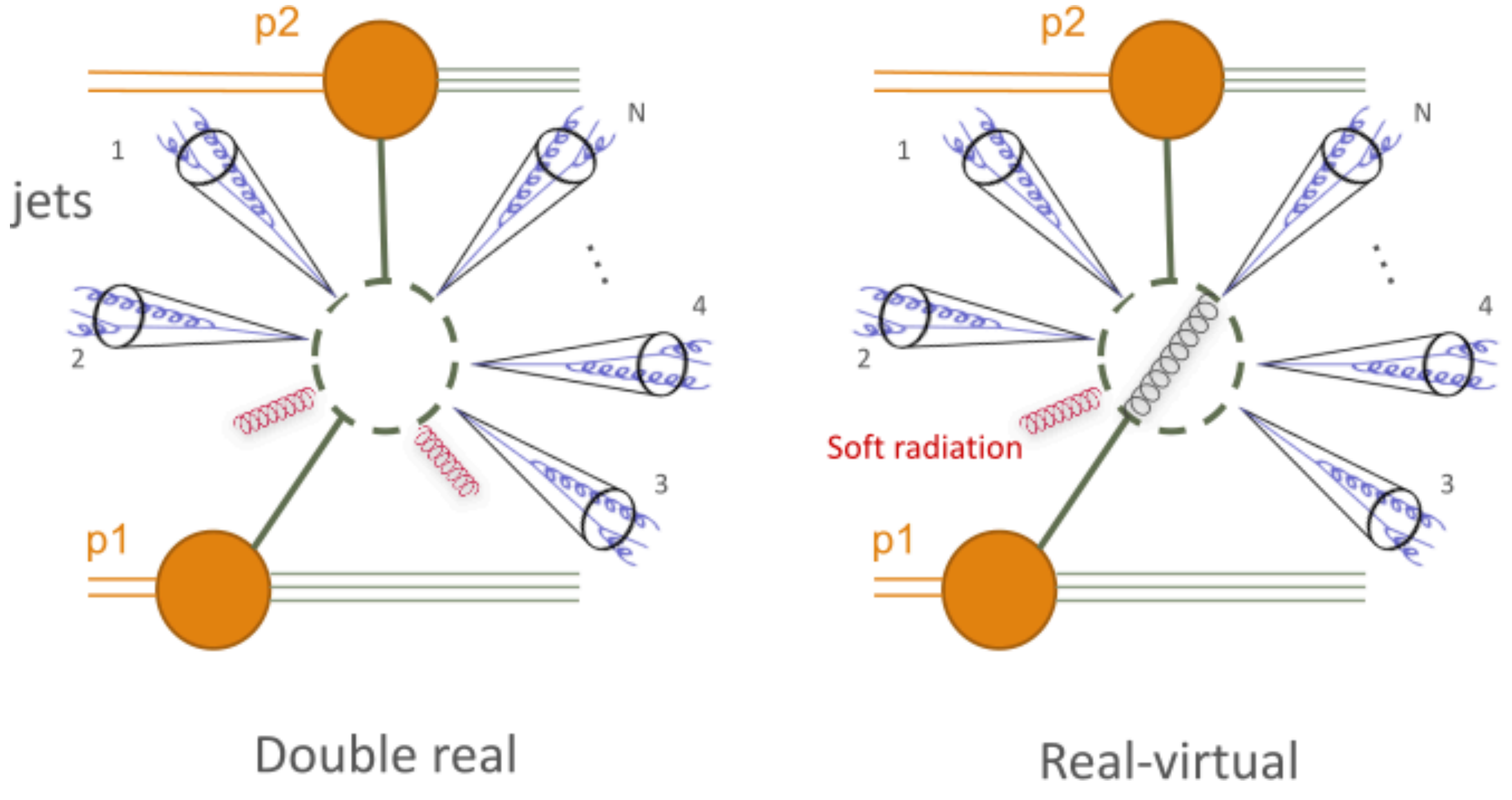
$$\int |\mathcal{M}|^2 \, F_J \, \mathrm{d}\phi_d = \int_0^\delta \left[|\mathcal{M}|^2 \, F_J \, \mathrm{d}\phi_d \right]_{\mathrm{simp}} + \int_\delta^1 |\mathcal{M}|^2 \, F_J \, \mathrm{d}\phi_4 + \mathcal{O}(\delta)$$

In both cases, one needs to know singular limits of amplitudes squared, and one should be able to integrate subtraction/slicing terms over singular parts of phase spaces. Integrals of subtraction and slicing terms contain infra-red divergencies that should cancel with similar divergencies in loop contributions.

Gehrmann, Glover, Czakon, Caola, Roentsch, K.M., Troscanyi, Somogyi, Del Duca, Duhr, Kardos, Magnea, Bertolotti, Pelliccioli, Uccirati, Torrielli, Signorile-Signorile, Catani, Grazzini, Boughezal, Petriello, Tackmann, Gaunt, Stahlhofner

In recent years, extensions of existing NNLO slicing and subtraction methods appeared, where such cancellations are demonstrated analytically for arbitrary collider processes.

Magnea, Bertolotti, Pelliccioli, Uccirati, Torrielli, Signorile-Signorile, Tagliabue, Devoto, Roetsch, Melnikov; Bell, Dehnadi, Mohrmann, Rahn, Pedron, Agarwal



These developments bring us one step closer to the formulation of an ultimate subtraction scheme at NNLO which will be amenable to a straightforward automation and will enable the construction of general-purpose codes, capable of computing real-emission contributions to arbitrary cross sections through NNLO.

A HIGHLY-DEVELOPED THEORY OF PARTONIC COLLISIONS

A highly-developed theory of partonic collisions, that can be used to describe complicated collider processes, is available.

Leading order computations are automated; it is a solved problem. Madgraph etc.

Modern NLO computations are possible for processes with up to 6 final-state particles. They incorporate electroweak corrections and are often matched to parton showers allowing one to simulate realistic events. They include realistic final states (for unstable intermediate particles) and all interferences between (resonance) signal and (non-resonance) background.

Worek, Pozzorini, Denner etc.; OpenLoops etc.

NNLO QCD computations have become available for many interesting processes. The limiting factors currently are availability of virtual loop amplitudes and the efficiency of implementation of subtraction schemes in numerical codes.

Gehrmann, Glover, Huss, Czakon, Mitov, Poncelet, Williams, Roentsch, Caola, Catani, Grazzini

The current focus is on computing two-loop loop amplitudes for proceses with three (some massive) final-state particles.

$$pp \to V + jj, \; pp \to VV + j, \; pp \to t\bar{t}j, \; pp \to t\bar{t}H$$

First N3LO QCD computations appeared (Higgs cross section and rapidity distribution in gluon fusion, Drell-Yan cross section and rapidity distirbutions). Amplitudes for 2->2 paronic processes are known; current frontier are 2->2 amplitudes with one massless and one massive final-state particle.

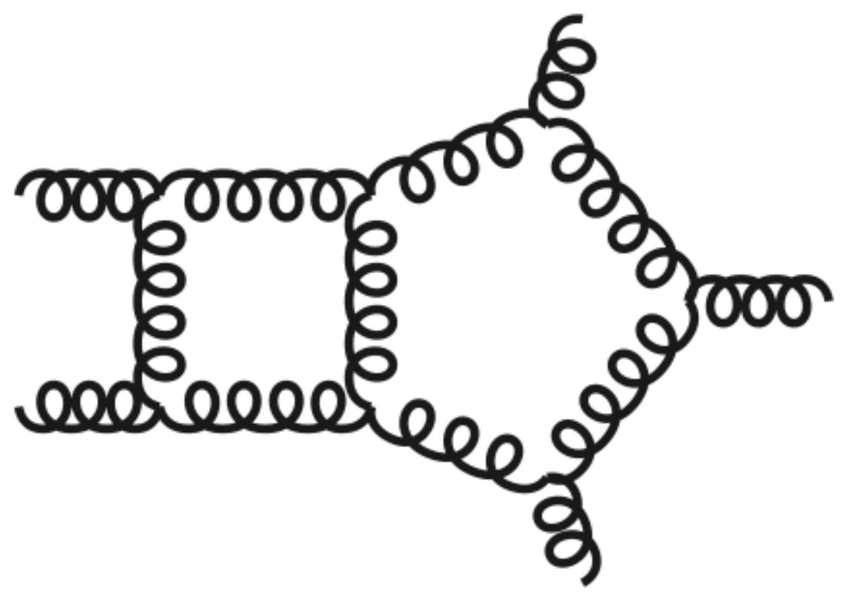
Anastasioiu, Duhr, Mistlberger, Gehrmann, Glover, Caola, Tancredi, Devoto, Buncioni

$$pp \to jj, pp \to \gamma\gamma \qquad pp \to Vj, pp \to Hj$$

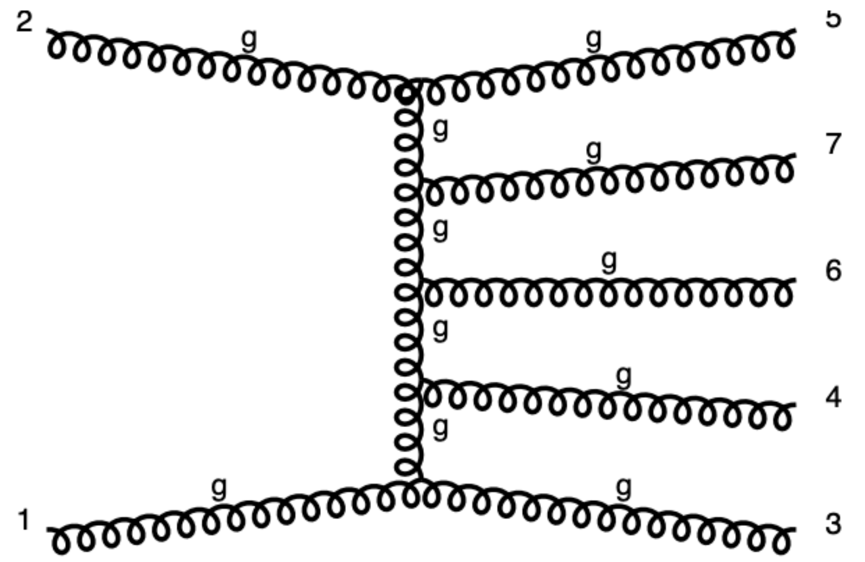
TECHNICAL PROGRESS LEADS TO BETTER PHENOMENOLOGY AND MORE PHYSICS OPPORTUNITIES

THE STRONG COUPLING CONSTANT AT THE HIGHEST ENERGIES

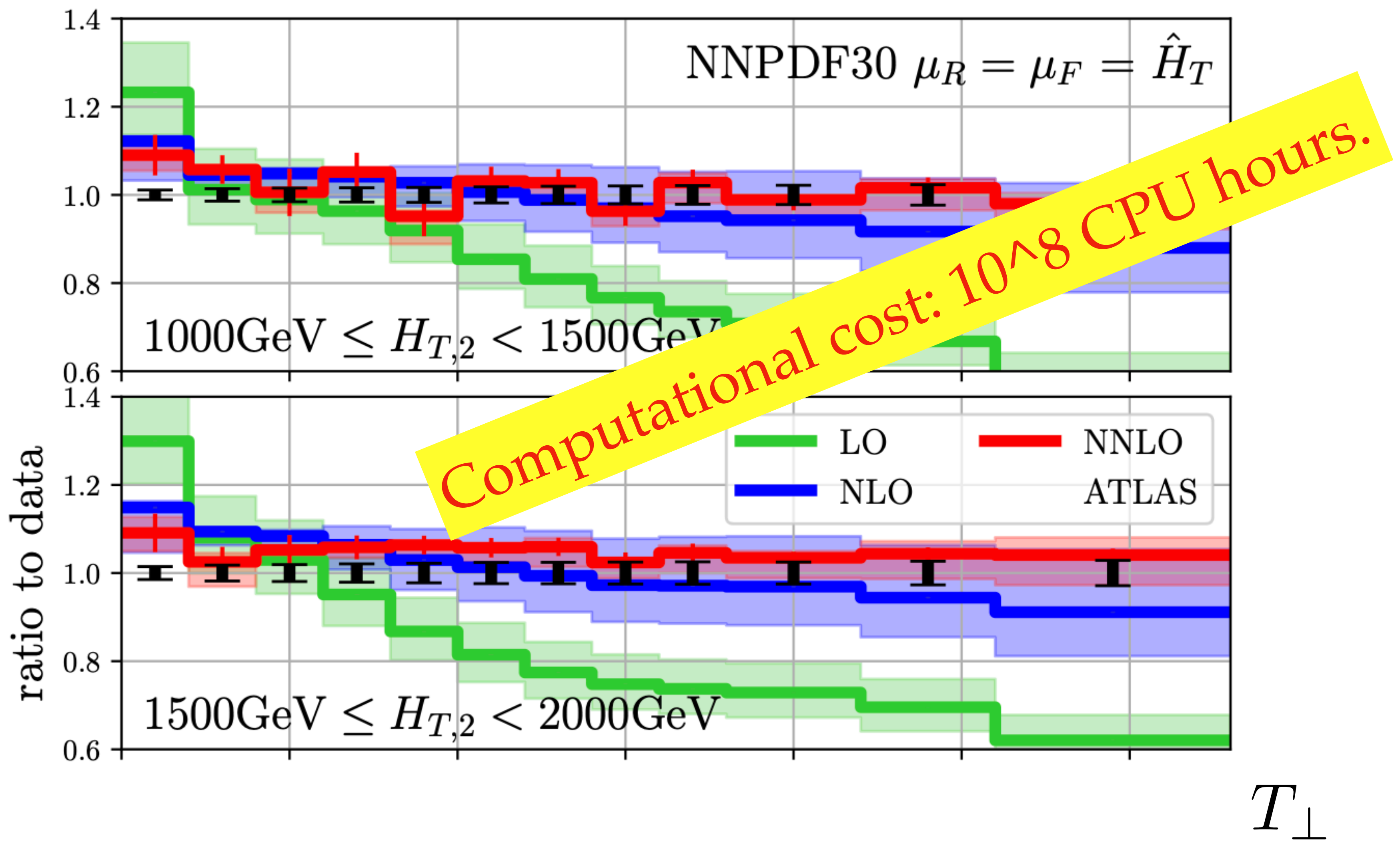
LHC experiments can measure the running of the strong coupling constant at very high energies. A useful observable is the transverse energy-energy correlator for 3j events. NLO results for this observable were known since quite some time. Pushing them to the next level — NNLO — was an enormous adventure.



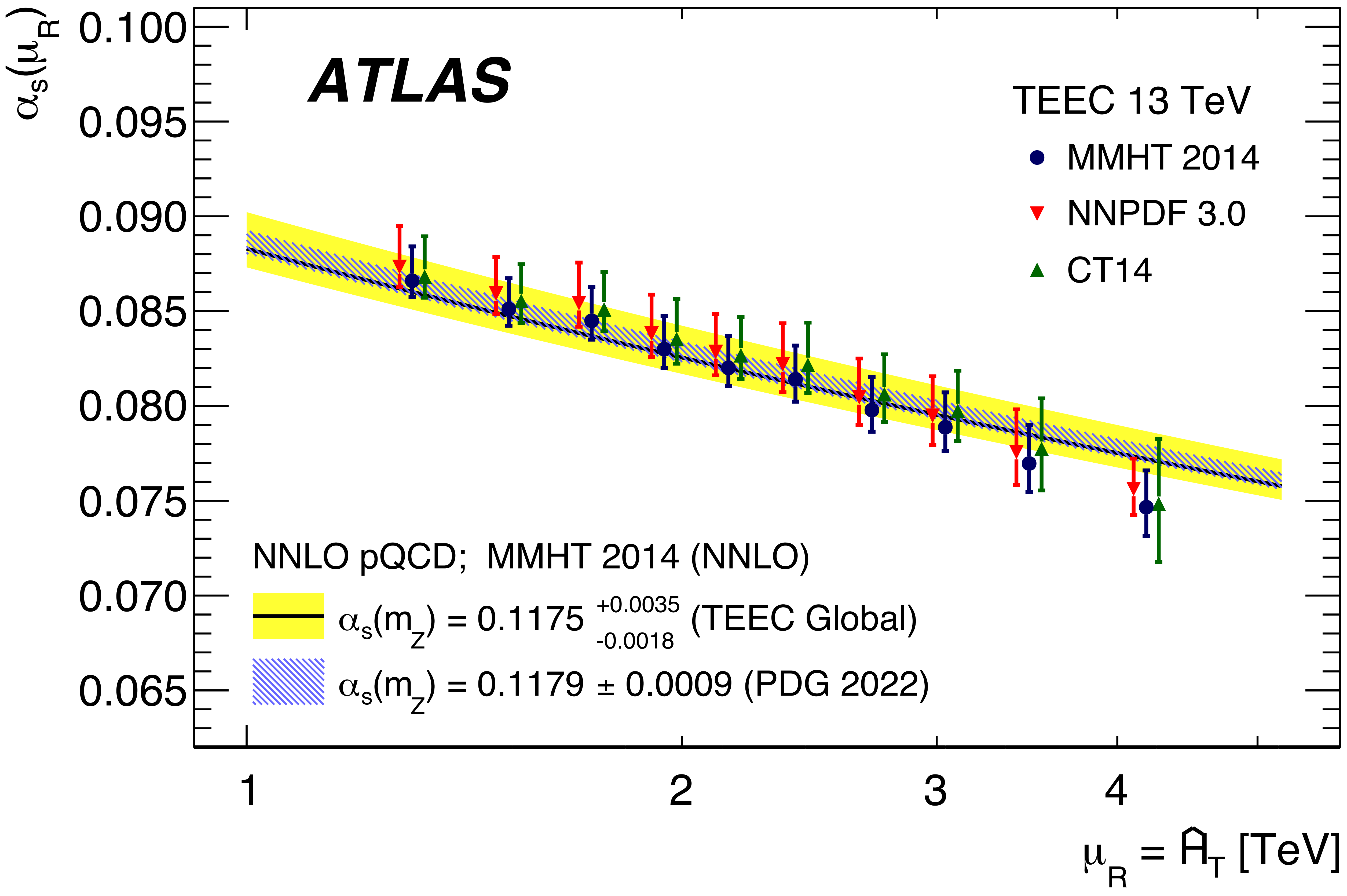
Two-loop amplitudes: Chicherin, Sotnikov, Gehrmann, Zhang, Henn, Wasser, Zola, Abreu et al.



Subtraction scheme: Czakon

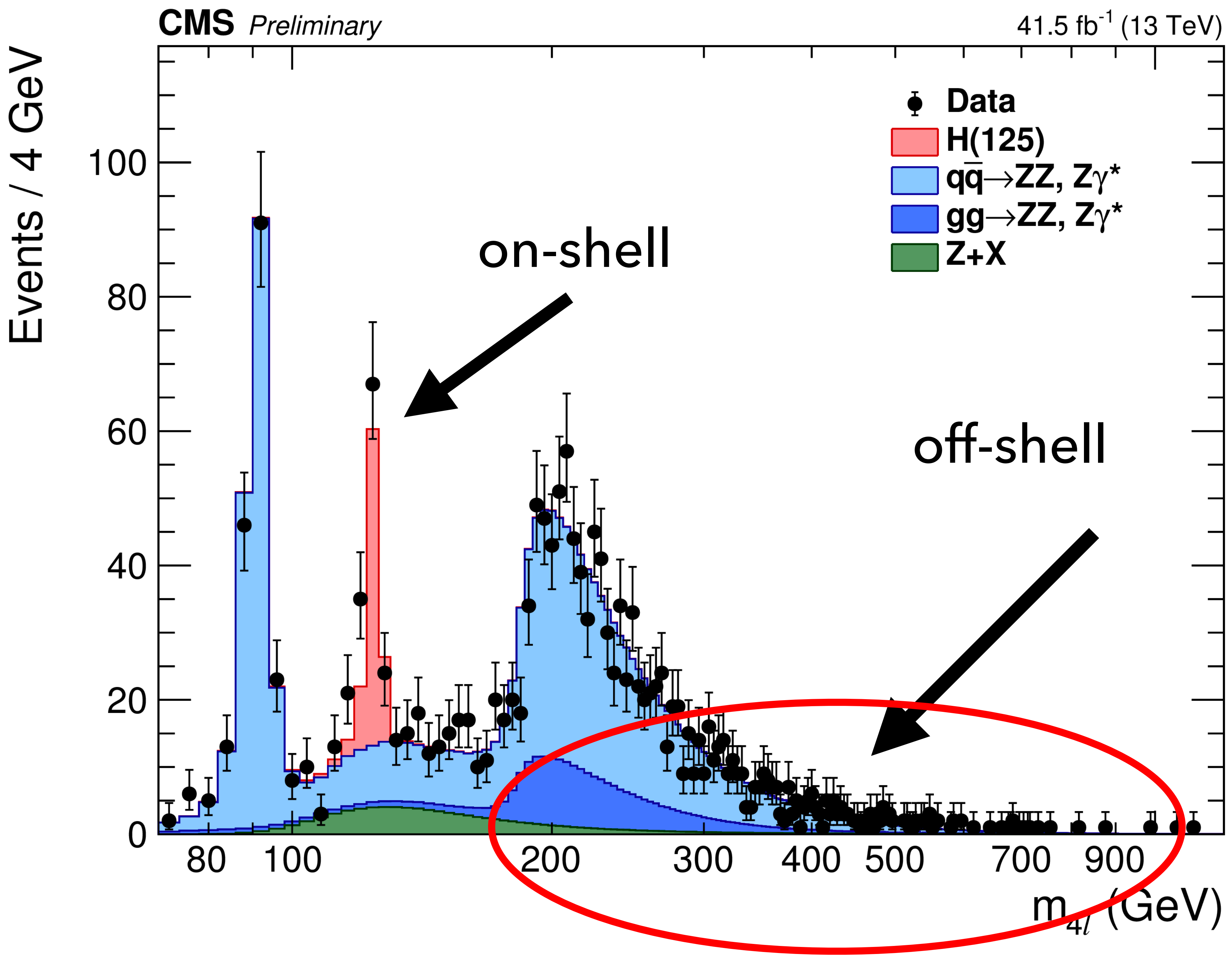


Alvarez, Cantero, Czakon, Lorente, Mitov, Poncelet

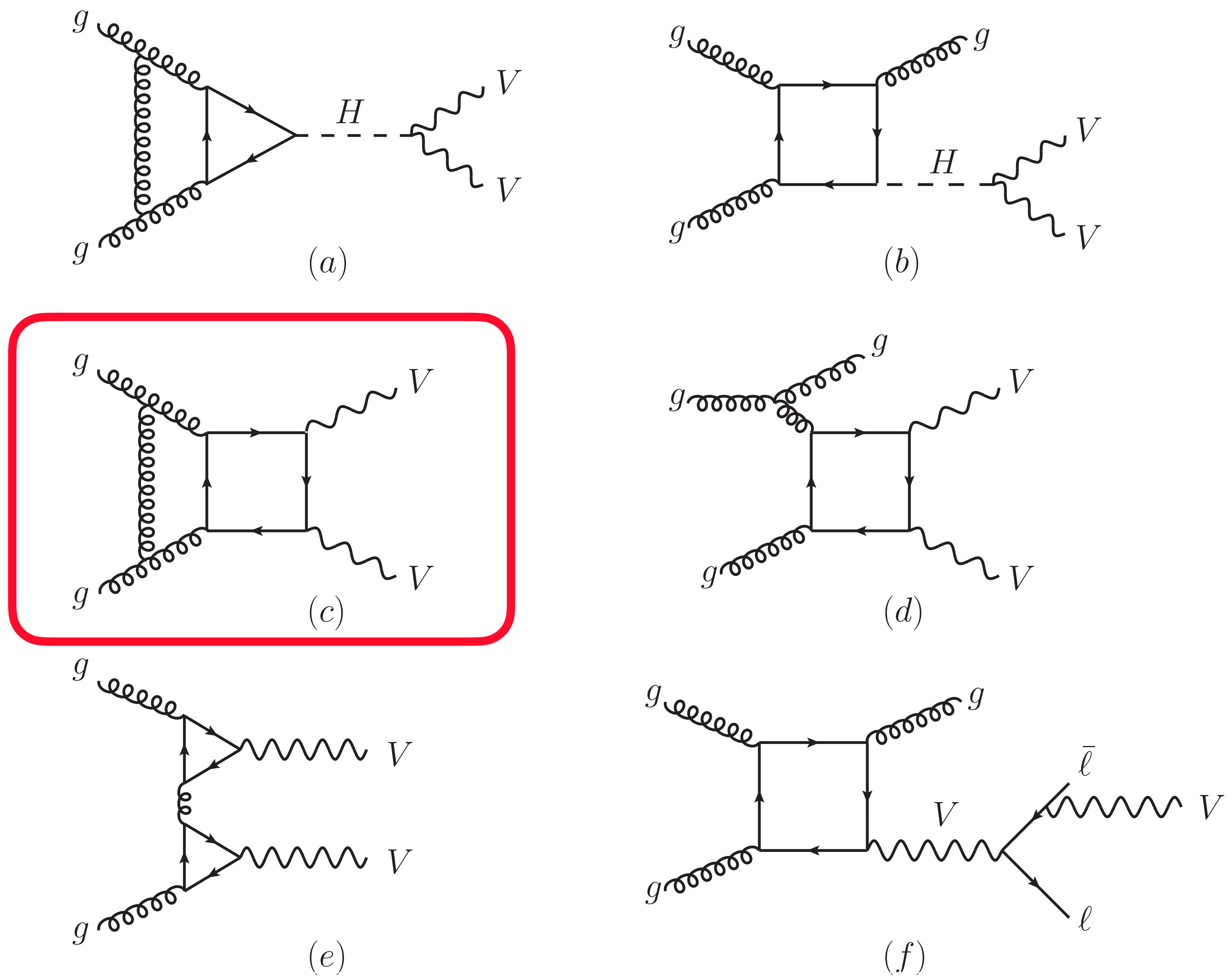


THE HIGGS WIDTH: FULL NNLO RESULTS FOR IRREDUCIBLE BACKGROUND

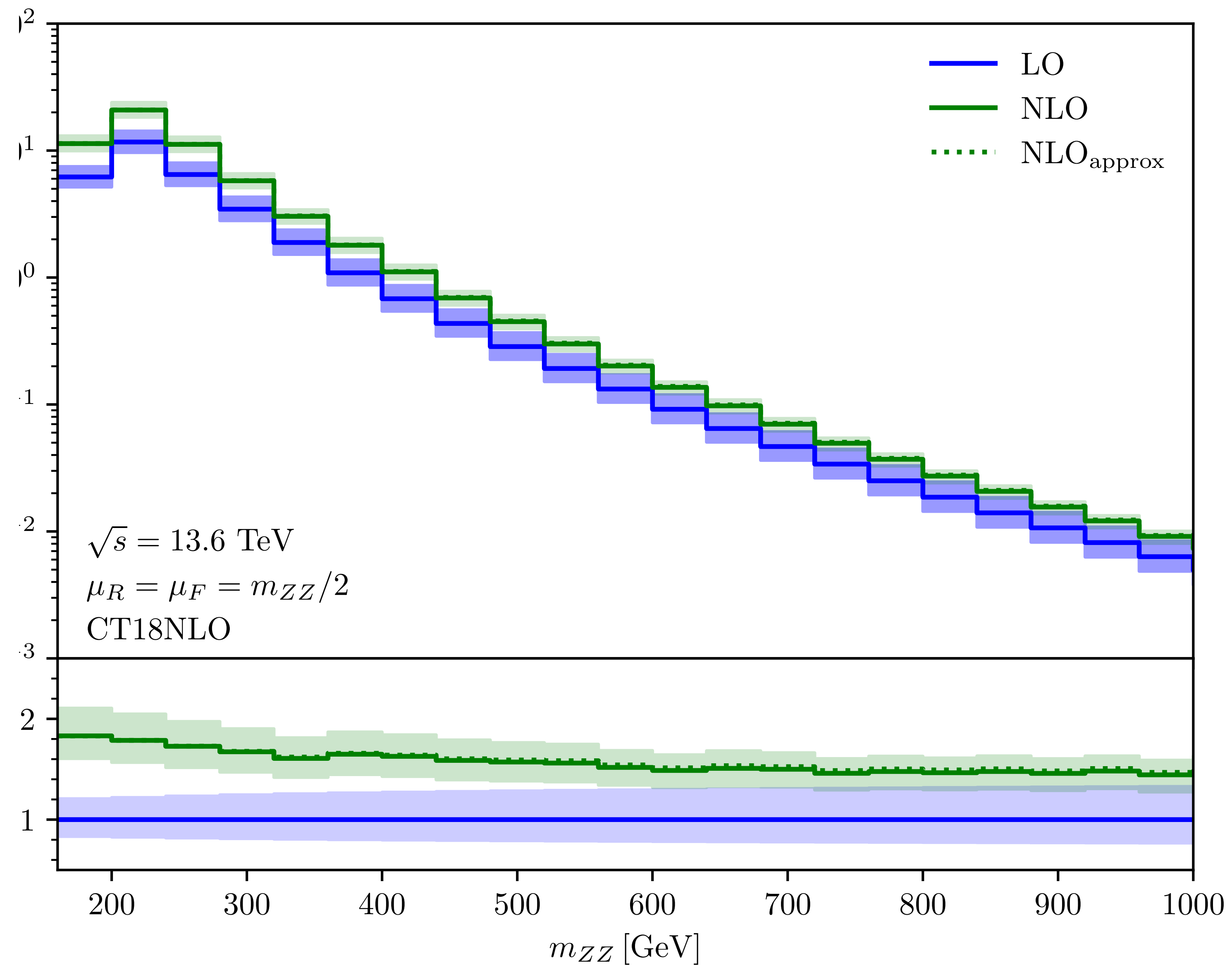
It is well-appreciated by now that one can extract the Higgs boson width from ZZ production using peculiar properties of the Higgs-boson off-shell contributions. Need to control the irreducible background; top-quark loop is (was) a challenge.



$$\Gamma_H \sim \frac{\sigma_{\rm off}}{\sigma_{\rm on}} \qquad \Gamma_H = 3.2^{+2.8}_{-2.2}\ \mathrm{MeV}$$



Both the signal and the irreducible background are part of the same process.



Class A$_h$: Both Z bosons couple directly to the same heavy top-quark loop. For these one- and two-loop contributions, we use the recent calculation [31] by some of us for which a combination of syzygy techniques [31, 36–40], finite field methods [41, 42], multivariate partial fractioning [43–47], and constructions of finite integrals were employed, and the resulting finite master integrals were evaluated numerically with `pySecDec` [48–50].

Agarwal, Jones, Kerner, von Manteuffel

STRONG COUPLING FROM Z TRANSVERSE MOMENTUM DISTRIBUTION

For a competitive measurement of the strong coupling at the LHC, one needs to find a quantity which

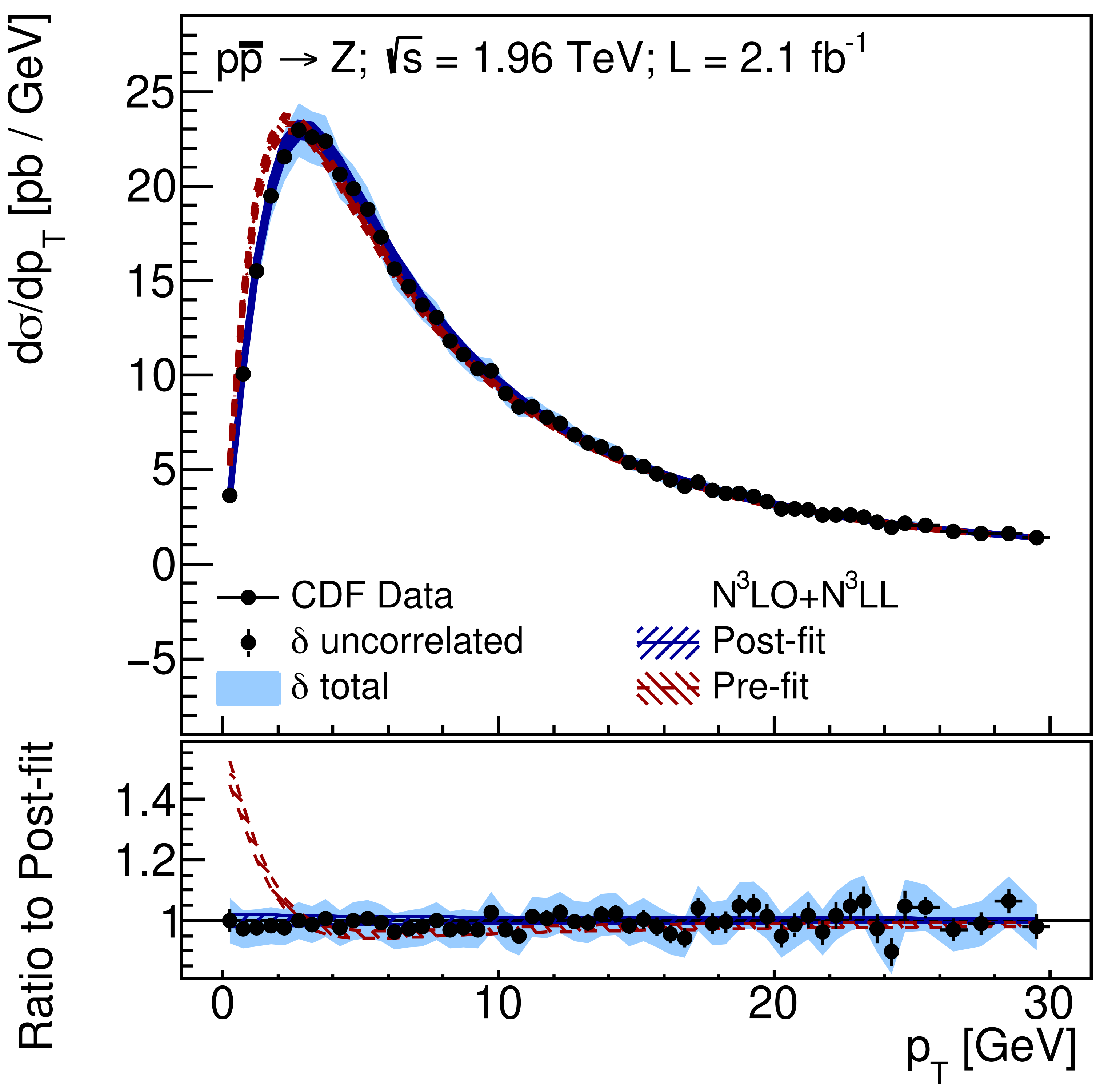
1) is proportional to the strong coupling constant;
2) can be predicted theoretically with a percent precision (NNLO and higher);
3) is independent (nearly independent) of poorly-known parton distribution functions;
4) refers to low(er) region of hard momentum region;
5) does not suffer from unknown non-perturbative effects.

Inclusive Z transverse momentum distribution seems to fit the bill.

$$\frac{\mathrm{d}\sigma_Z}{\sigma_z \mathrm{d}p_\perp} \sim \frac{\alpha_s(p_\perp)}{2\pi p_\perp} \ln \frac{M_Z}{p_\perp}$$

ATLAS followed up on the proposal and obtained a very precise value of the strong coupling constant which is very well-compatible with the world average.

$\alpha_s(m_z) = 0.1183 \pm 0.0009$ ATLAS, 8 TeV data

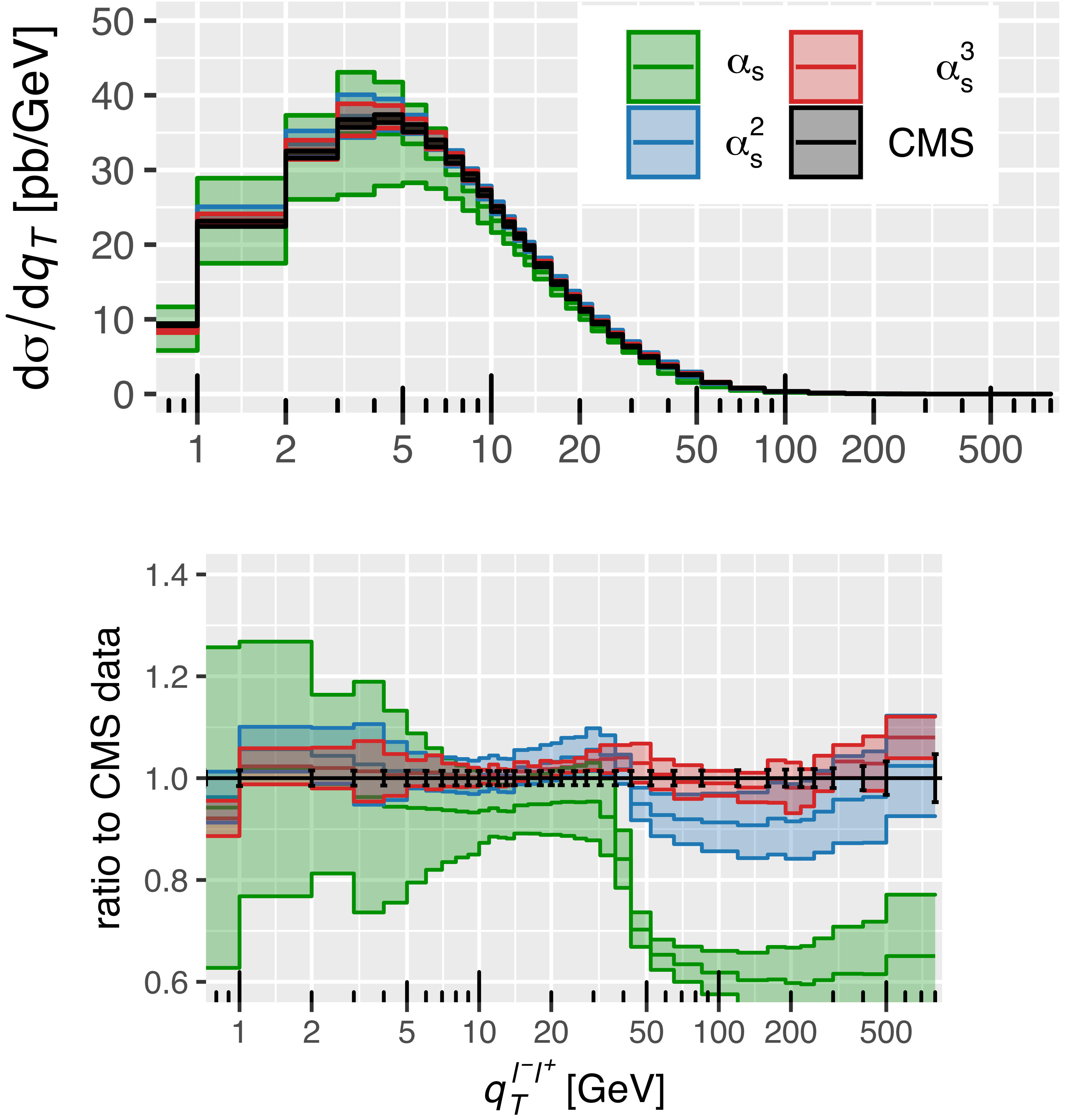


Camarada, Ferrera, Schott

STRONG COUPLING FROM Z TRANSVERSE MOMENTUM DISTRIBUTION

A percent-level prediction for Z transverse momentum distribution e requires us to employ some of the most sophisticated theoretical computations ever performed in pQCD.

1) N3LO QCD predictions for the inclusive Z-boson production cross section and rapidity distribution;
2) NNLO QCD predictions for Z+jet production;
3) state-of-the-art transverse momentum resummation,   that describe Z-boson transverse momentum distribution at small pt ;
4) electroweak corrections to Z+jet production;
5) advanced knowledge of parton distribution functions;
6) models for non-perturbative smearing at small transverse momenta.



Duhr, Mistlberger, X. Chen, Gehrmann, Gehrmann-de Ridder, Glover, Zhu, Yang, Huss, Vita, Ebert, Luou,. Boughezal, Focke, Liu, Petriello, Ellis, Giele, Campbell et al.

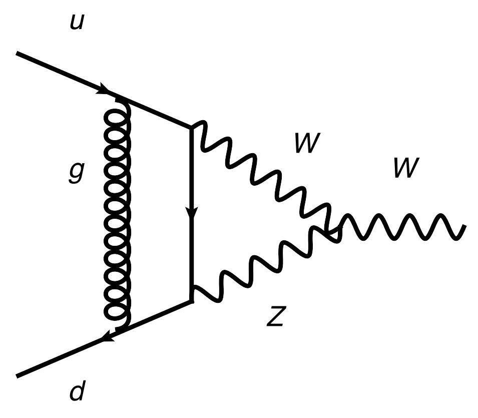
$\alpha_s(m_z) = 0.1183 \pm 0.0009$ ATLAS, 8 TeV data

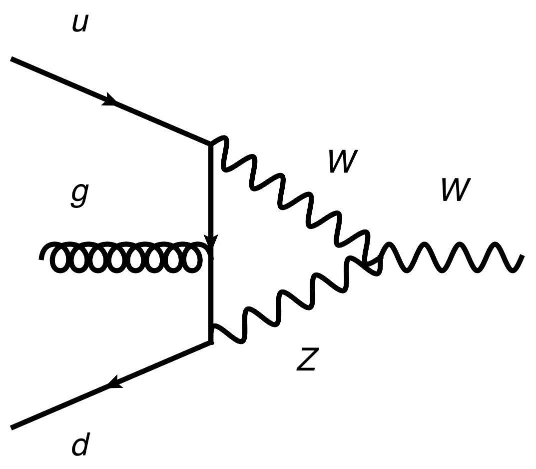
Neumann and Campbell

W-MASS AND MIXED QCD-ELECTROWEAK CORRECTIONS

To minimize the impact of QCD theory on the determination of the W-mass, models for vector boson production are tuned using Z-production data and then used to describe the W case. It becomes important to carefully study all effects that distinguish between Z and W production and electroweak and mixed electroweak -QCD corrections is an important example of such effects.

$$m_W^{\rm meas} = \frac{\langle p_\perp^{l,W}\rangle^{\rm meas}}{\langle p_\perp^{l,Z}\rangle^{\rm meas}}\, m_Z\, C_{\rm th}. \qquad C_{\rm th} = \frac{m_W}{m_Z}\frac{\langle p_\perp^{l,Z}\rangle^{\rm th}}{\langle p_\perp^{l,W}\rangle^{\rm th}}$$





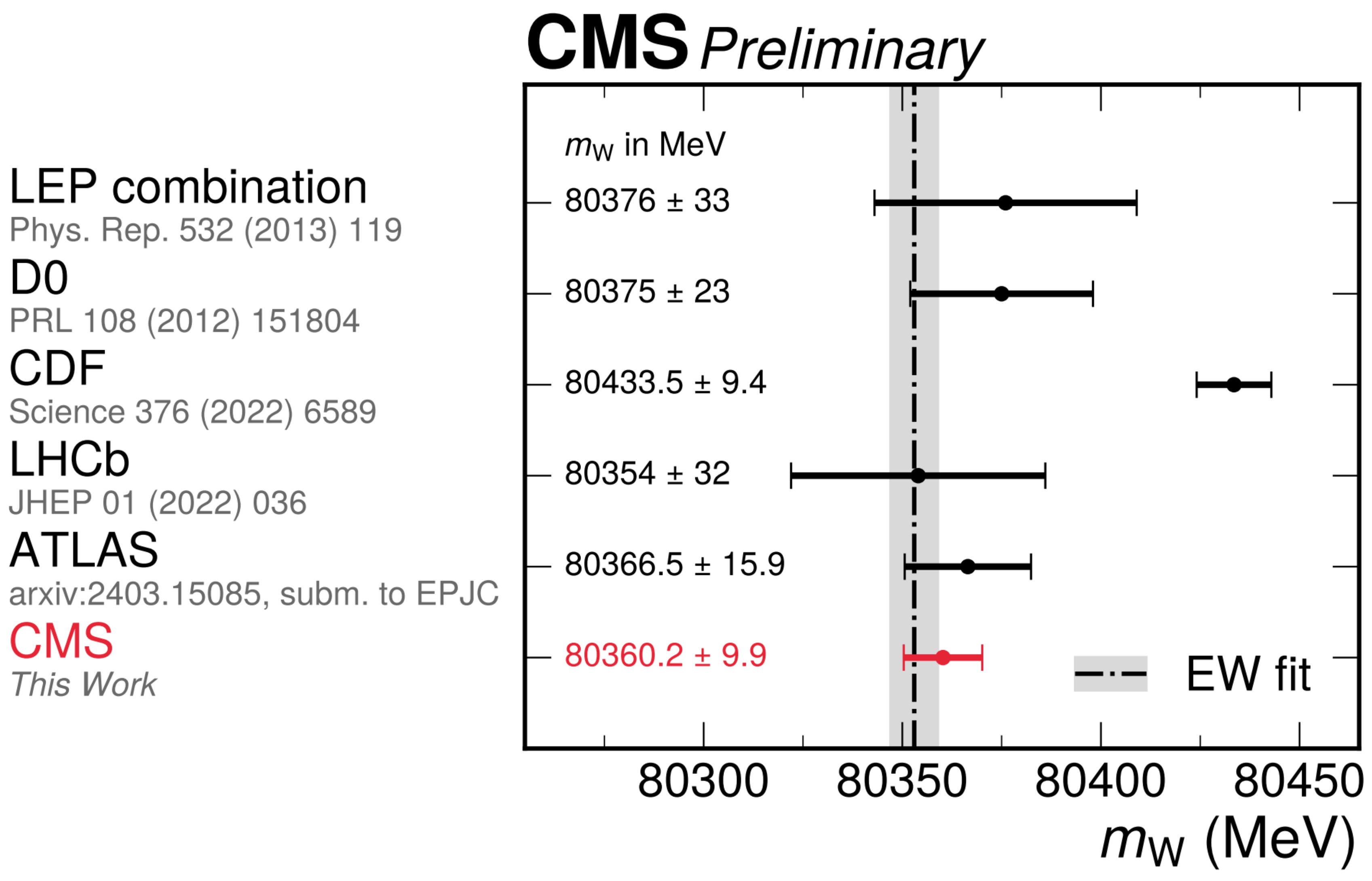
A better theory changes the theoretical correction factor and leads to changes in the extracted value of the W mass.

$$\frac{\delta m_W^{\rm meas}}{m_W^{\rm meas}} = \frac{\delta C_{\rm th}}{C_{\rm th}} = \frac{\delta\langle p_\perp^{l,Z}\rangle^{\rm th}}{\langle p_\perp^{l,Z}\rangle^{\rm th}} - \frac{\delta\langle p_\perp^{l,W}\rangle^{\rm th}}{\langle p_\perp^{l,W}\rangle^{\rm th}}$$

No fiducial cuts: $\Delta m_W = m_W - m_W^{EW} = 7\ {\rm MeV}$

- QCD-electroweak effects are more important than the electroweak ones;
- Compensation mechanism between W and Z distribution is important; in first moments taken separately are close to 50 MeV;
- PDF uncertainty has a very minor impact on these shifts;

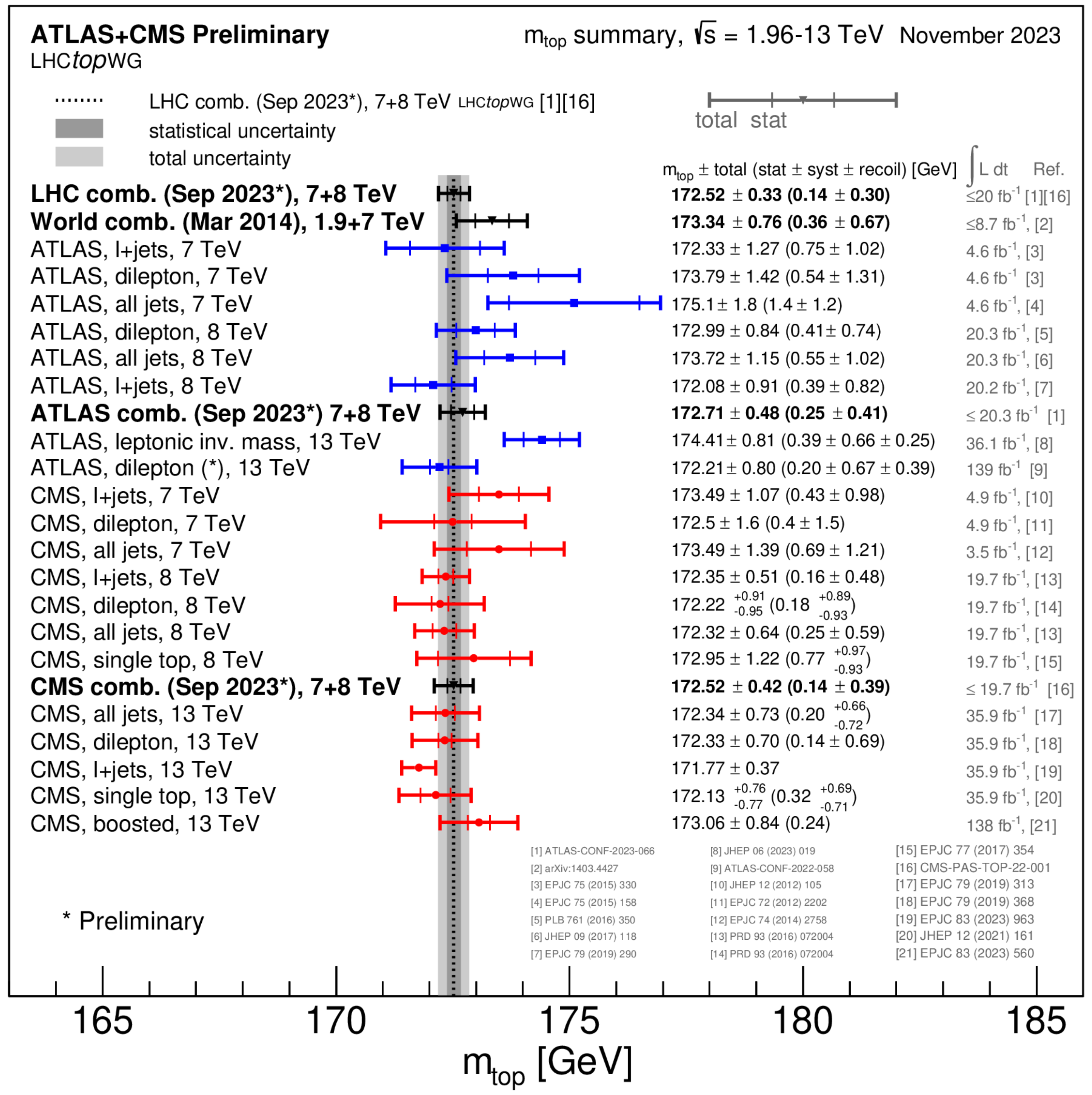
ATLAS cuts: $\Delta m_W = m_W - m_W^{EW} = 17\ {\rm MeV}$

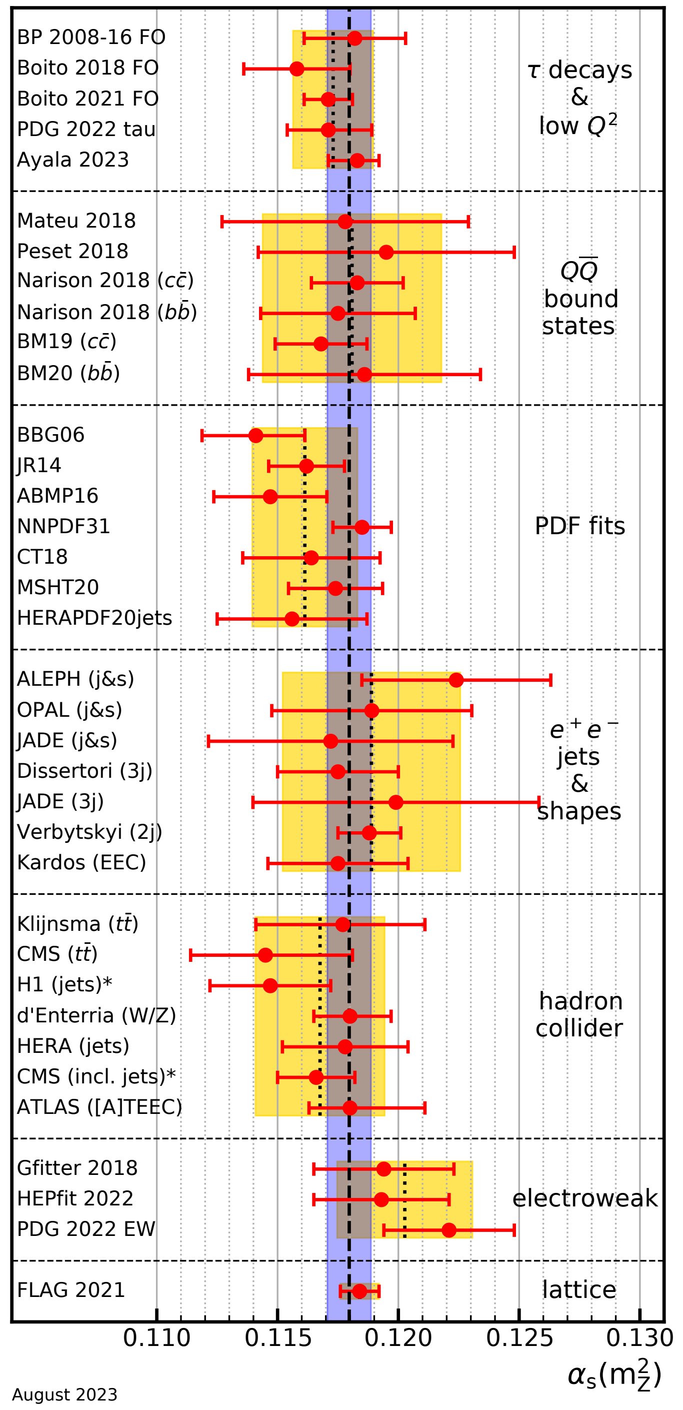


Behring, Buccioni, Caola, Delto, Melnikov, Jaquier, Roentsch

THE CONCEPTUAL PROBLEM OF NON-PERTURBATIVE POWER CORRECTIONS

Modelling non-perturbative effects with parton showers is not satisfactory for high-precision observables. It is known to cause significant confusion.





$$m_t = 172.52 \pm 0.33 \text{ GeV}$$

$$\alpha_s(M_z) = 0.118 \pm 0.001$$

$$\mathrm{d}\sigma_{\text{hard}} = \int \mathrm{d}x_1\, \mathrm{d}x_2 f_i(x_1) f_j(x_2)\, \mathrm{d}\sigma_{ij}(x_1, x_2, \{p_{\text{fin}}\})\, O_J(\{p_{\text{fin}}\})\left(1 + \mathcal{O}(\Lambda^n_{\text{QCD}}/Q^n)\right)$$

POWER CORRECTIONS

Can one learn something relevant about these effects from perturbation theory given all the advances that we have had in this field?

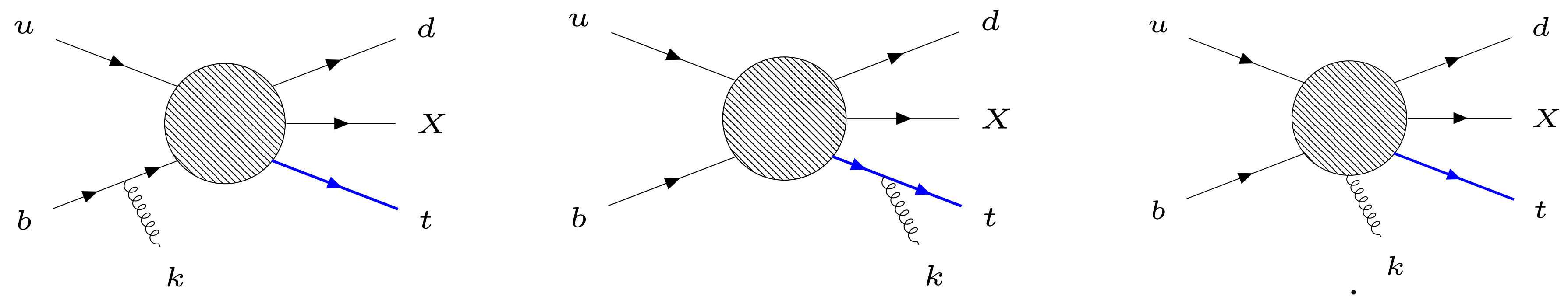
A recent discussion of inter-dependences between the perturbative evolution of parton showers and the hadronization models through a shower infra-red cut-off is an interesting example of this.

Hoang, Jin, Plätzer, Samitz

Furthermore, since Feynman integrals run over all momenta, including the soft ones, one can use Feynman diagrams to estimate the sensitivity of cross sections and observables to these problematic integration regions.

The famous Kinoshita-Lee-Nauenberg infra-red cancellation, as well as the idea of renormalons and its connection to QCD with a (fake) gluon mass can be interpreted in this way.

Calculation of linear $\mathcal{O}(\Lambda_{\rm QCD})$ non-perturbative corrections in the context of renormalon models can be simplified using Low-Burnett-Kroll next-to-soft-emission theorem and some tricks from the perturbative toolbox.



$$|\mathcal{M}(..,k)|^2 = -J^\mu J_\mu F_{\rm LO}(q_t,p_b,q_d,p_u) - J_\mu L^\mu F_{\rm LO}(q_t,p_b,q_d,p_u) + \mathcal{O}(k^0)$$

The approach based on renormalons has its limitations but it also leads to important insights into non-perturbative effects that are listed below:

1) one cannot determine the pole mass of the top quark from top production cross section with a precision better than $\mathcal{O}(\Lambda_{\rm QCD})$;

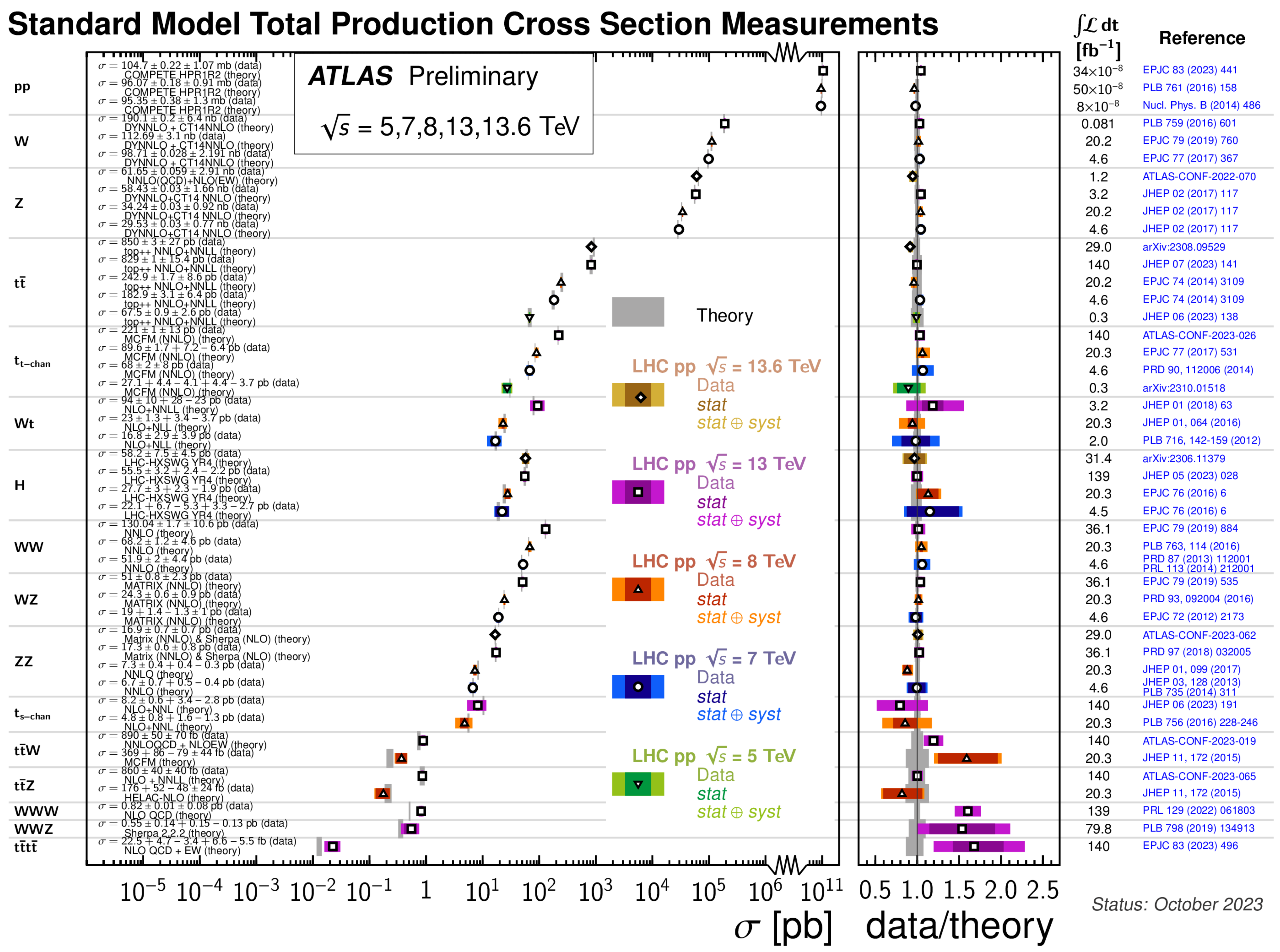
2) even basic kinematic distributions in top-production processes receive linear power corrections independent of the top mass parameter used; these power corrections are not described by parton showers;

3) polarization effects in top quark production processes are affected by linear power corrections (in the narrow width approximation);

4) in electron-positron collisions, non-perturbative corrections to shape variables in 3-jet and 2-jet regions are different, in variance with the standard assumption that are made when fitting the strong coupling constant α_s. Similar effects should exist at hadron colliders.

Ferrario Ravasio, Limatola, Nason, Caola, Melnikov, Ozcelik, Makaroc

HL-LHC AS AN ULTIMATE PRECISION MACHINE

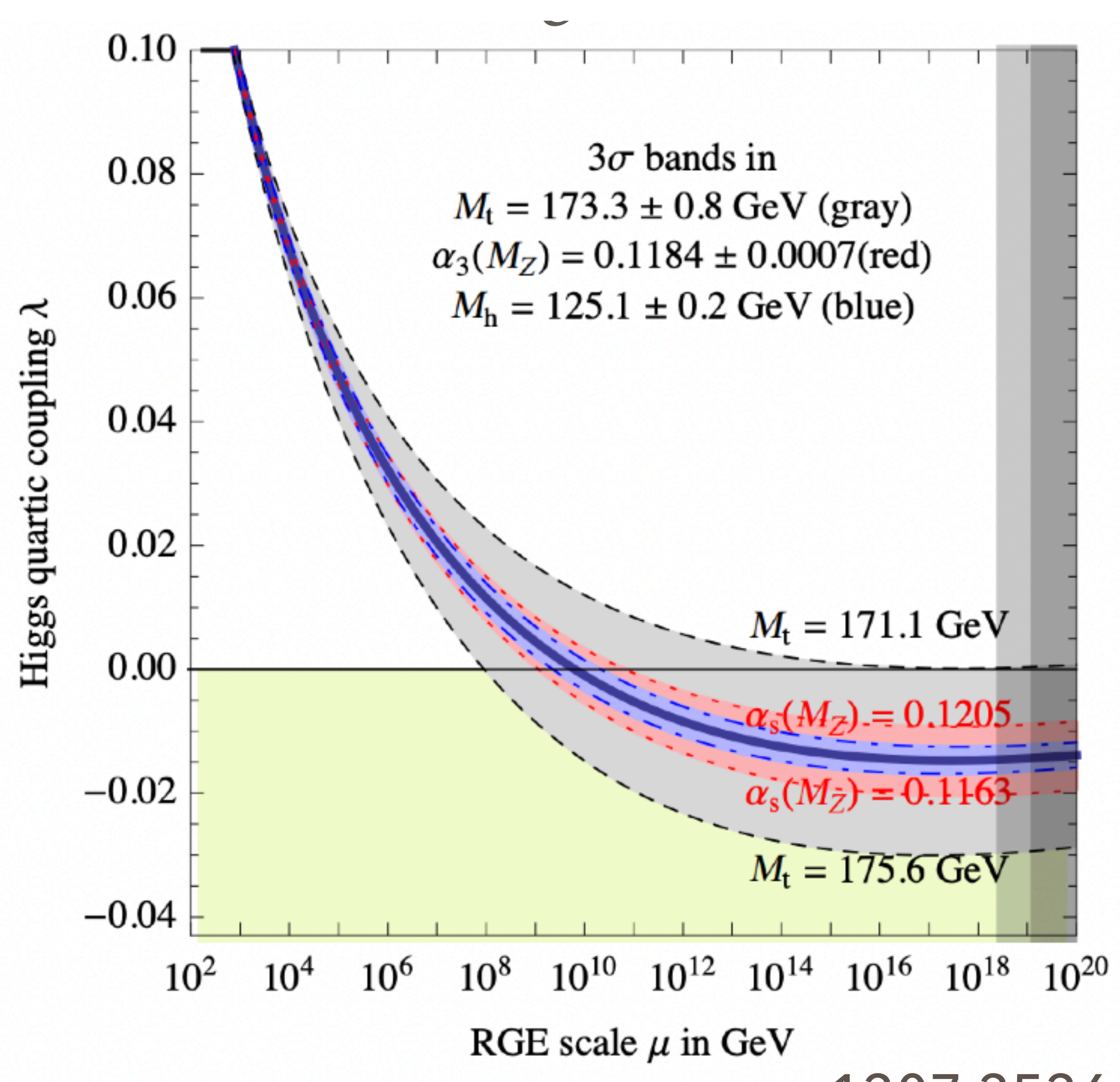


LHC AS THE ULTIMATE PRECISION MACHINE

Measurement of the top quark mass which combines the high(est) precision with solid estimates of non-perturbative effects;

Measurement of the W-boson mass with precision comparable to outcomes of the precision electroweak fit;

Measurement of the strong coupling constant at the LHC measurements with the below 1 percent precision;



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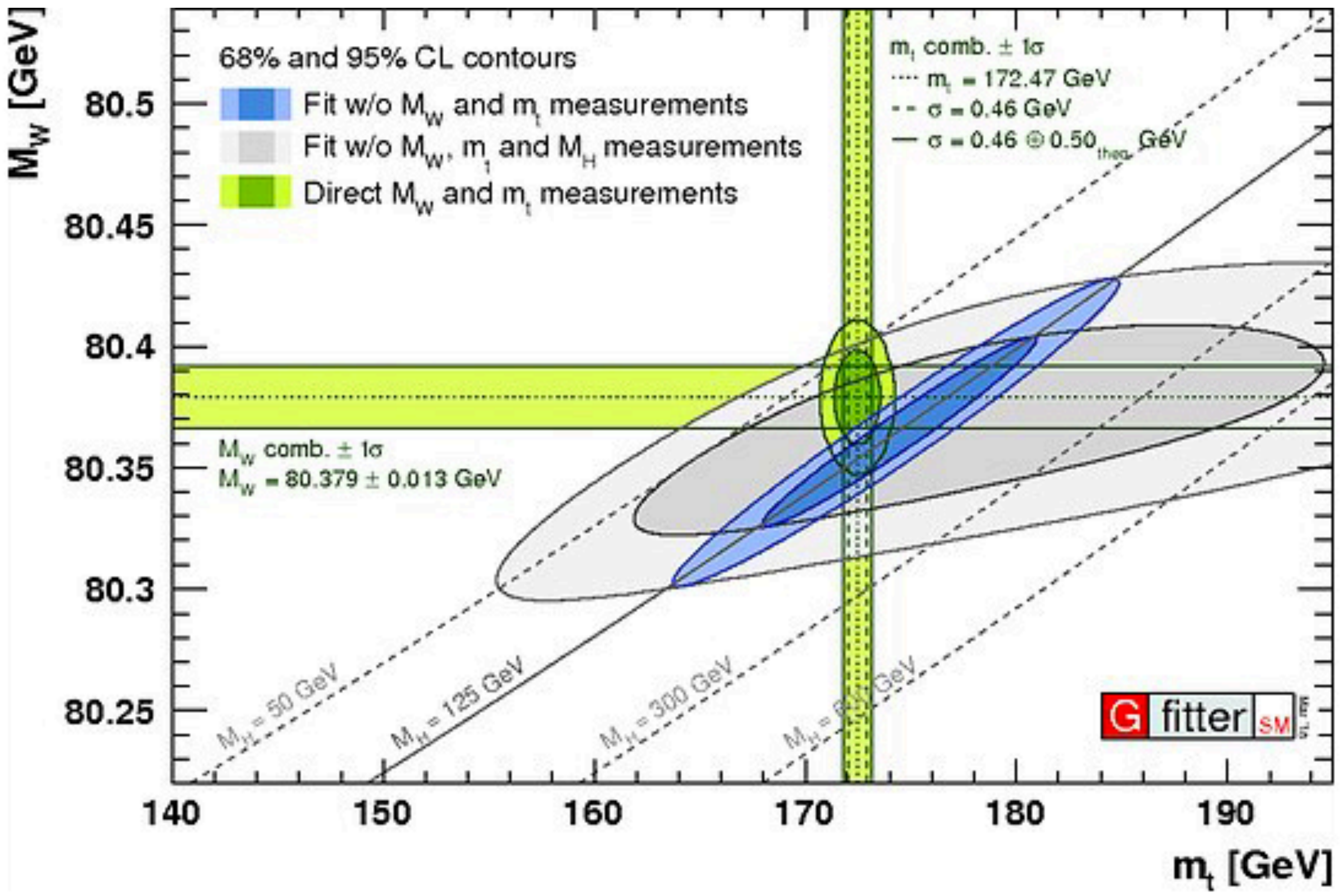
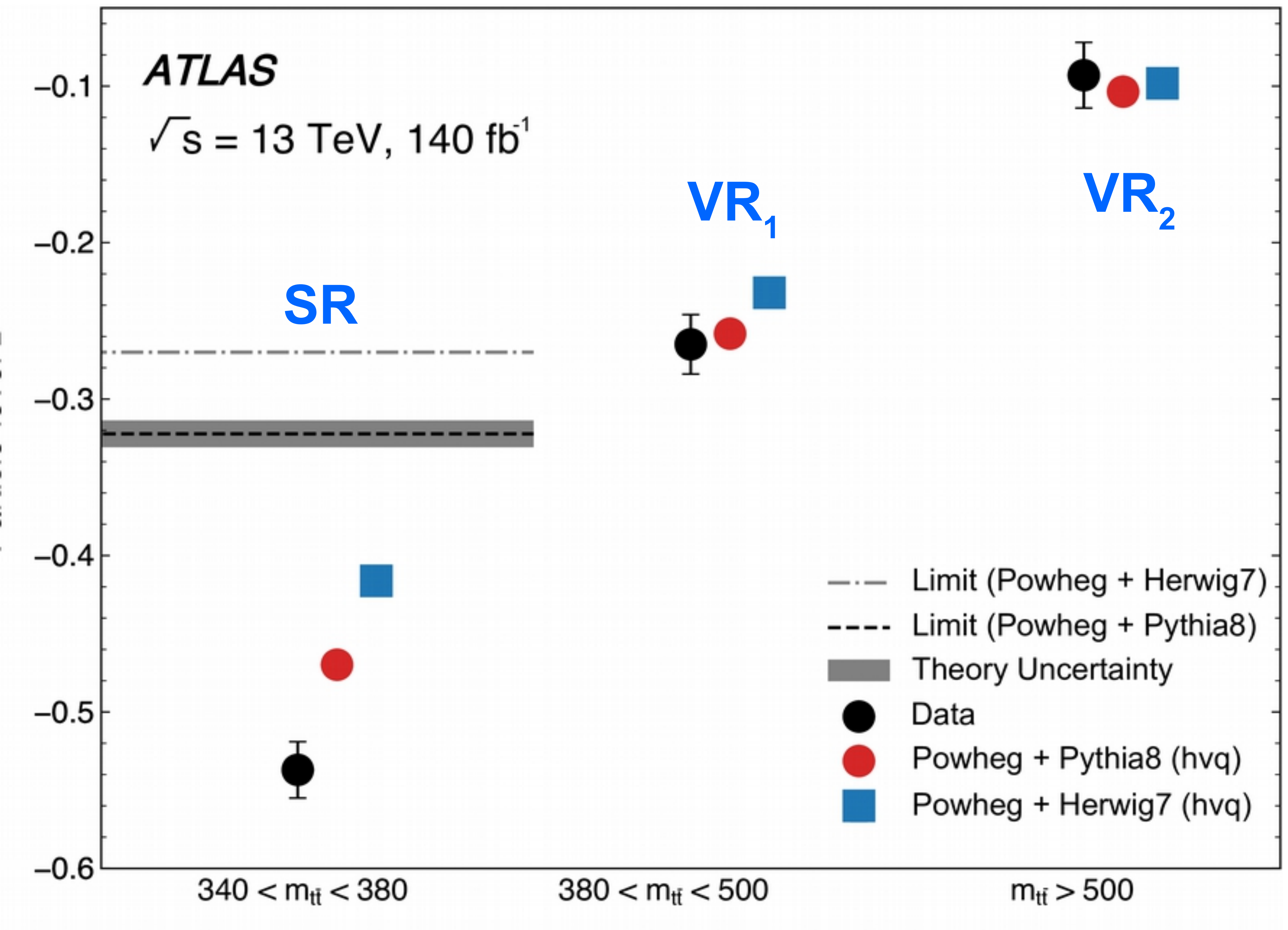
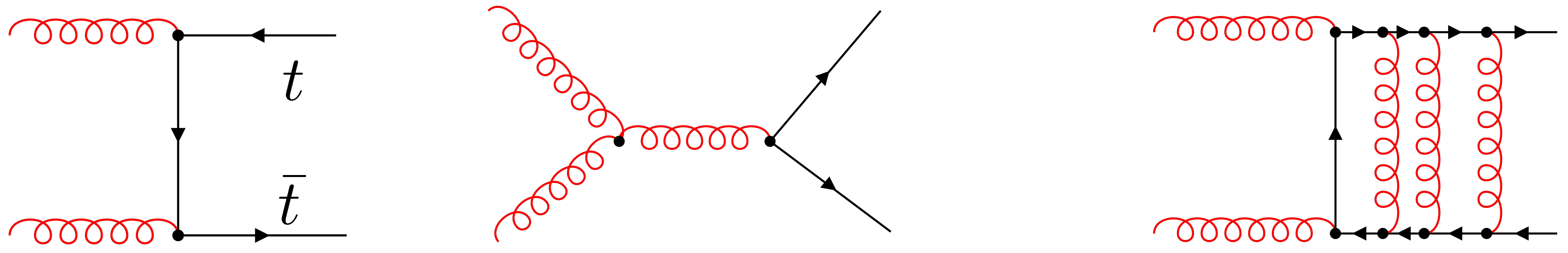


TOP QUARK SPIN-CORRELATIONS /ENTANGLEMENT/THRESHOLD BEHAVIOUR

Recently, top-quark spin correlations as a function of top pair invariant mass were studied at the LHC. An enhancement of correlations in the threshold region was observed. The threshold region is special — it is a place were bound state effects of top and anti-top exist; they can be described in QCD from first principles. Can effects of top-antitop bound states be observed at the LHC?



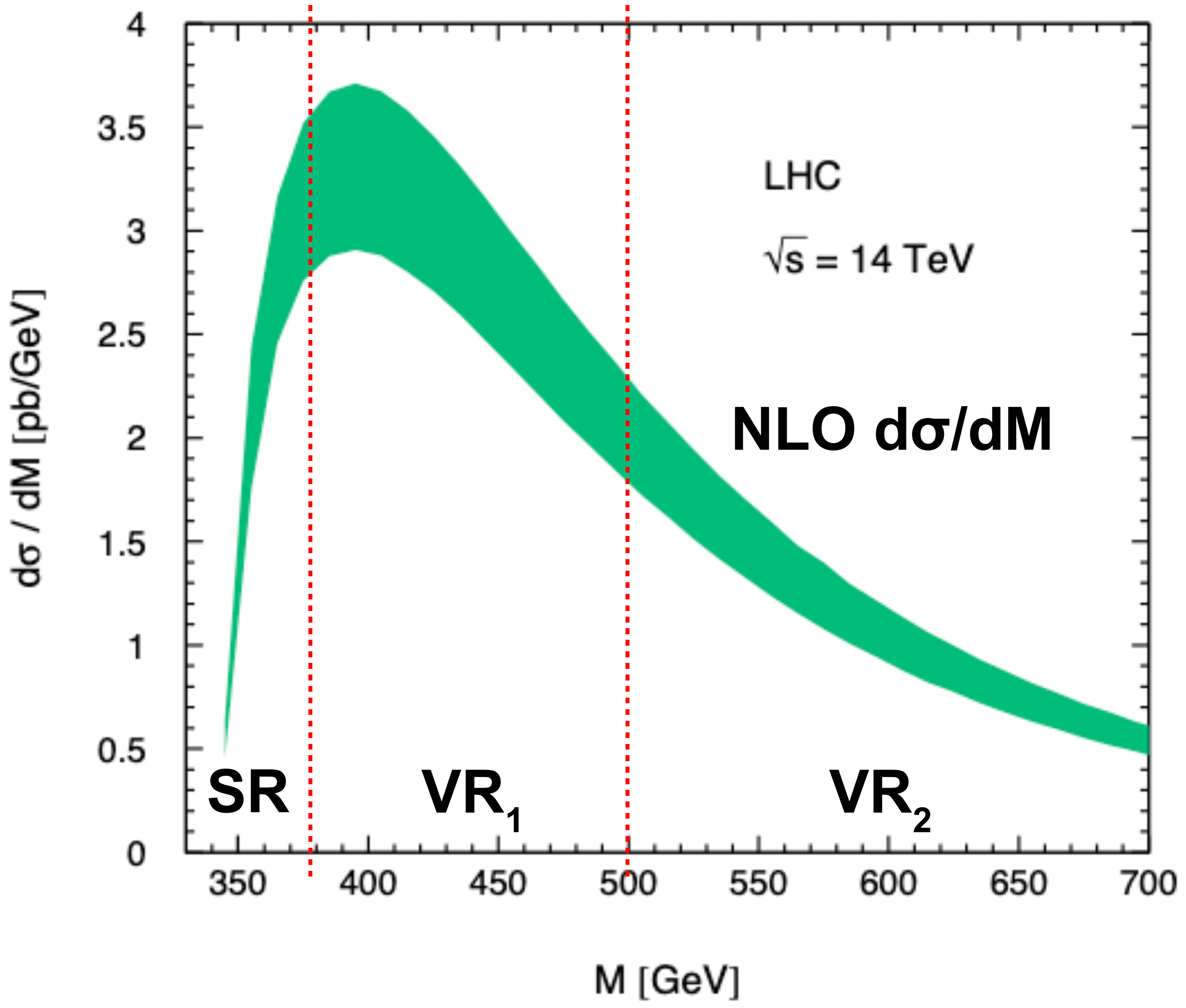
Nature 633 (2024) 542

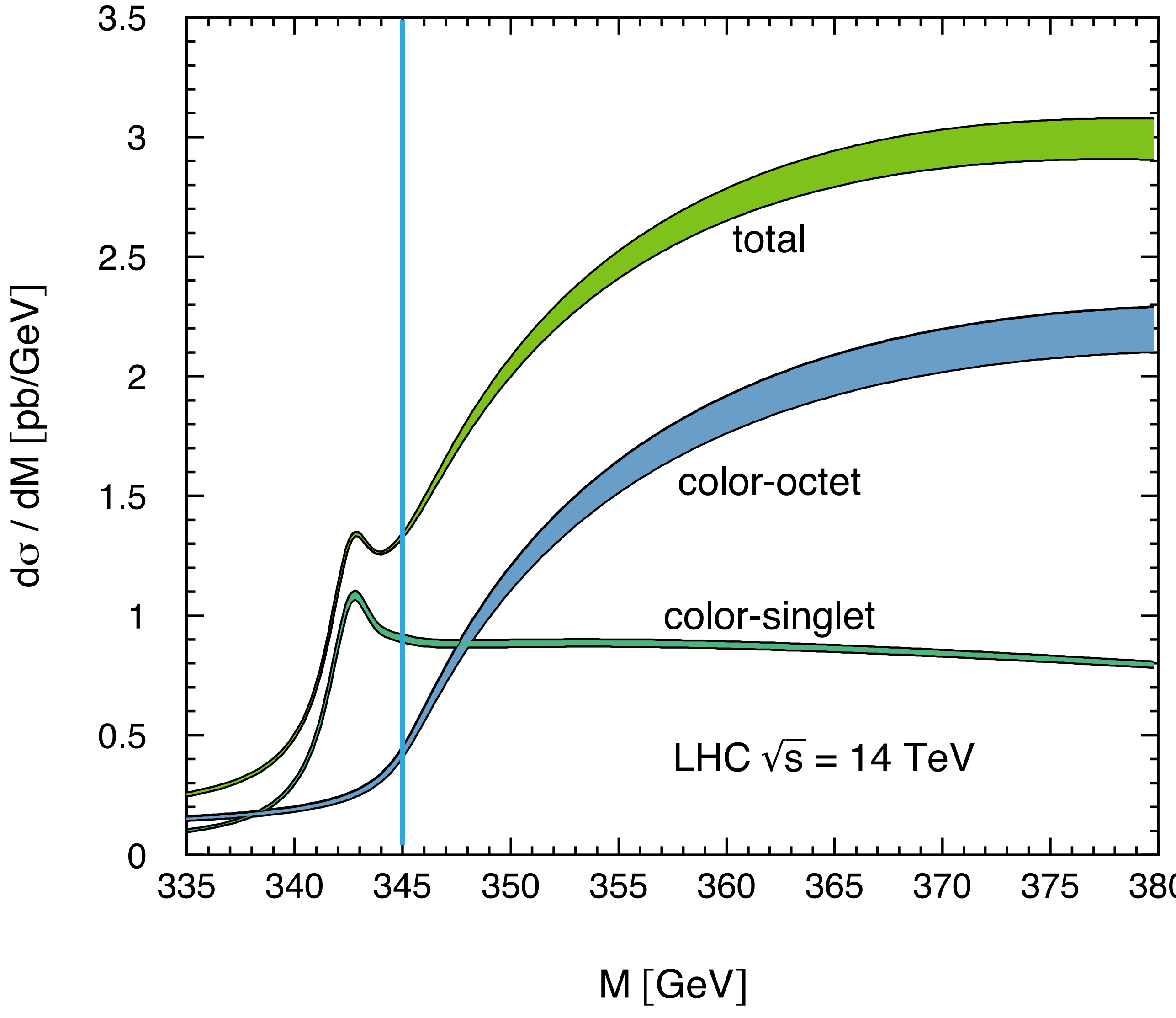


$$\Rightarrow \sum_{n=0}^{\infty} \frac{\alpha_s^n}{\beta_{t\bar{t}}^n}$$

$$\Psi_{t\bar{t}} = \frac{1}{\sqrt{2}} \left(|+-\rangle - |-+\rangle \right)$$

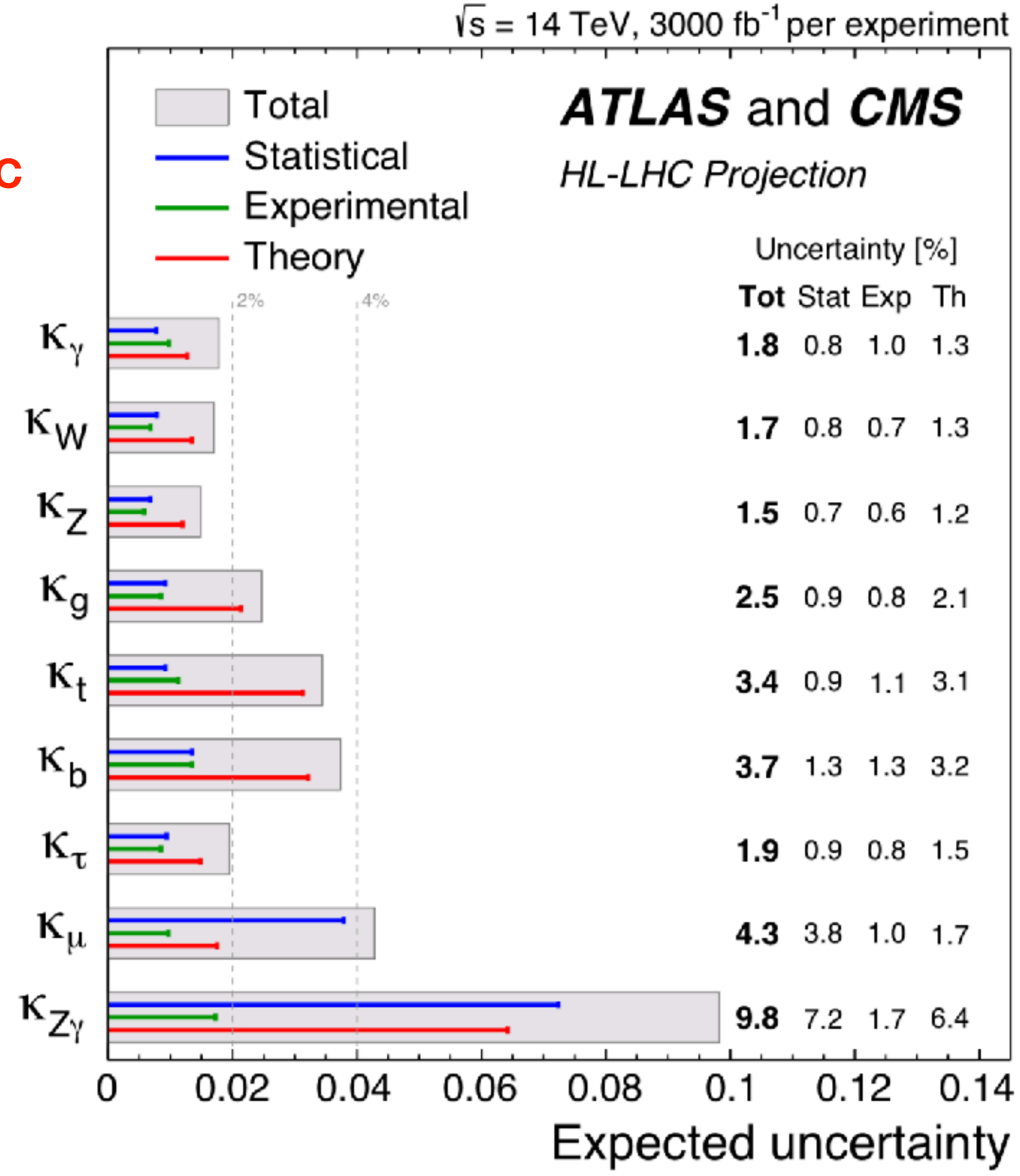
At the threshold, top quarks in a color-singlet channel will have the zero-spin wave function which means that spins of top and anti-top are fully correlated.





THE ULTIMATE PRECISION FOR HIGGS BOSON COUPLINGS

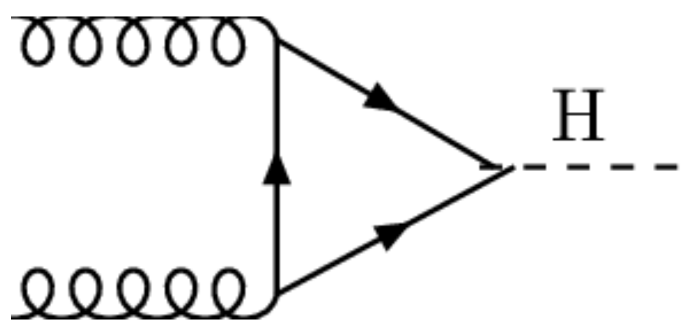
	ATLAS - CMS Run 1 combination	ATLAS Run 2	CMS Run 2	Current precision	HL-LHC
κ_γ	13%	1.04 ± 0.06	1.10 ± 0.08	6%	1.8%
κ_W	11%	1.05 ± 0.06	1.02 ± 0.08	6%	1.7%
κ_Z	11%	0.99 ± 0.06	1.04 ± 0.07	6%	1.5%
κ_g	14%	0.95 ± 0.07	0.92 ± 0.08	7%	2.5%
κ_t	30%	0.94 ± 0.11	1.01 ± 0.11	11%	3.4%
κ_b	26%	0.89 ± 0.11	0.99 ± 0.16	11%	3.7%
κ_τ	15%	0.93 ± 0.07	0.92 ± 0.08	8%	1.9%
κ_μ	-	$1.06^{+0.25}_{-0.30}$	1.12 ± 0.21	20%	4.3%
$\kappa_{Z\gamma}$	-	$1.38^{0.31}_{-0.36}$	1.65 ± 0.34	30%	9.8%
B_{inv}		$< 11\,\%$	$< 16\,\%$	11%	2.5%
		Nature 607, 52-59 (2022)	Nature 607, 60-68 (2022)		



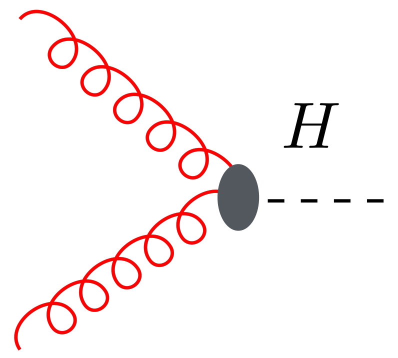
TH Uncertainties dominant (assumed to be 1/2 of Run 2)

Lectures by Marumi Kado at Maria Laach Summer School, 2024.

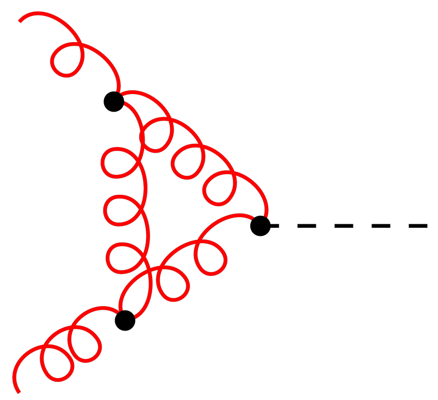
THE SIMPLEST HIGGS BOSON PRODUCTION CROSS SECTION

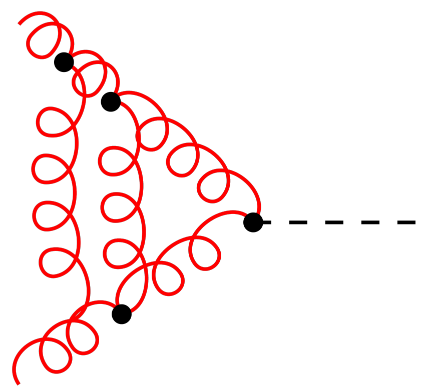


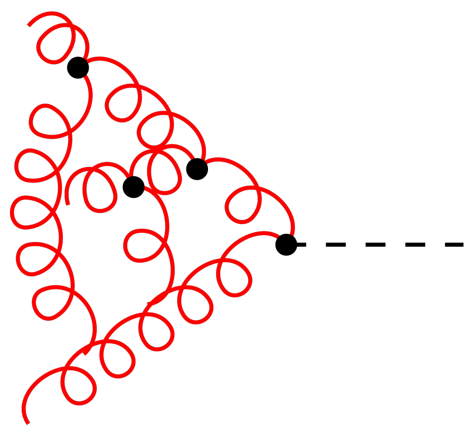
$\Rightarrow$



$$\boxed{\sigma = 48.58\,\text{pb}^{+2.22\,\text{pb}\,(+4.56\%)}_{-3.27\,\text{pb}\,(-6.72\%)}\ (\text{theory}) \pm 1.56\,\text{pb}\,(3.20\%)\ (\text{PDF}+\alpha_s)\,.}$$





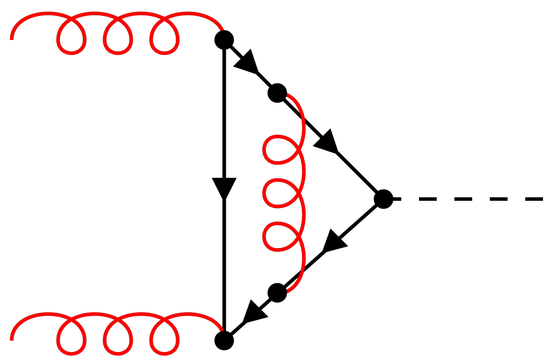


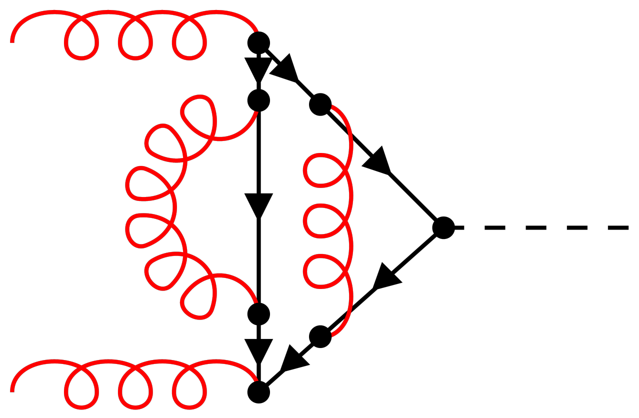
$$\begin{aligned}
48.58\,\text{pb} = \;& 16.00\,\text{pb} && (+32.9\%) && (\text{LO, rEFT})\\
&+20.84\,\text{pb} && (+42.9\%) && (\text{NLO, rEFT})\\
&-\;2.05\,\text{pb} && (-4.2\%) && ((t,b,c),\ \text{exact NLO})\\
&+\;9.56\,\text{pb} && (+19.7\%) && (\text{NNLO, rEFT})\\
&+\;0.34\,\text{pb} && (+0.7\%) && (\text{NNLO},\ 1/m_t)\\
&+\;2.40\,\text{pb} && (+4.9\%) && (\text{EW, QCD-EW})\\
&+\;1.49\,\text{pb} && (+3.1\%) && (\text{N}^3\text{LO, rEFT})
\end{aligned}$$

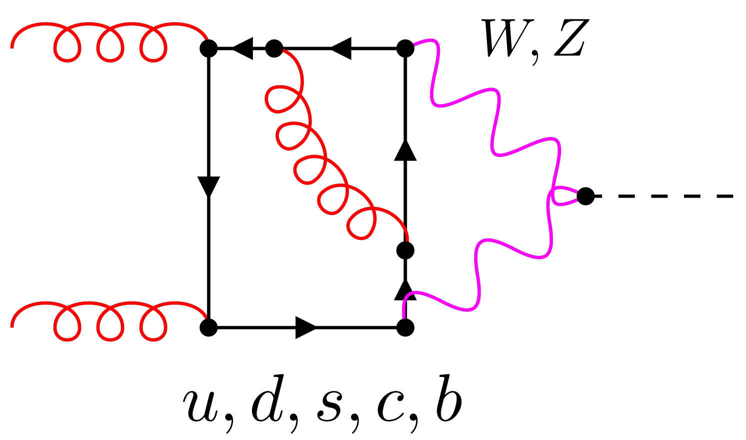
Anastasiou, Duhr, Dulat, Furlan, Gehrmann, Herzog, Lazopoulos, Mistlberger

δ(scale)	δ(trunc)	δ(PDF-TH)	δ(EW)	$\delta(t,b,c)$	$\delta(1/m_t)$
+0.10 pb −1.15 pb	±0.18 pb	±0.56 pb	±0.49 pb	±0.40 pb	±0.49 pb
+0.21% −2.37%	±0.37%	±1.16%	±1%	±0.83%	±1%

Mistlberger, Bonetti, Tancredi, K.M., Becchetti, Bonciani, Del Duca, Hirschi, Moriello, Czakon, Eschment, Schellenberger , Niggetiedt, Poncelet







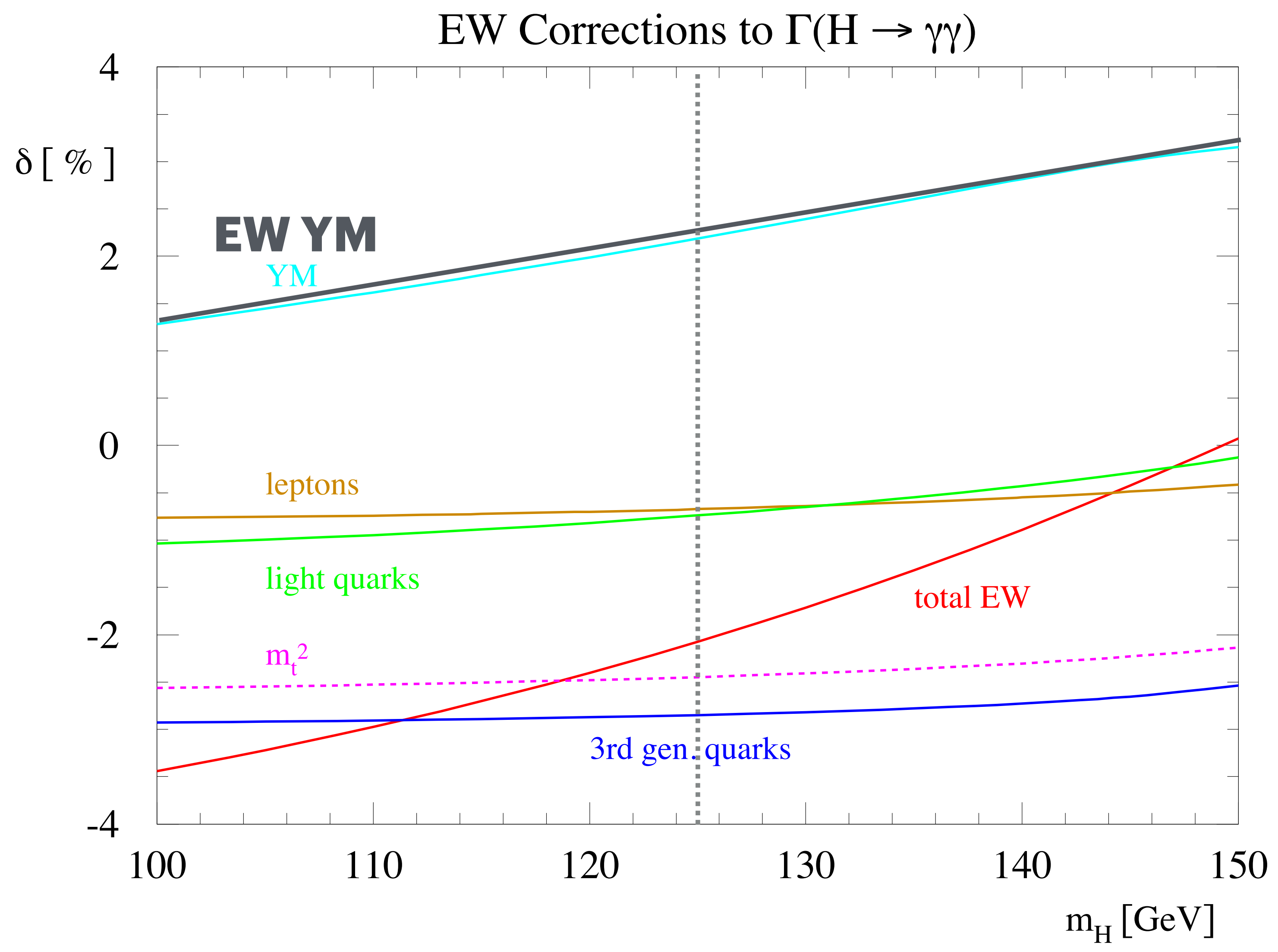
THE HIGHEST PRECISION FROM CROSS SECTION RATIOS?

One can reduce theoretical uncertainties by considering ratios of cross sections and other observables since common uncertainties may cancel out. The usual problem with ratios is that it is unclear to what extent the good things keep happening in fiducial regions. However, computations for fiducial, realistic cross sections have come a long way, so probably one should take advantage of this.

$$\frac{\sigma(gg \to H \to \gamma\gamma)}{\sigma(gg \to H \to ZZ^* \to 4l)} = \frac{\Gamma(H \to \gamma\gamma)}{\Gamma(H \to ZZ^* \to 4l)}$$

These widths are affected by QCD and EW radiative corrections. QCD corrections are tiny and are known to very high orders. It would be interesting to "observe" (highly-nontrivial) electroweak corrections to be in agreement with the SM.

$$\frac{\Gamma(H \to \gamma\gamma)}{\Gamma(H \to ZZ^* \to 4l)} \approx R_{\mathrm{LO}}\left(1 + 0.02_{\mathrm{QCD}} - 0.04_{\mathrm{EW}}\right)$$



Degrassi, Maltoni

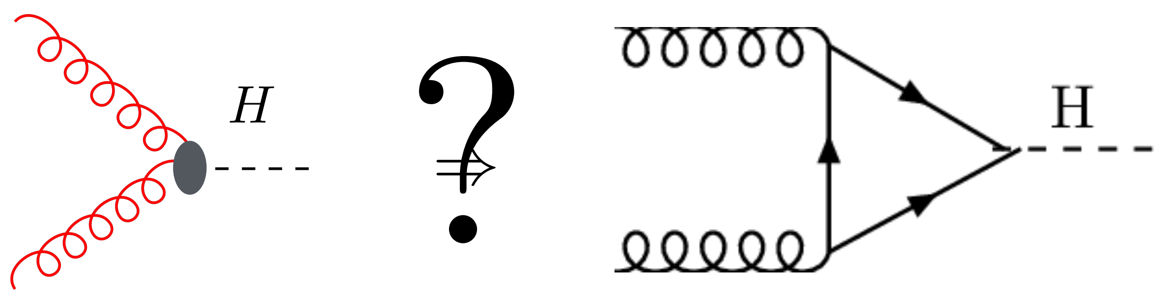
POI	Scenario	$\Delta_{\mathrm{tot}}/\sigma_{\mathrm{SM}}$	$\Delta_{\mathrm{stat}}/\sigma_{\mathrm{SM}}$	$\Delta_{\mathrm{exp}}/\sigma_{\mathrm{SM}}$	$\Delta_{\mathrm{sig}}/\sigma_{\mathrm{SM}}$	$\Delta_{\mathrm{bkg}}/\sigma_{\mathrm{SM}}$
$\sigma_{\mathrm{ggF}}^{ZZ}$	HL-LHC S1	$^{+0.044}_{-0.044}$	$^{+0.016}_{-0.016}$	$^{+0.031}_{-0.034}$	$^{+0.019}_{-0.017}$	$^{+0.018}_{-0.016}$
	HL-LHC S2	$^{+0.034}_{-0.034}$	$^{+0.016}_{-0.016}$	$^{+0.027}_{-0.027}$	$^{+0.010}_{-0.009}$	$^{+0.010}_{-0.009}$
$B_{\gamma\gamma}/B_{ZZ}$	HL-LHC S1	$^{+0.061}_{-0.057}$	$^{+0.020}_{-0.019}$	$^{+0.053}_{-0.049}$	$^{+0.018}_{-0.017}$	$^{+0.016}_{-0.014}$
	HL-LHC S2	$^{+0.045}_{-0.042}$	$^{+0.020}_{-0.019}$	$^{+0.037}_{-0.035}$	$^{+0.011}_{-0.011}$	$^{+0.010}_{-0.009}$
B_{WW}/B_{ZZ}	HL-LHC S1	$^{+0.065}_{-0.061}$	$^{+0.019}_{-0.018}$	$^{+0.042}_{-0.038}$	$^{+0.036}_{-0.034}$	$^{+0.028}_{-0.027}$
	HL-LHC S2	$^{+0.049}_{-0.047}$	$^{+0.019}_{-0.018}$	$^{+0.036}_{-0.034}$	$^{+0.020}_{-0.018}$	$^{+0.019}_{-0.018}$
$B_{\tau\tau}/B_{ZZ}$	HL-LHC S1	$^{+0.066}_{-0.062}$	$^{+0.024}_{-0.024}$	$^{+0.043}_{-0.038}$	$^{+0.033}_{-0.033}$	$^{+0.029}_{-0.026}$
	HL-LHC S2	$^{+0.053}_{-0.050}$	$^{+0.024}_{-0.024}$	$^{+0.037}_{-0.035}$	$^{+0.023}_{-0.022}$	$^{+0.019}_{-0.017}$
B_{bb}/B_{ZZ}	HL-LHC S1	$^{+0.118}_{-0.105}$	$^{+0.038}_{-0.037}$	$^{+0.053}_{-0.048}$	$^{+0.058}_{-0.052}$	$^{+0.080}_{-0.069}$
	HL-LHC S2	$^{+0.092}_{-0.084}$	$^{+0.038}_{-0.037}$	$^{+0.046}_{-0.043}$	$^{+0.036}_{-0.032}$	$^{+0.061}_{-0.054}$

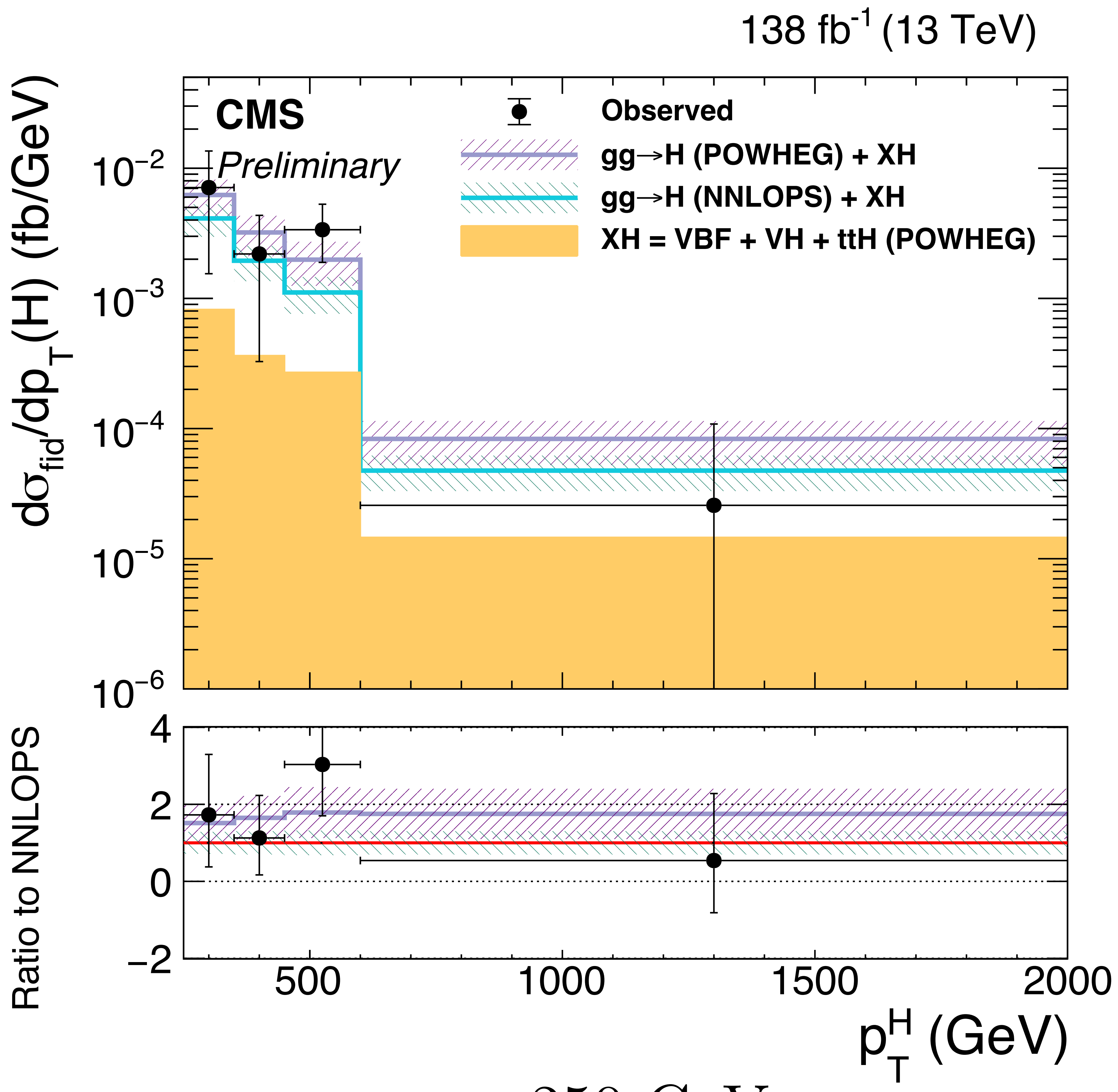
ATLAS HL projections on branchings

A DIFFERENT TAKE ON PRECISION

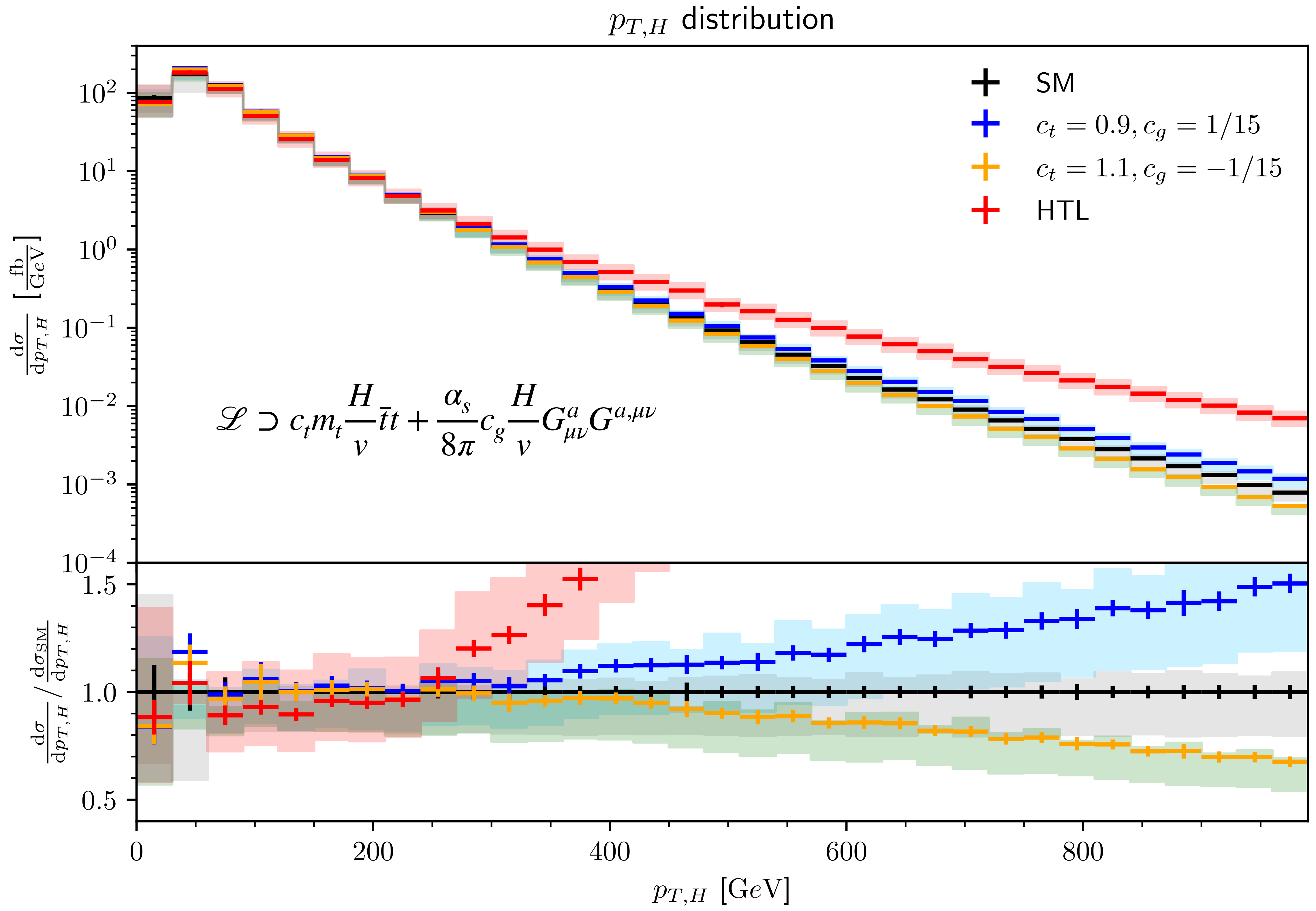
Measurements of Higgs production at high transverse momentum are very interesting from the BSM/EFT viewpoint. Such measurements are still statistically limited but we do not see very large deviations which tells us that the Higgs is indeed produced through a top quark loop, without substantial ultra-short-range component.

$$d\sigma_{\rm fid} = 1.96^{+0.86}_{-0.69}\ {\rm fb} \qquad d\sigma_{\rm fid,SM} = 1.2 \pm 0.2\ {\rm fb}$$





$$p_{\perp,H} > 250\ {\rm GeV}$$



Campillo Aleveria, Heinrich, Kerner, Kunz

SUMMARY

Perturbative QCD is a well-developed theory whose role, in the context of the LHC physics, is to facilitate interpretation of experimental results in terms of parameters that appear in the Lagrangian of the SM or its extensions.

Continuous methodological progress in perturbative QCD allows us to describe collider processes of ever increasing complexity with higher and higher precision.

State-of-the-art calculations at next-to-leading and next-to-next-to-leading orders in perturbative QCD remain very challenging, but are becoming more and more manageable. The focus is slowly shifting towards the next perturbative order, N3LO.

These impressive successes of the perturbative approach to hadron collisions, emphasize the need of a systematic understanding of non-perturbative power corrections at hadron colliders. Without it, further meaningful improvements in ultra-precise determinations of physical parameters (the top quarks mass, the strong coupling constant etc.) may not be possible, in spite of being statistically achievable.

Continued progress with perturbative computations in QCD, as well as a better understanding of non-perturbative effects that hopefully can be achieved, will allow for many exciting physics studies at high-luminosity LHC, that range from ultra-high precision measurements of fundamental parameters of the SM to detailed explorations of physics at high transverse momenta and invariant masses.